Q: Why do the Earth & planets have ionospheres? magnetospheres?

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w/ liberal "borrowing" from Fuller-Rowel, Solomon, Sojka, Lean, Vasylinunas, Bagenal, Luhman

Heliophysics chain

Q: Why do the Earth & planets have ionospheres?

A: Because of the **Sun's** corona (its EUV & X-rays)

Q: Why does the Sun have a corona?

A: Because of its magnetic field (and its heating)

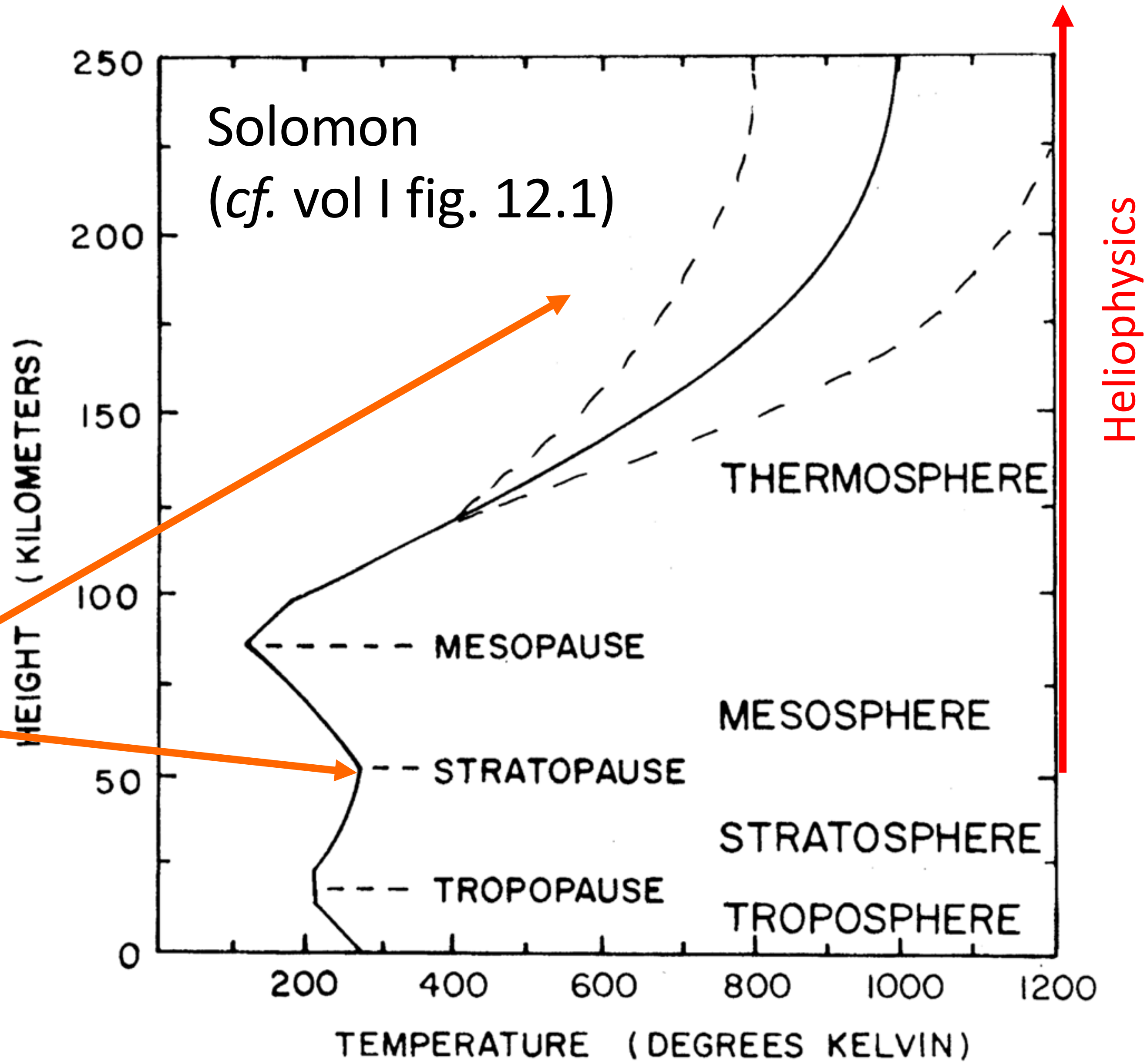
Q: Why does the Sun have a magnetic field?

A: Because of its dynamo

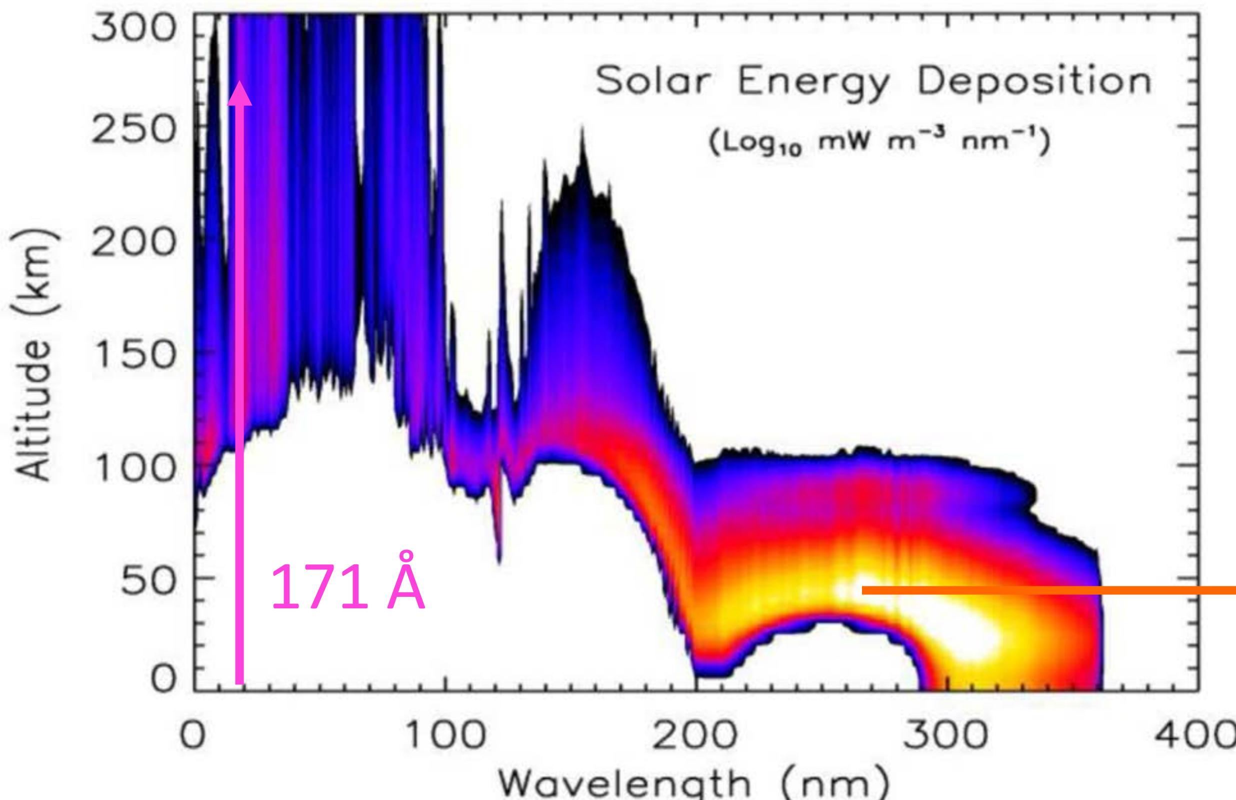
Earth's neutral atmosphere

	Earth
N_2	77%
O_2	21%
H_2O	1%
CO_2	0.03%

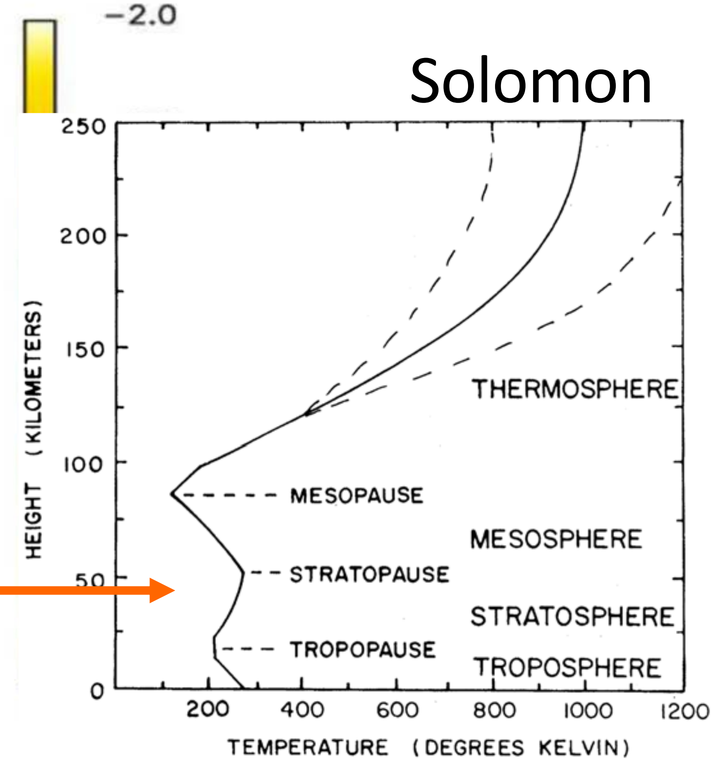
??

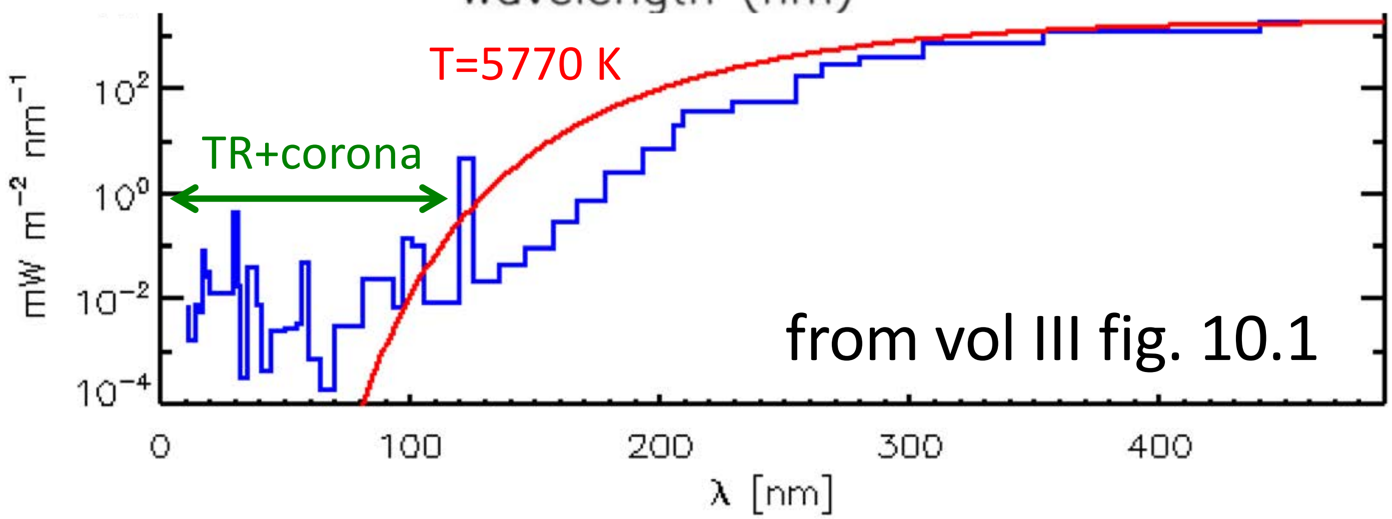


vol. III fig. 13.3



Solomon





from vol III fig. 10.1

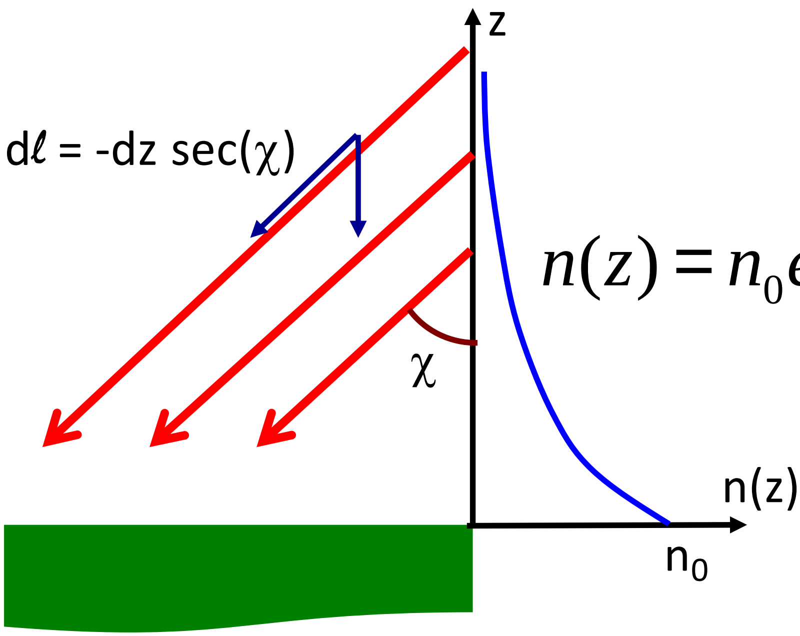
Fate of a photon

w/ absorption x-section σ

optical path $\tau(x)$ = avg. # absorbers in cylinder w/ x-section σ

Prob. of survival: $$P(x) = \exp\left[-\int^{x} \sigma n(\ell)\, d\ell \right]$$

$\tau=1$ ➔ 1 absorber: mean-free path



$$n(z) = n_0 e^{-z/H}$$

$$\tau(z) = \int_z^{\infty} \sigma n(z') \sec(\chi)\, dz'$$

$$= \sigma n_0 \sec(\chi) \int_z^{\infty} e^{-z'/H}\, dz'$$

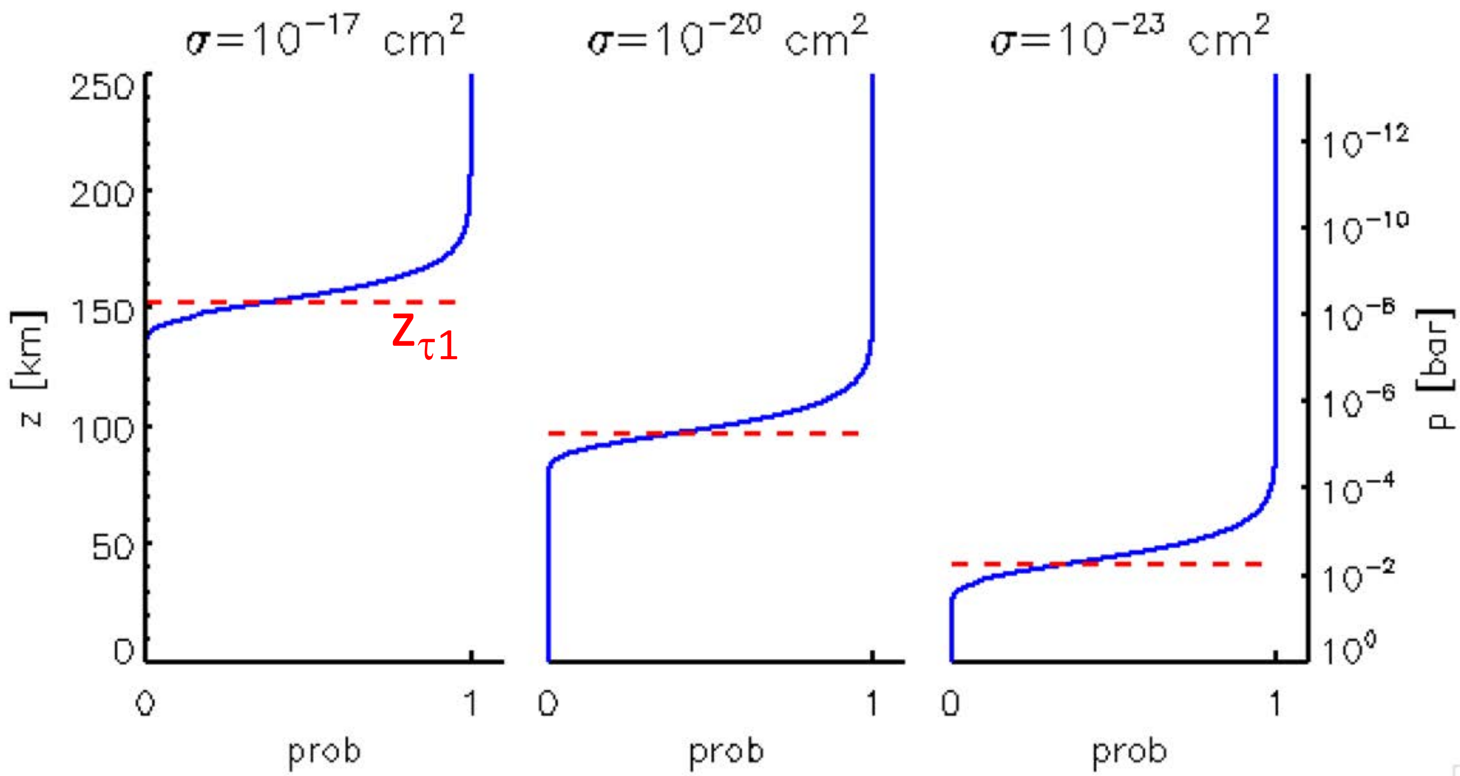
$$\tau(z) = \sigma n_0 H \sec(\chi)\, e^{-z/H} = e^{-(z-z_{\tau 1})/H}$$

height of $\tau=1$: $$z_{\tau 1} = H \ln\left[\sigma n_0 H \sec(\chi)\right]$$

Prob. of survival: $$P(z) = e^{-\tau(z)} = \exp\left[-e^{-(z-z_{\tau 1})/H} \right]$$

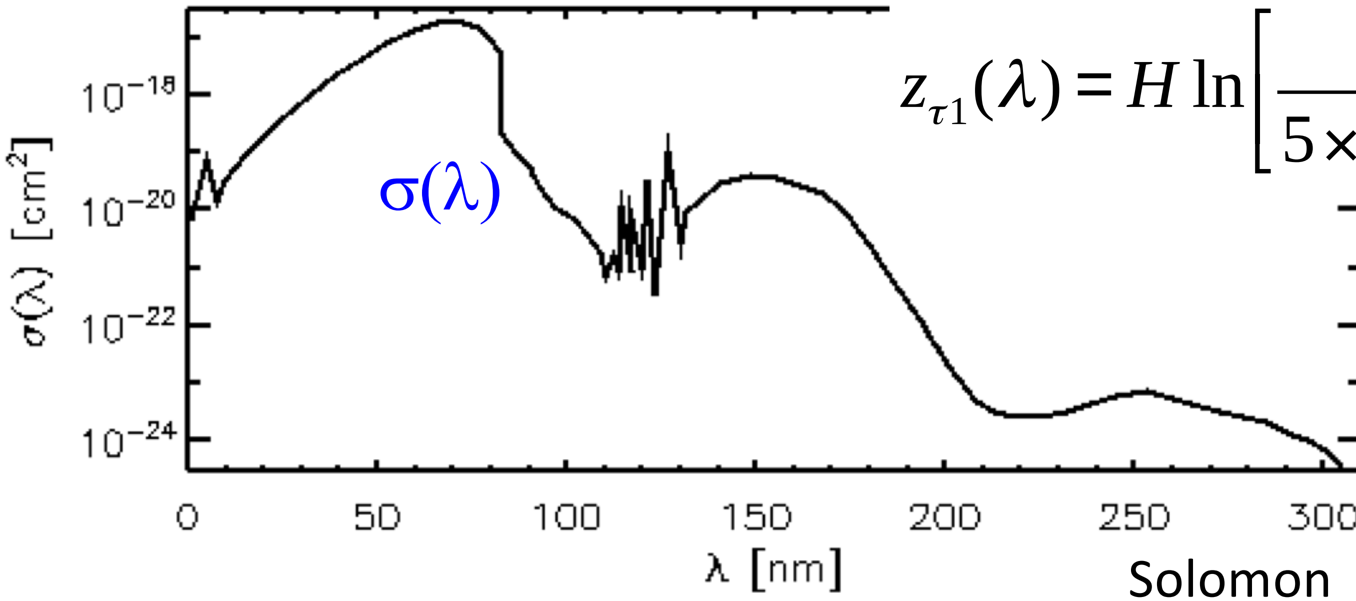
$$e^{z_{\tau 1}/H} = \sigma n_0 H \sec(\chi) = \frac{\sigma n_0 kT}{\bar{m}g}\sec(\chi) = \sigma\frac{p_0}{\bar{m}g}\sec(\chi) = \frac{\sigma}{\sigma_0}\sec(\chi)$$

$$\sigma_0 = \frac{\bar{m}g}{p_0} = \frac{5\times 10^{-23}\,\mathrm{g}\cdot 980\,\mathrm{cm/s^2}}{10^6\,\mathrm{erg/cm^3}} = 5\times 10^{-26}\,\mathrm{cm^2}$$

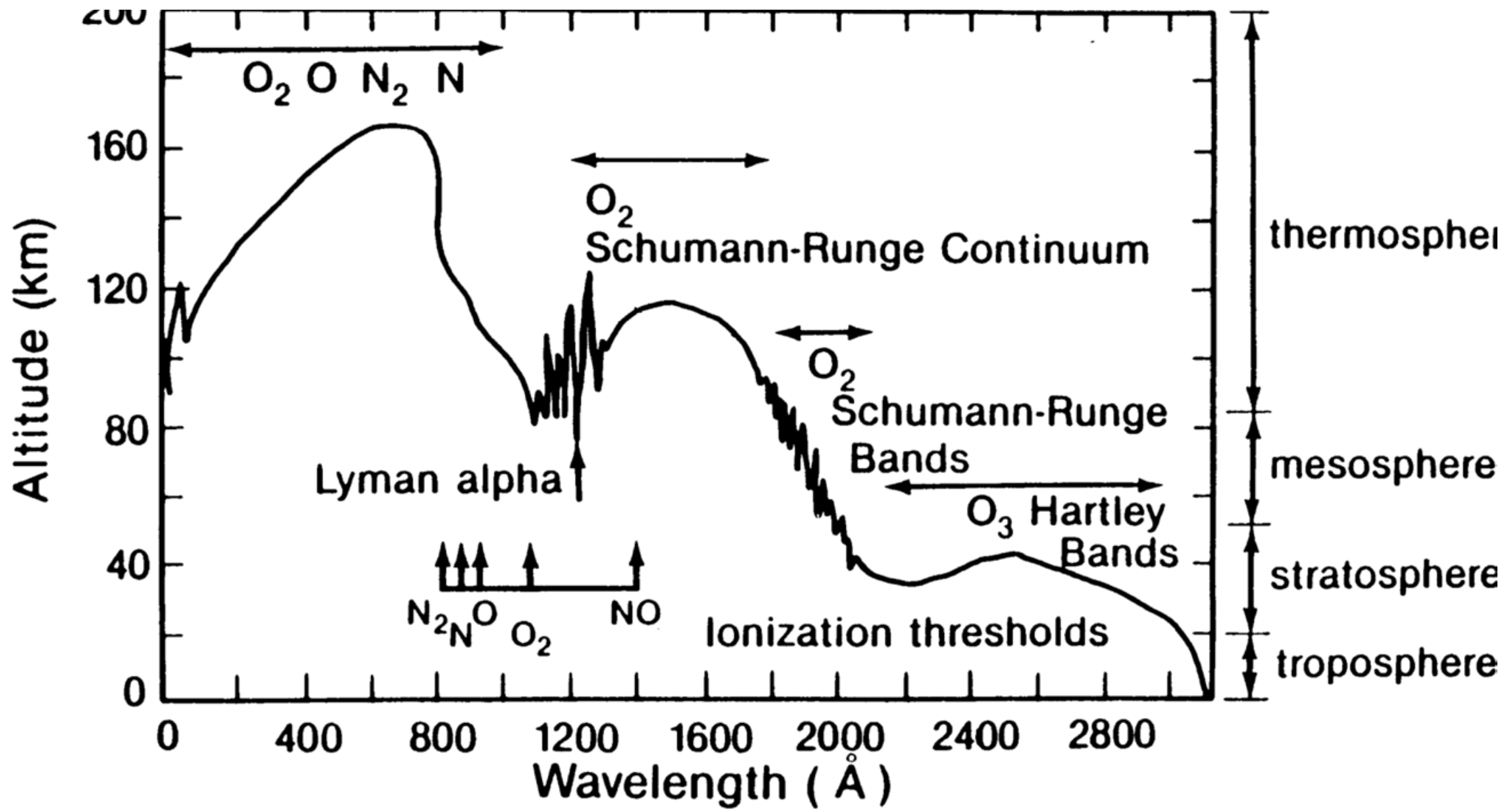


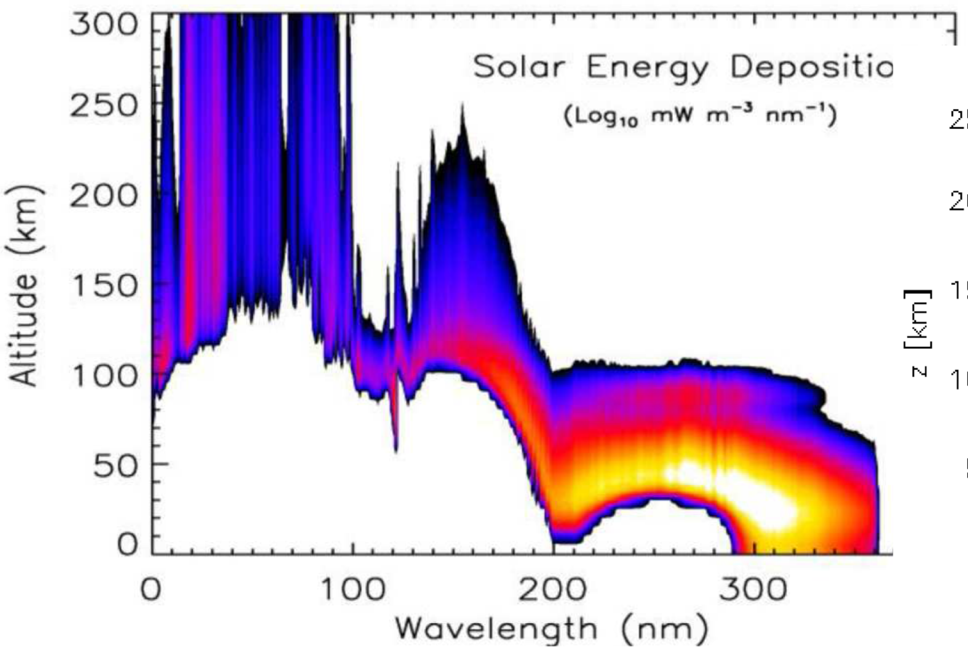
$$P(z) = e^{-\tau(z)} = \exp\left[-e^{-(z-z_{\tau 1})/H}\right]$$

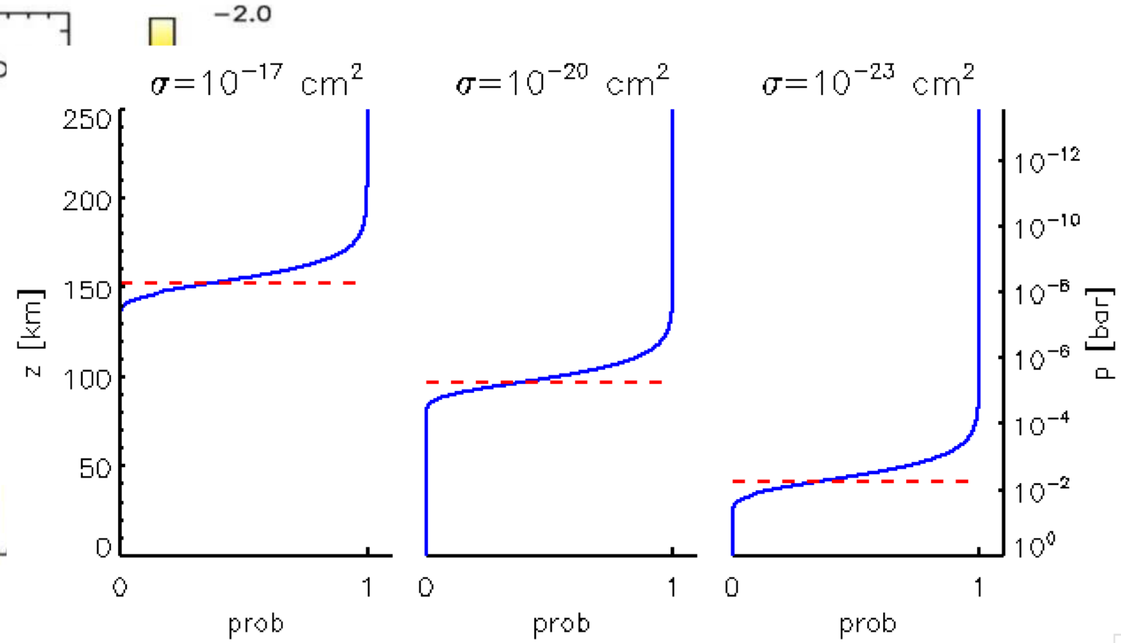
$$z_{\tau 1}(\lambda) = H \ln\left[\frac{\sigma(\lambda)}{5\times 10^{-26}\,\mathrm{cm}^2}\right]$$

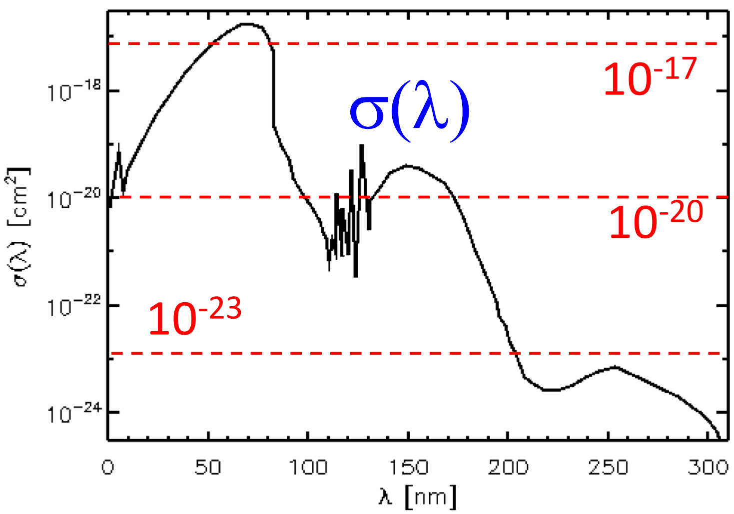


Solomon









$$P[z(\lambda)] = \exp\left[-e^{-[z - z_{\tau 1}(\lambda)]/H} \right]$$

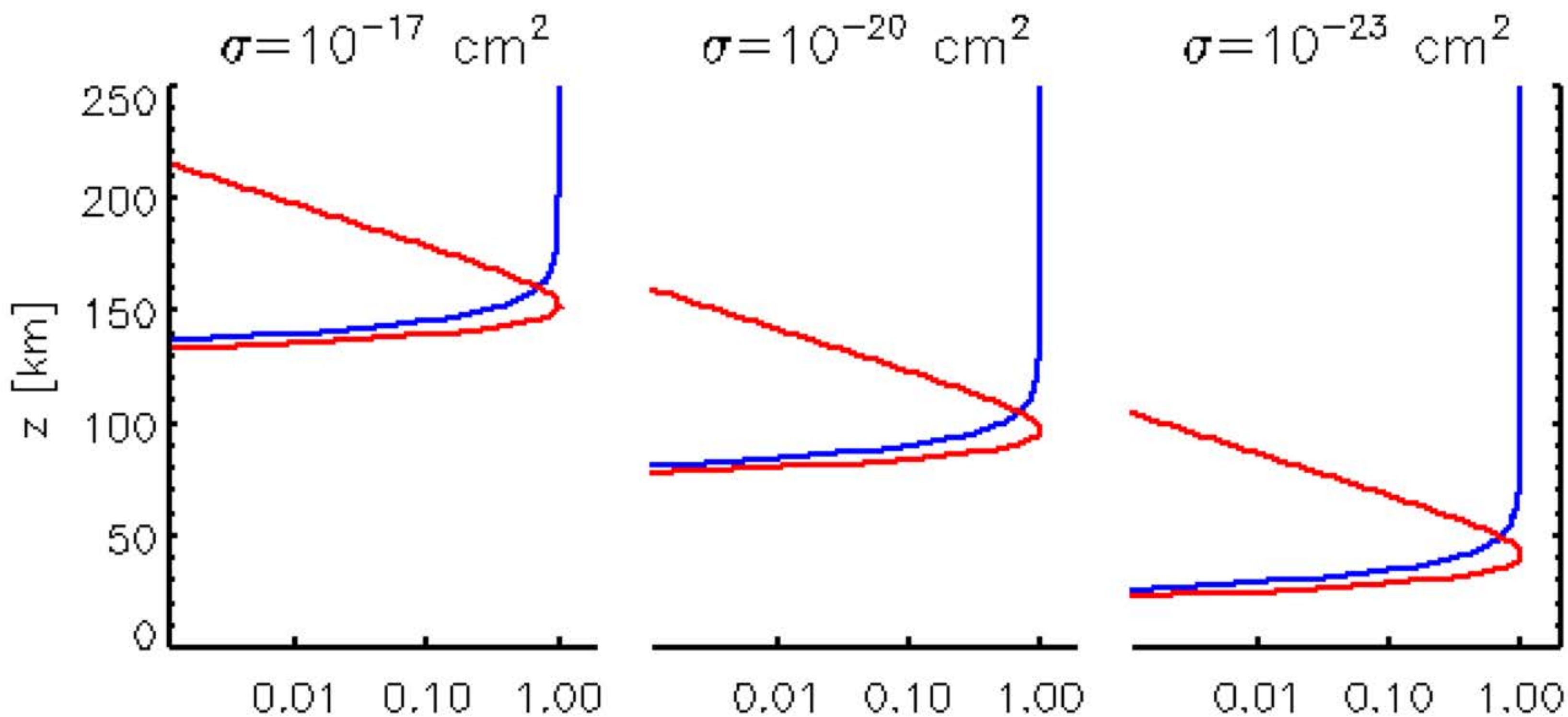
$$z_{\tau 1}(\lambda) = H \ln\left[\frac{\sigma(\lambda)}{5 \times 10^{-26}\,\mathrm{cm}^2} \right]$$

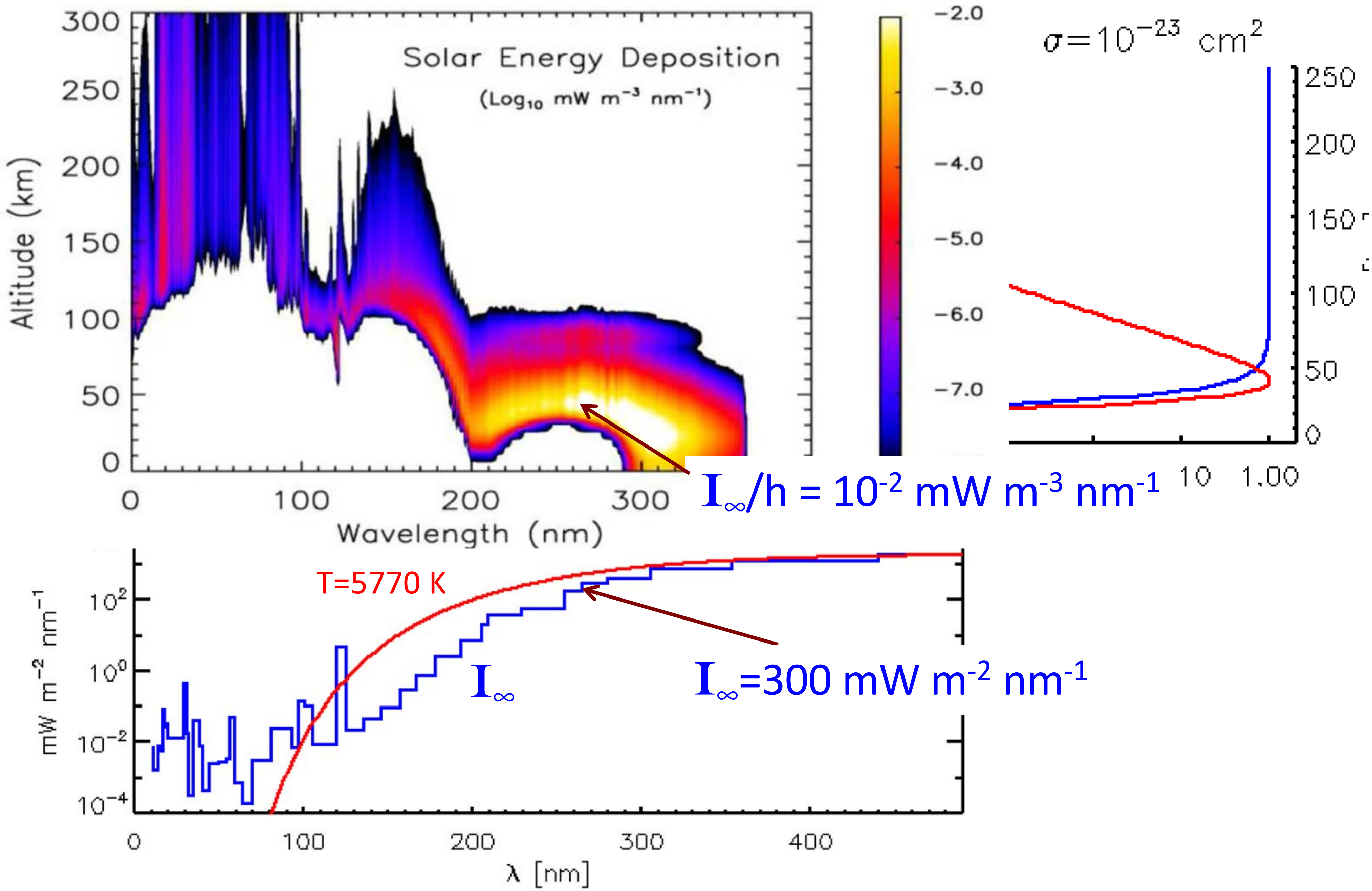
Radiation intensity & heating

Energy flux: $$I(z) = I_\infty P(z) = I_\infty \exp\left[-e^{-(z-z_{\tau 1})/H} \right]$$

Energy deposition: $$\frac{dI}{dz} = \frac{I_\infty}{H} \exp\left[-e^{-(z-z_{\tau 1})/H} - \frac{z-z_{\tau 1}}{H} \right]$$

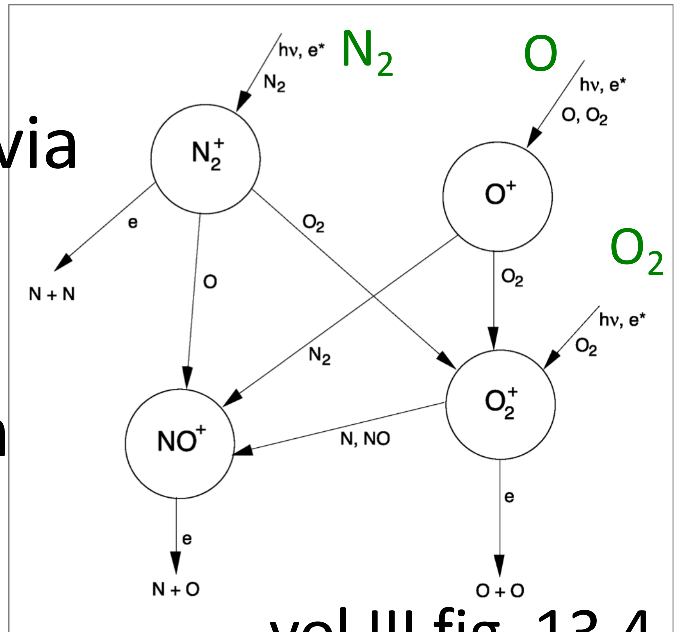
Chapman layer



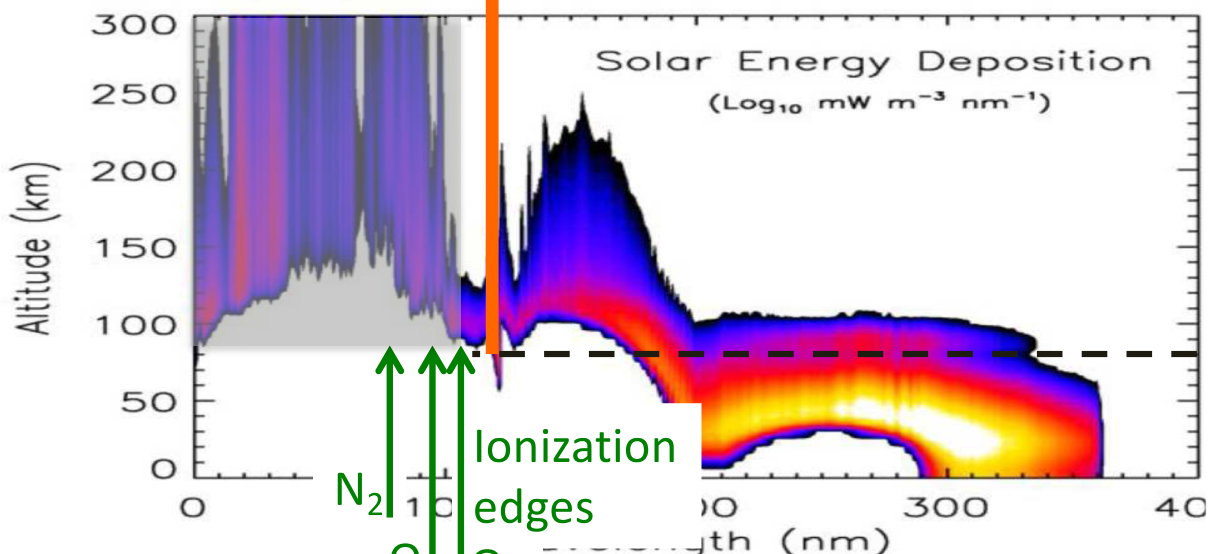


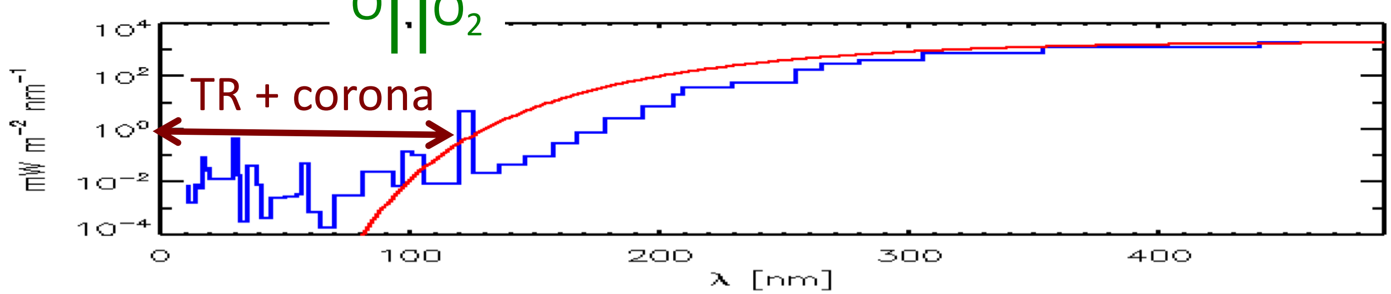
Solar Energy Deposition
(Log10 mW m-3 nm-1)
Altitude (km)
Wavelength (nm)
σ=10-23 cm2
Iinf/h = 10-2 mW m-3 nm-1
T=5770 K
Iinf
Iinf=300 mW m-2 nm-1
mW m-2 nm-1
λ [nm]

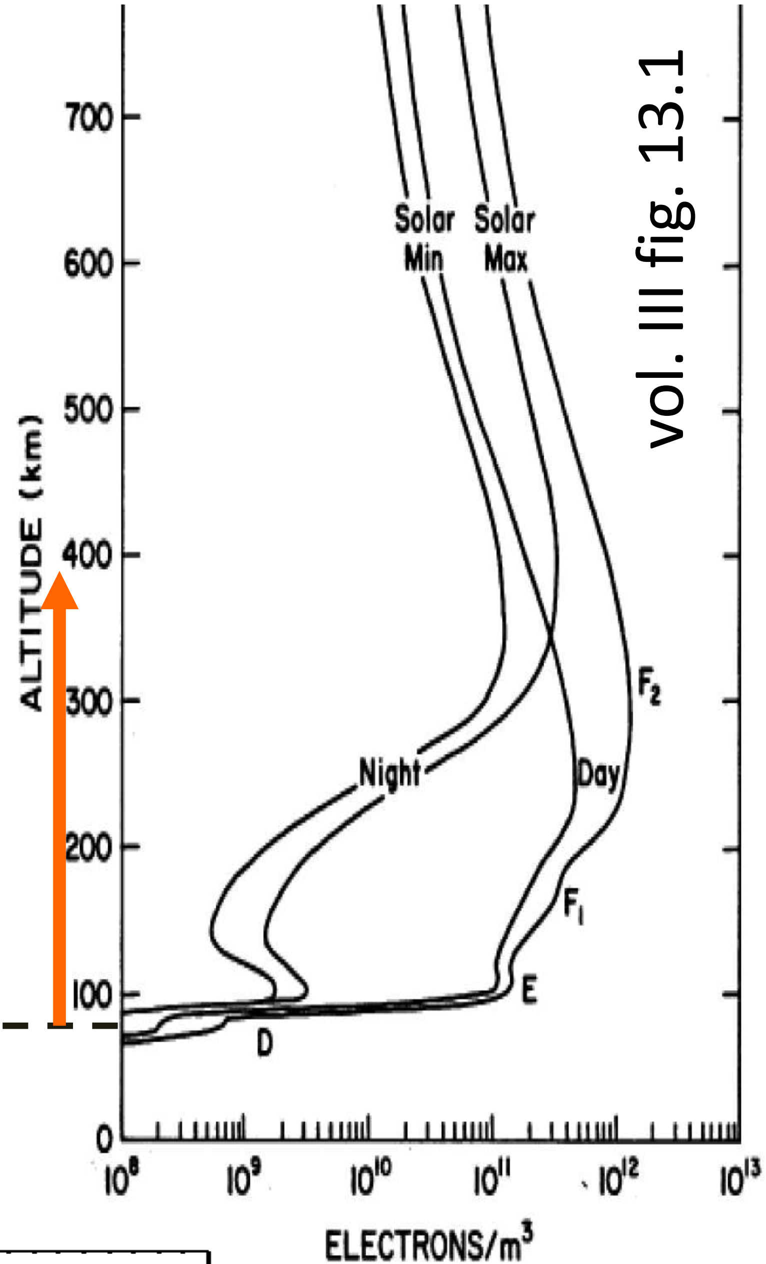
Absorption via ionization creates ion/electron pairs



vol III fig. 13.4







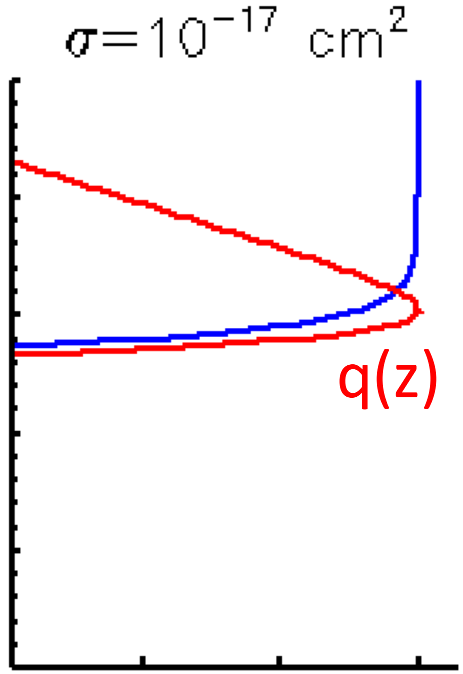
vol. III fig. 13.1

Rate of photo-ionization (per volume)

= Electron production rate:

$$q(z) = \sigma_{\text{ion}} n(z)\, F(z) = \sigma_{\text{ion}} n(z)\, F_\infty P(z)$$

$$= \sigma_{\text{ion}} n_0 F_\infty \exp\left[-e^{-(z-z_{\tau 1})/H} - \frac{z}{H} \right]$$



Electron destruction by recombination with +'ve ions @ rate

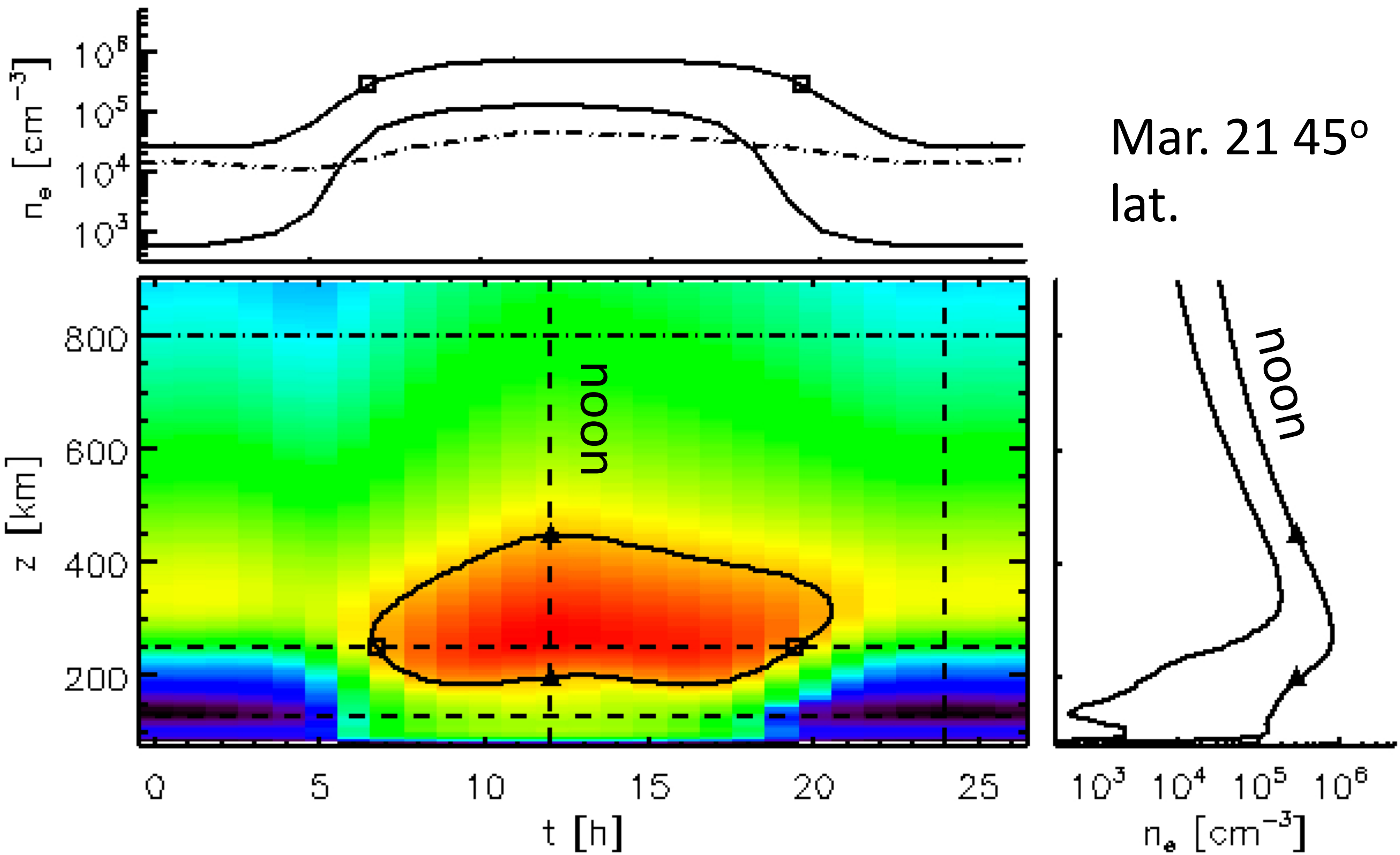
$$L = \alpha\, n_e n_i \approx \alpha\, n_e^2$$

Assuming neutrality

Production balances destruction: q=L

$$n_e(z) = \sqrt{\frac{q(z)}{\alpha(z)}}$$

Q: why does n_e vary more lowdown?

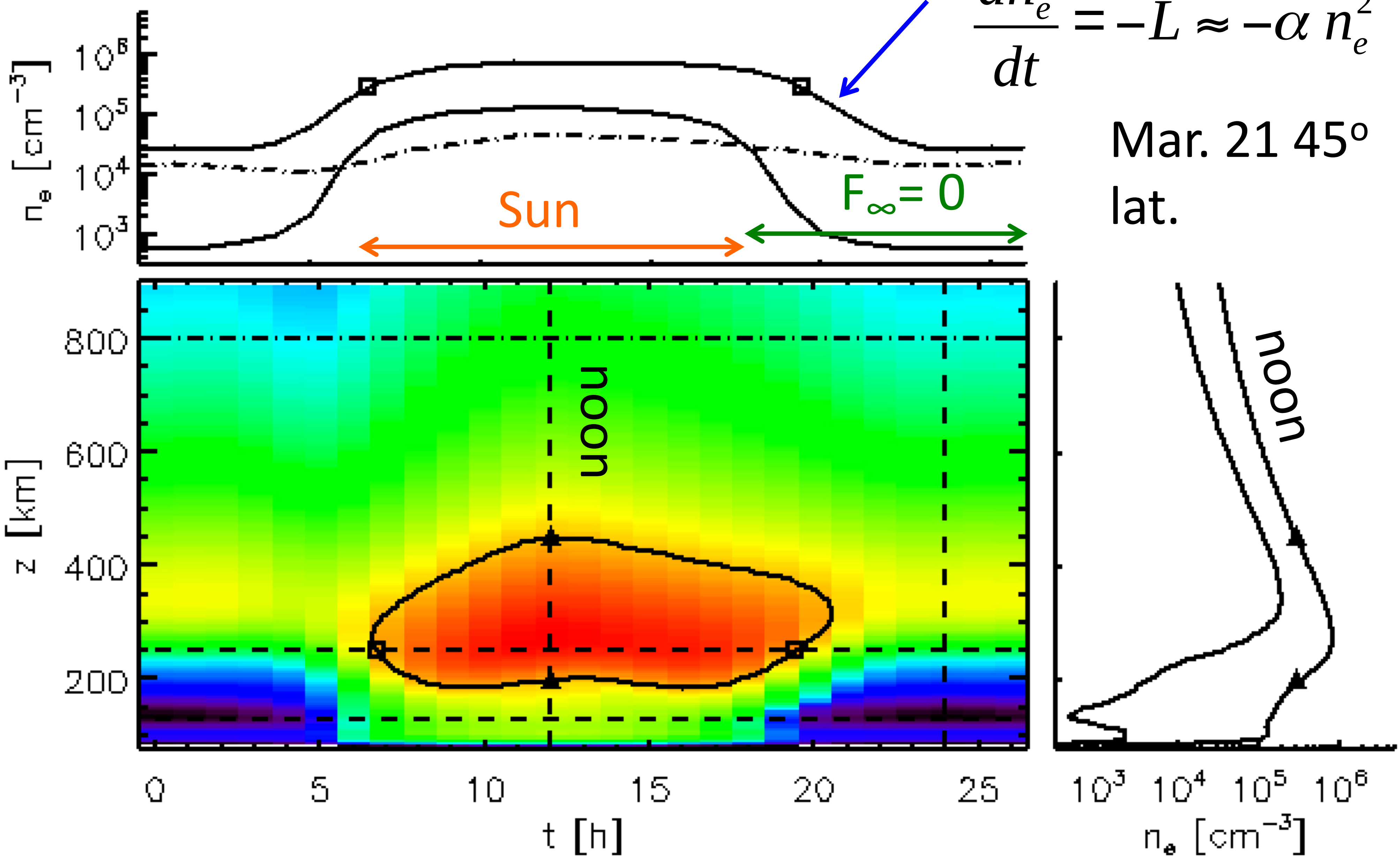


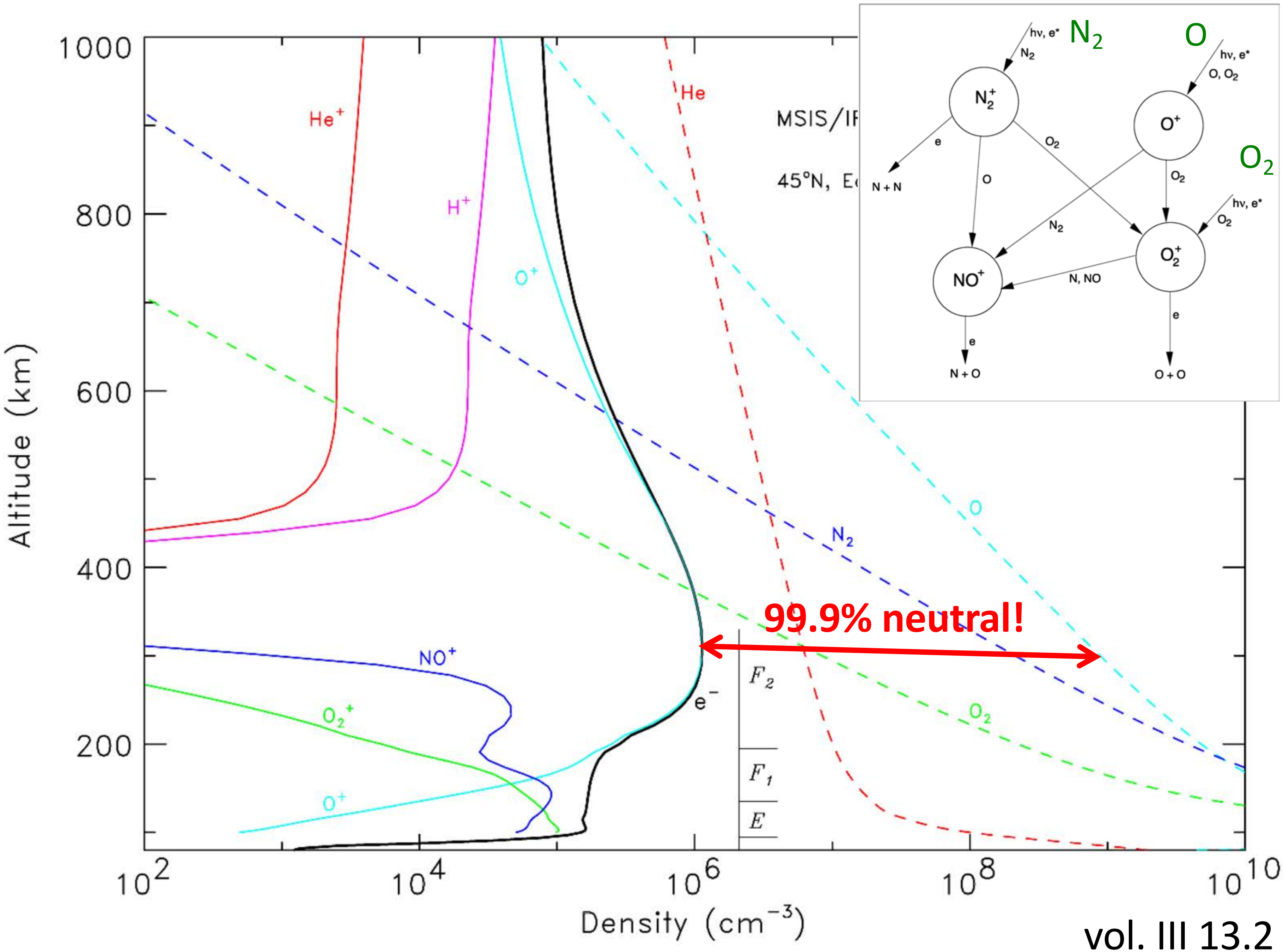
Mar. 21 45^o lat.

Production shut off
– recombination removes electrons

$$\frac{dn_e}{dt} = -L \approx -\alpha\, n_e^2$$

Mar. 21 45° lat.

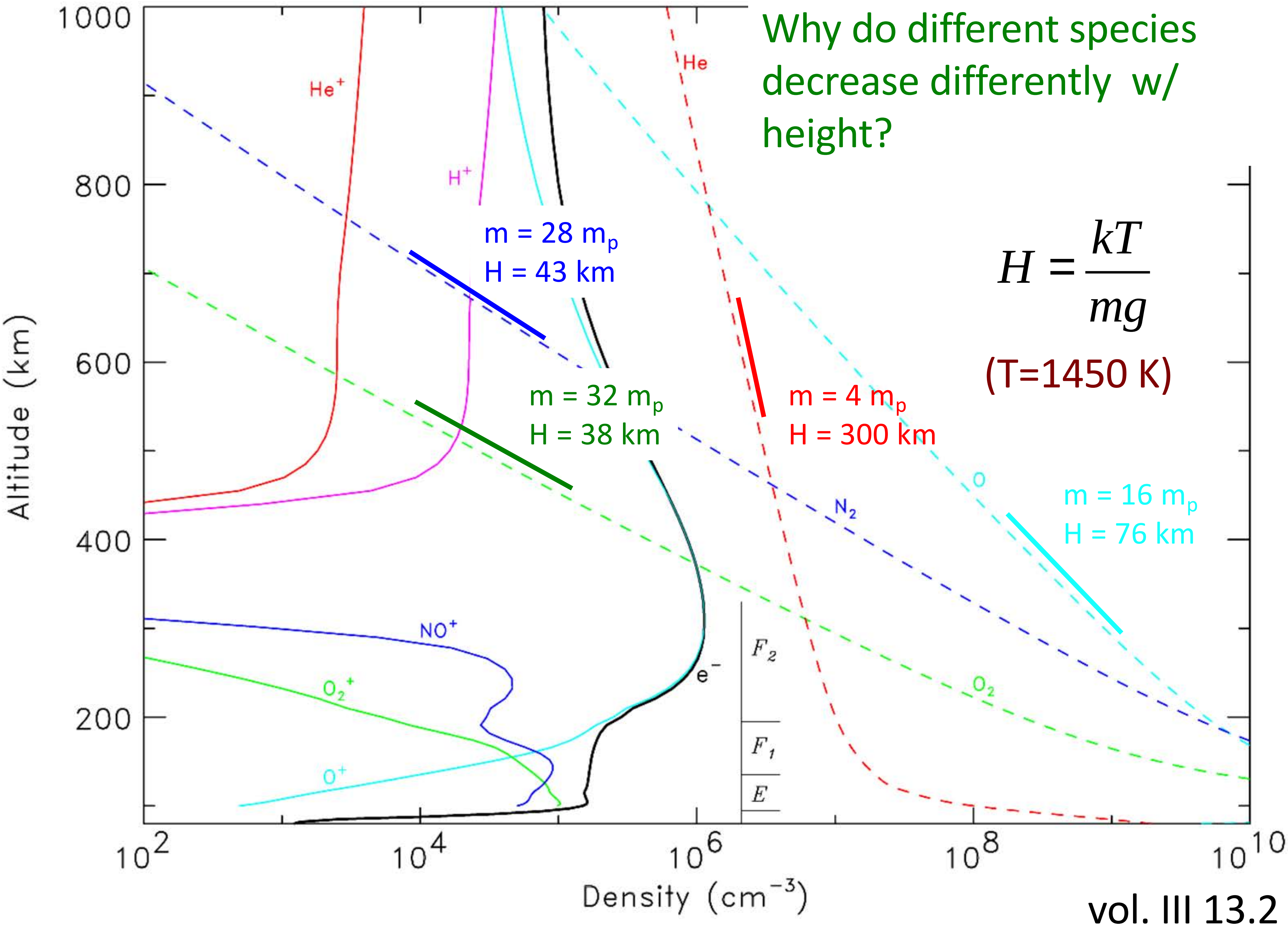




1000
800
600
400
200
Altitude (km)
10^2
10^4
10^6
10^8
10^{10}
Density (cm^{-3})
He^+
H^+
O^+
He
MSIS/IF
45°N, E
O
N_2
O_2
NO^+
O_2^+
O^+
e^-
F_2
F_1
E
99.9% neutral!
N_2
O
O_2
hν, e*
N_2^+
O^+
O_2^+
NO^+
N + N
N + O
O + O
O, O_2
N, NO
e

vol. III 13.2

Why do different species decrease differently w/ height?



$$H = \frac{kT}{mg}$$

(T=1450 K)

vol. III 13.2

Ionospheric plasma

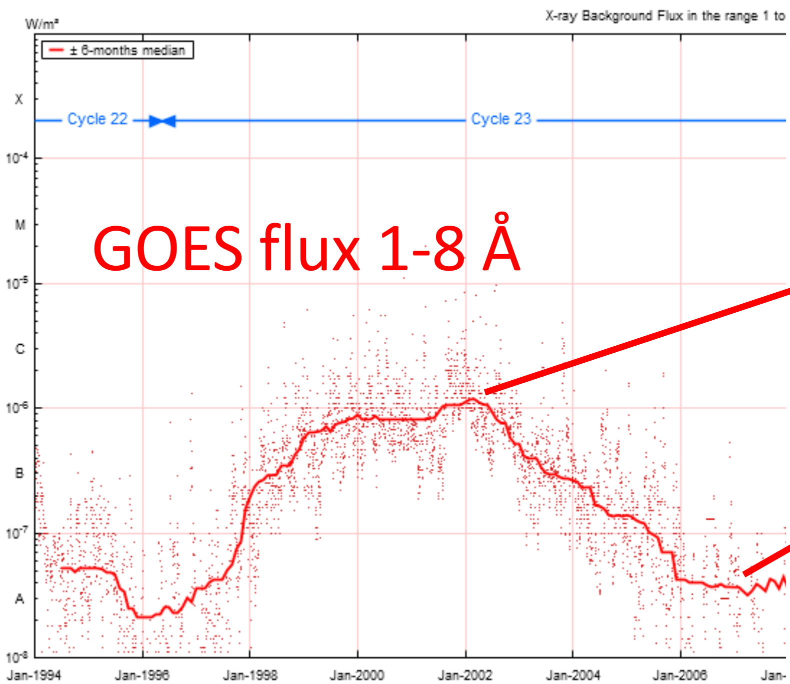
- ions/e^- form plasma – conducting fluid
- Neutrals: separate fluid
- Continual creation/destruction couples fluids – created ``drag force'' between them

A plasma with electron density n_e (cm^{-3}) screens out E fields w/ f < its plasma frequency

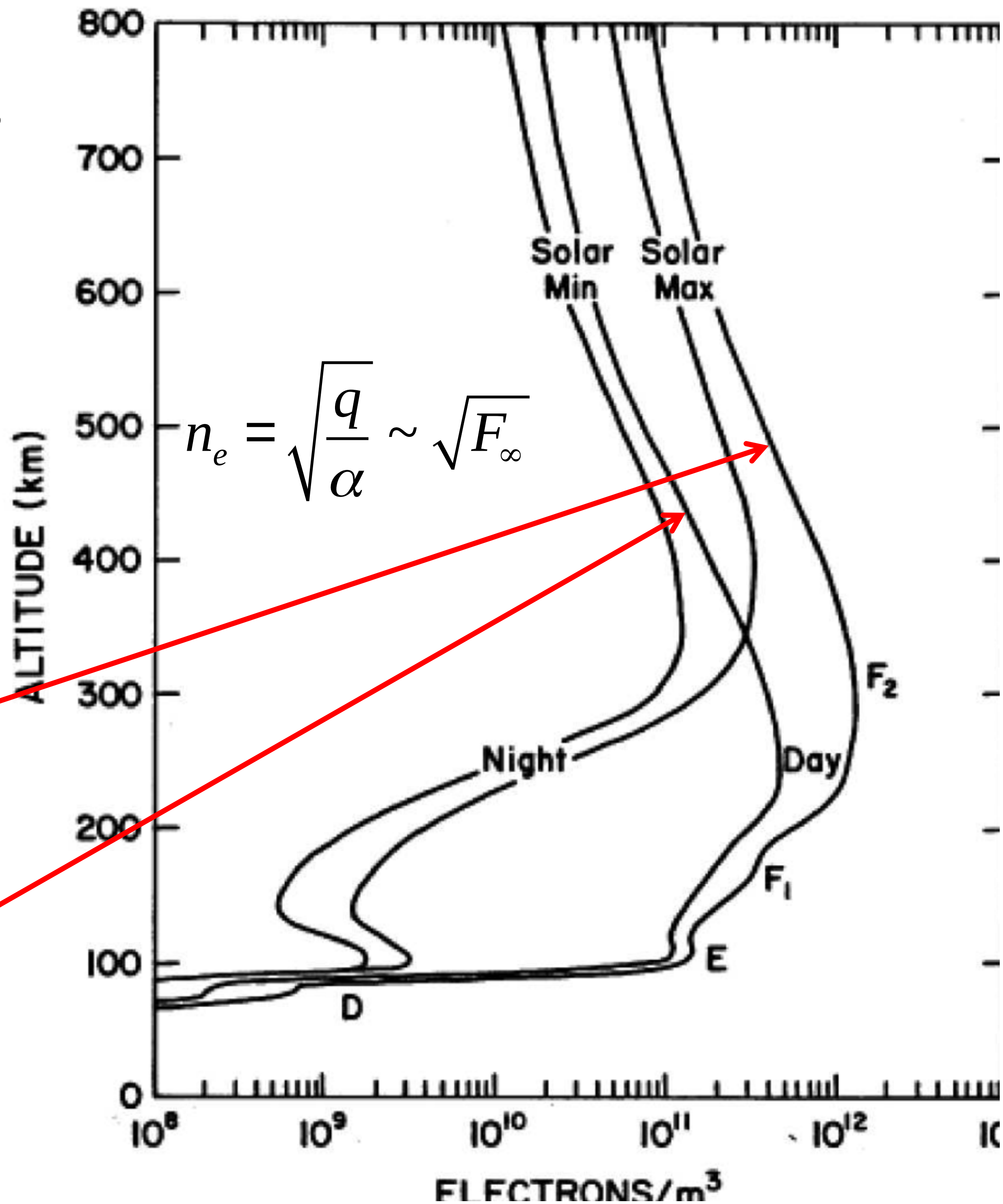
$$f_p = \sqrt{\frac{e^2 n_e}{\pi m_e}} = 10^4\,\mathrm{Hz}\; n_e^{1/2}$$

Q: what is the lowest freq. solar radio emission we can observe from the ground?

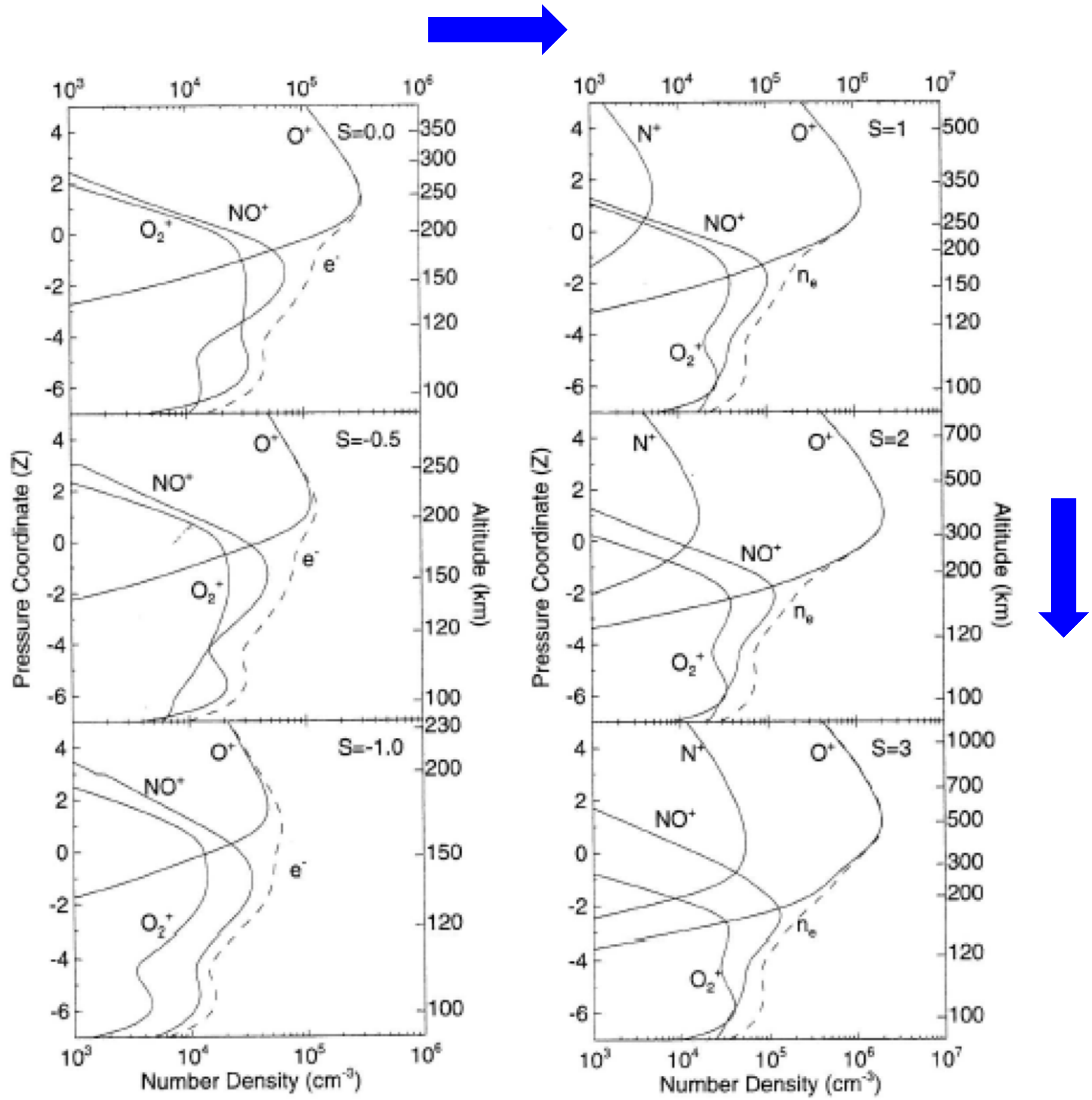
Corona varies – ionosphere varies



$$n_e = \sqrt{\frac{q}{\alpha}} \sim \sqrt{F_\infty}$$



increasing coronal flux

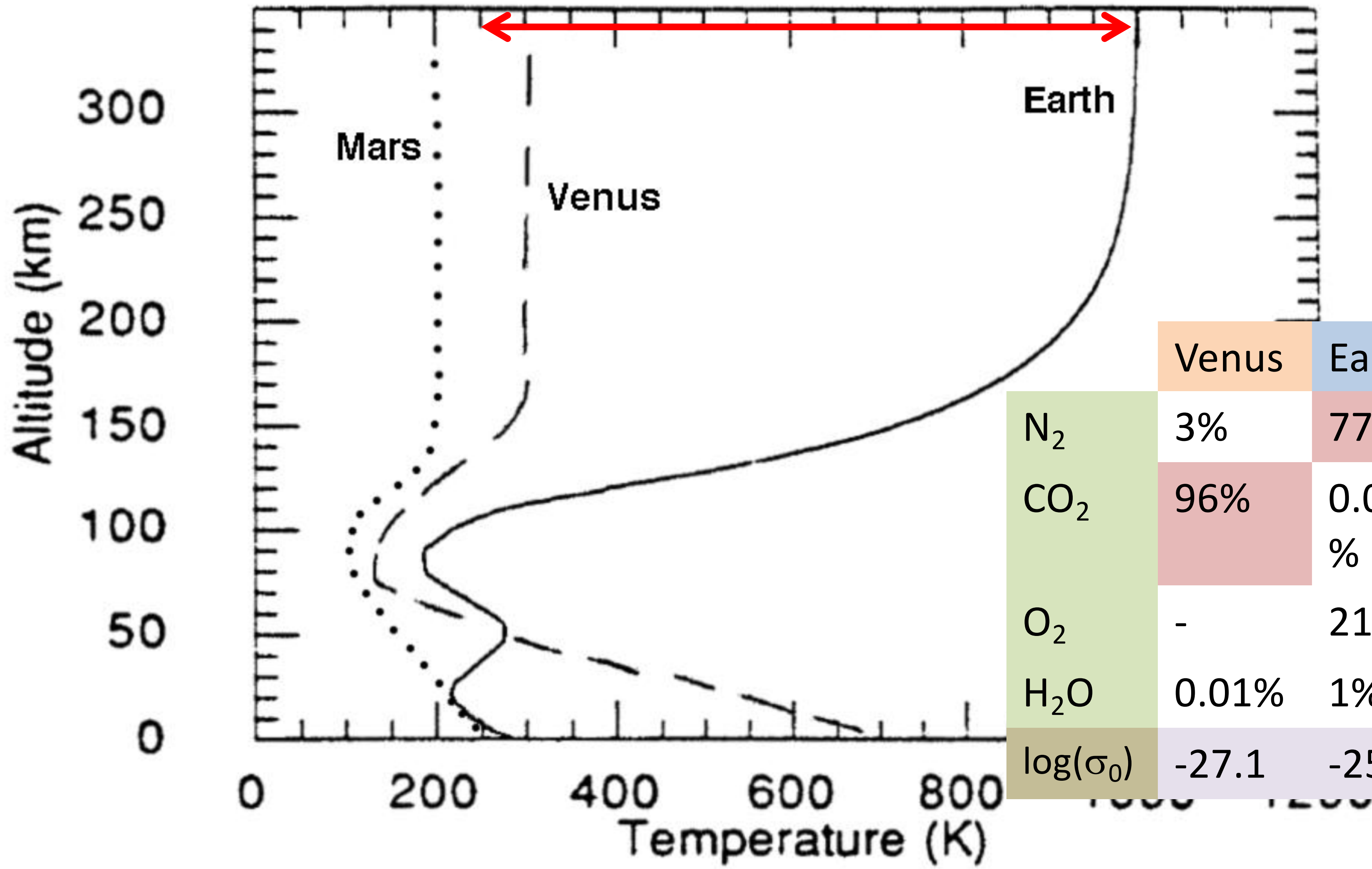


vol. III 14.4

Other planets… other atmospheres

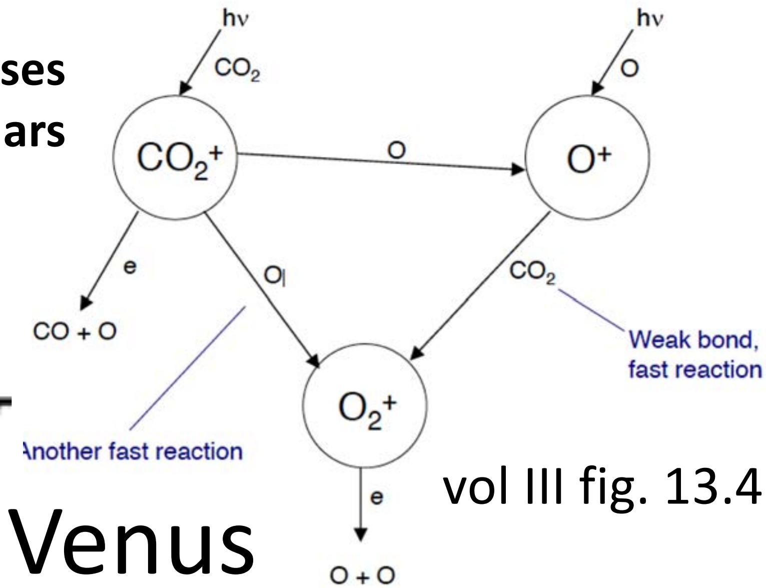
Why is Earth's thermosphere so hot?

Solomon



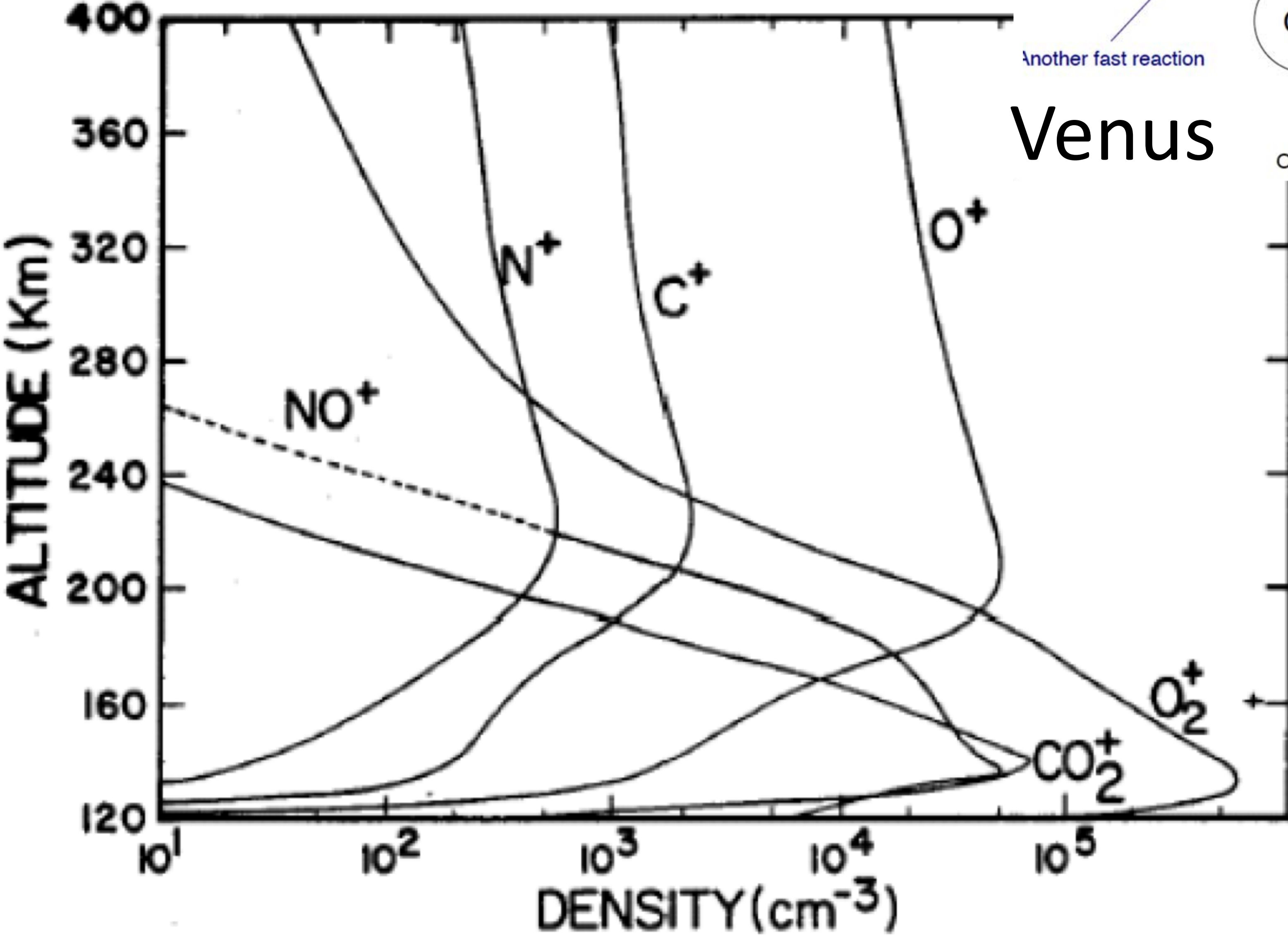
	Venus	Earth	Mars
N_2	3%	77%	3%
CO_2	96%	0.03 %	95%
O_2	-	21%	-
H_2O	0.01%	1%	-
$\log(\sigma_0)$	-27.1	-25.3	-23.4

Principal Ionization Processes on Venus & Mars



vol III fig. 13.4

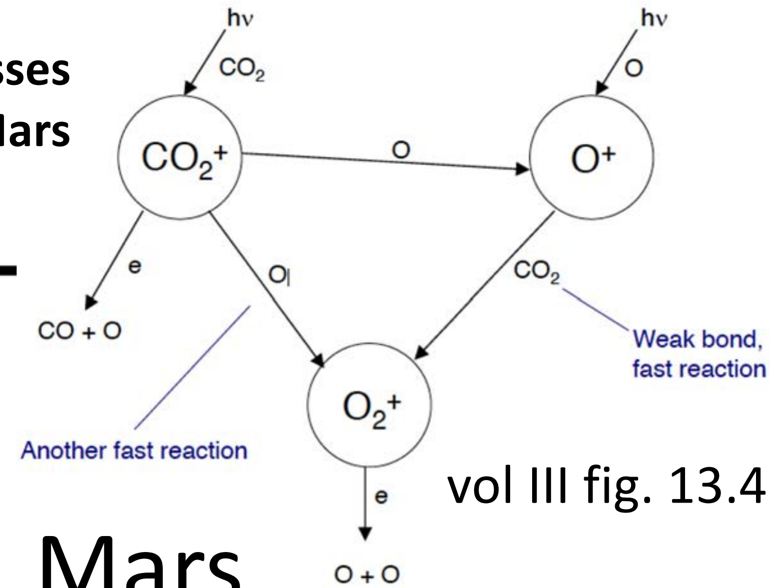
Solomon $\chi=60^o$



Venus

	Venus
N_2	3%
CO_2	96%
O_2	-
H_2O	0.01%

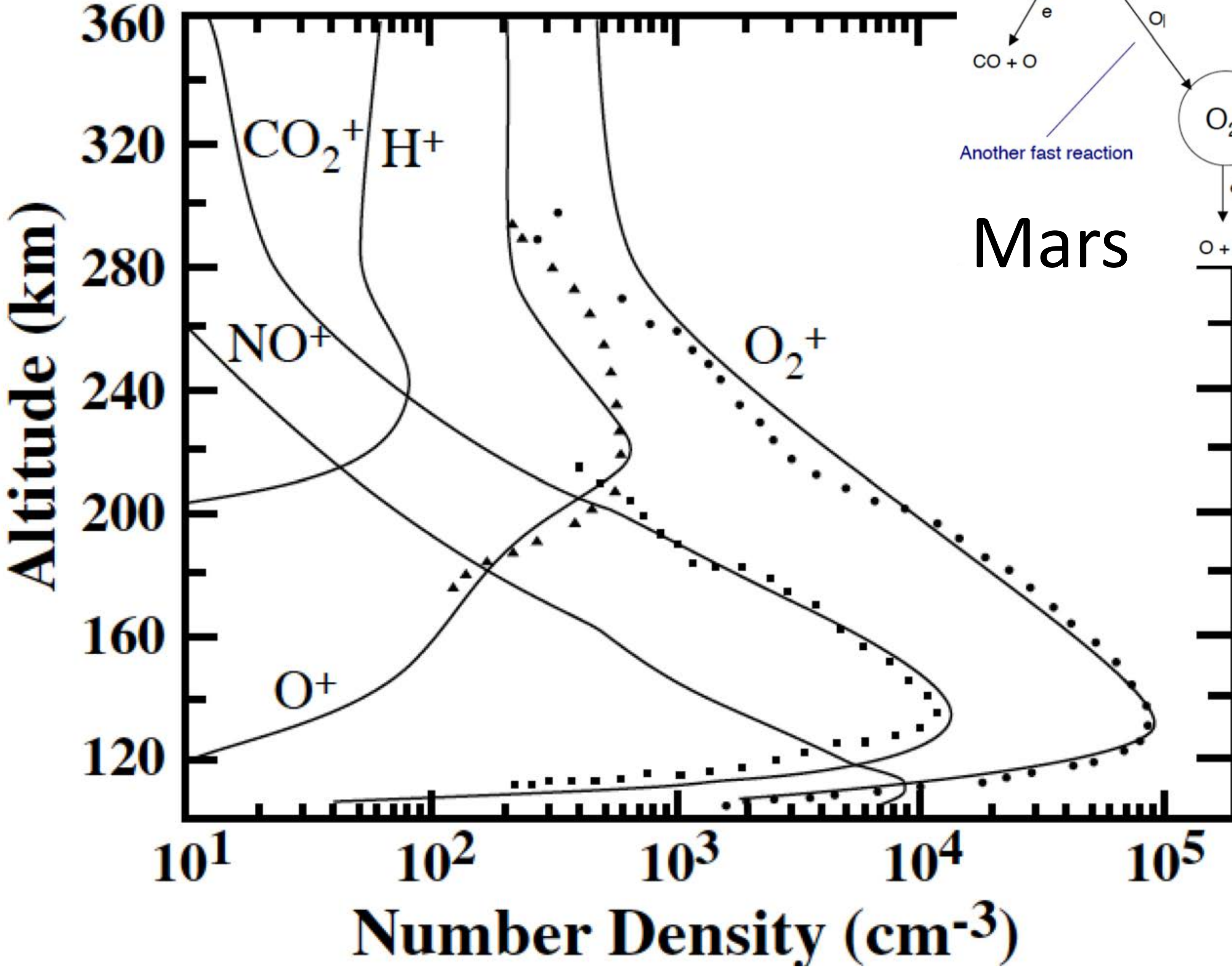
Principal Ionization Processes on Venus & Mars



vol III fig. 13.4

Mars

vol III fig. 13.6

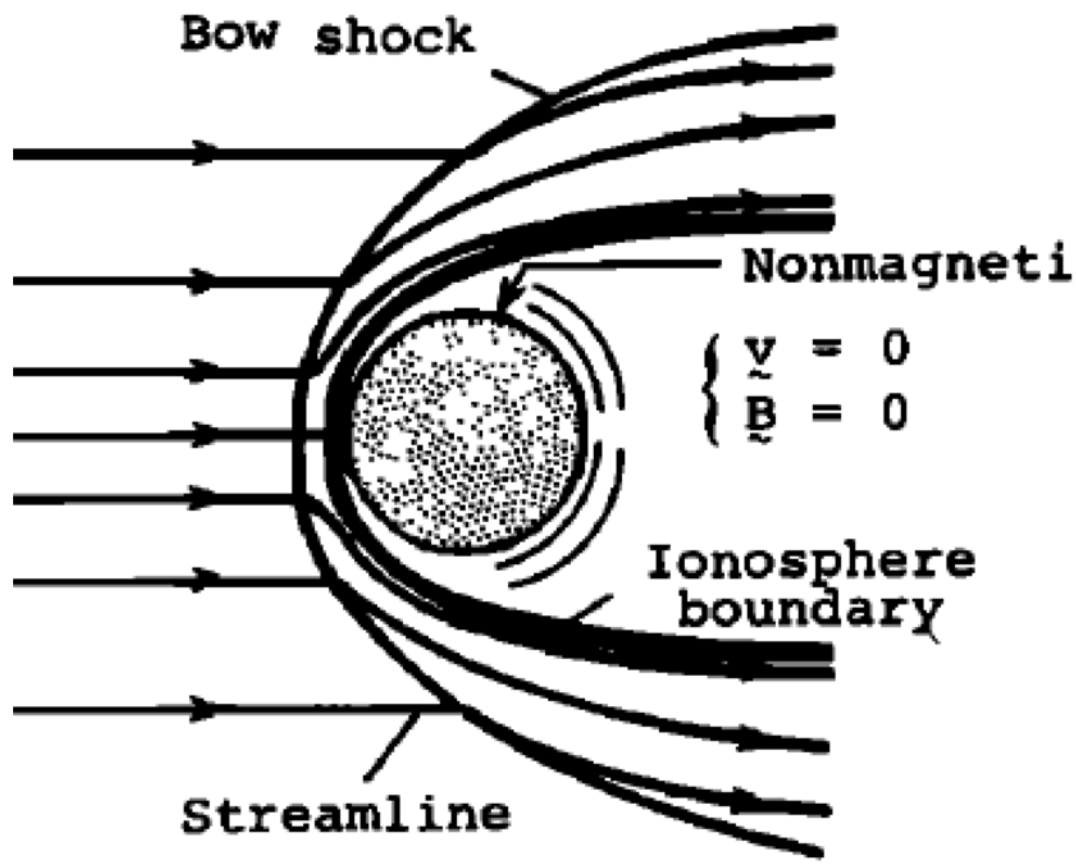


	Mars
N_2	3%
CO_2	95%
O_2	-
H_2O	-

Venus or Mars

- No dynamo – no **B**
- Ionosphere ➔ conducting bdry
- SW– w/ **B** – can't penetrate
- Supersonic flow deflected by obstacle
- Bow shock forms

Spreiter & Stahara 1980

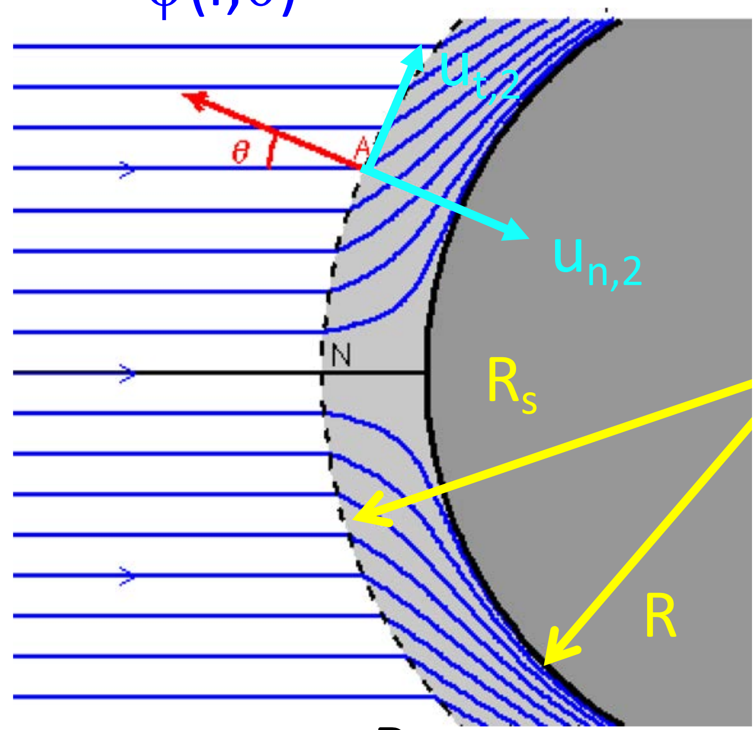


Simple picture of bow shock

- Ignore pressure from SW **B**
- SW: $u_\infty/c_{s,\infty}$, ρ_∞, $M_\infty \gg 1$
- Standing shock ~ sphere radius= R_s
- Post-shock flow
  - v. ubsonic – $M \ll 1$
  - ➔ ~ incompressible w/ $u_r(R) = 0$

$$\mathbf{u} = \nabla\psi \times \nabla\phi$$

$$\psi(r,\theta) = C\left(\frac{r^4}{R^4} - \frac{R^2}{r^2}\right)\sin^2\theta$$ Lighthill 1957



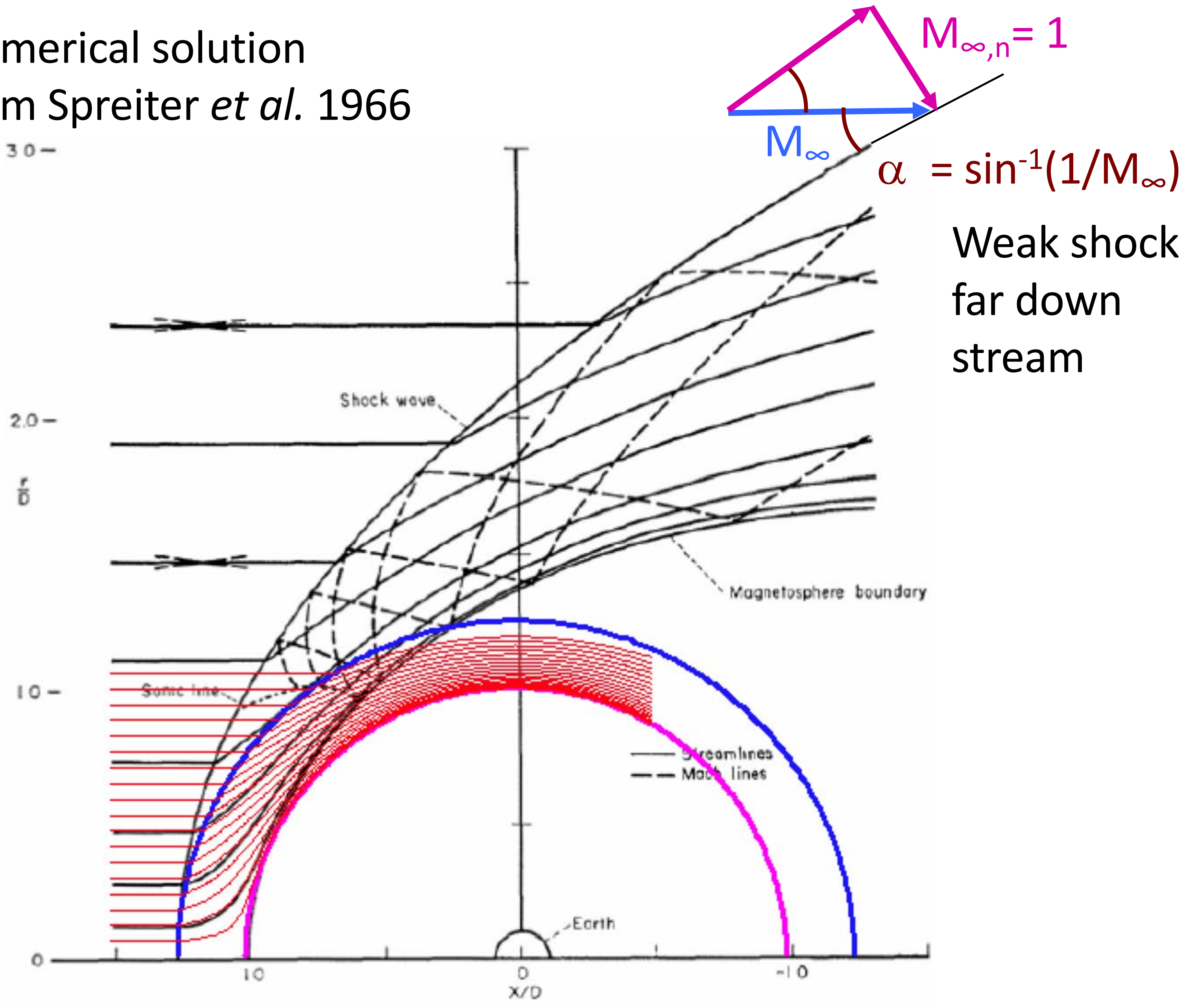
  - $u_{n,2} = u_{n,1}/4$, $u_{t,2} = u_{t,1}$

$$\frac{u_{r,2}}{\cos\theta} = 2\frac{CR_s^2}{R^4}\left(1 - \frac{R^5}{R_s^5}\right) = -\frac{1}{4}u_\infty \qquad \frac{u_{\theta,2}}{\sin\theta} = -\frac{CR_s^2}{R^4}\left(4 + \frac{R^5}{R_s^5}\right) = u_\infty$$

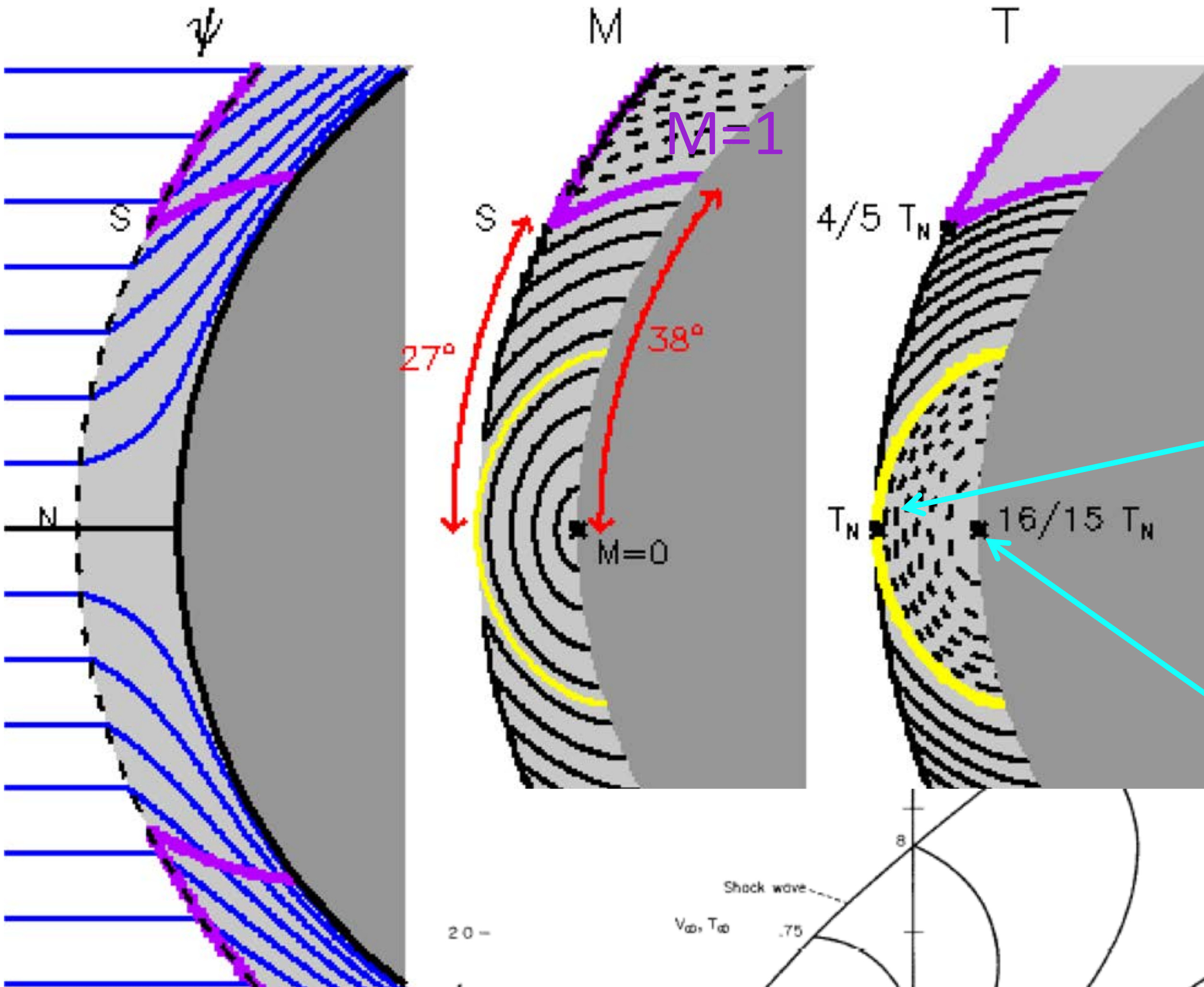
Bow shock

$$R_s = \left(\frac{3}{2}\right)^{2/5} R = 1.18\, R$$

Numerical solution
from Spreiter *et al.* 1966



Weak shock
far down
stream



Shock partially thermalizes flow KE of SW:

- Nose point (normal)

$$T_N = \frac{3}{8} \cdot \frac{\frac{1}{2} m u_\infty^2}{k_B}$$

- Stagnation point

$$T_s = \frac{16}{15} T_N = \frac{2}{5} \cdot \frac{\frac{1}{2} m u_\infty^2}{k_B}$$

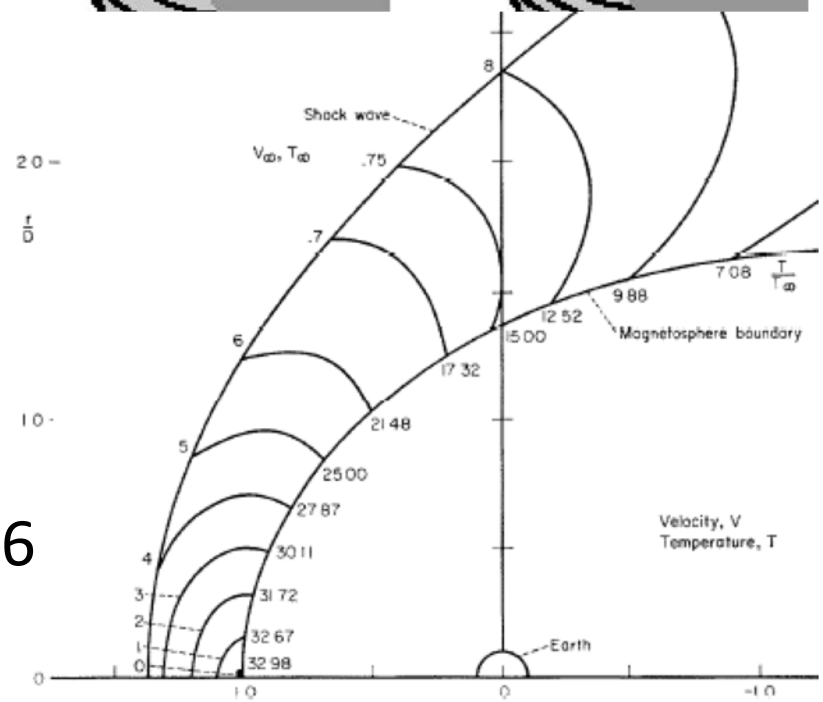
u_∞ = 400 km/s

➔ T_N = 3.6 MK

➔ T_s = 3.8 MK

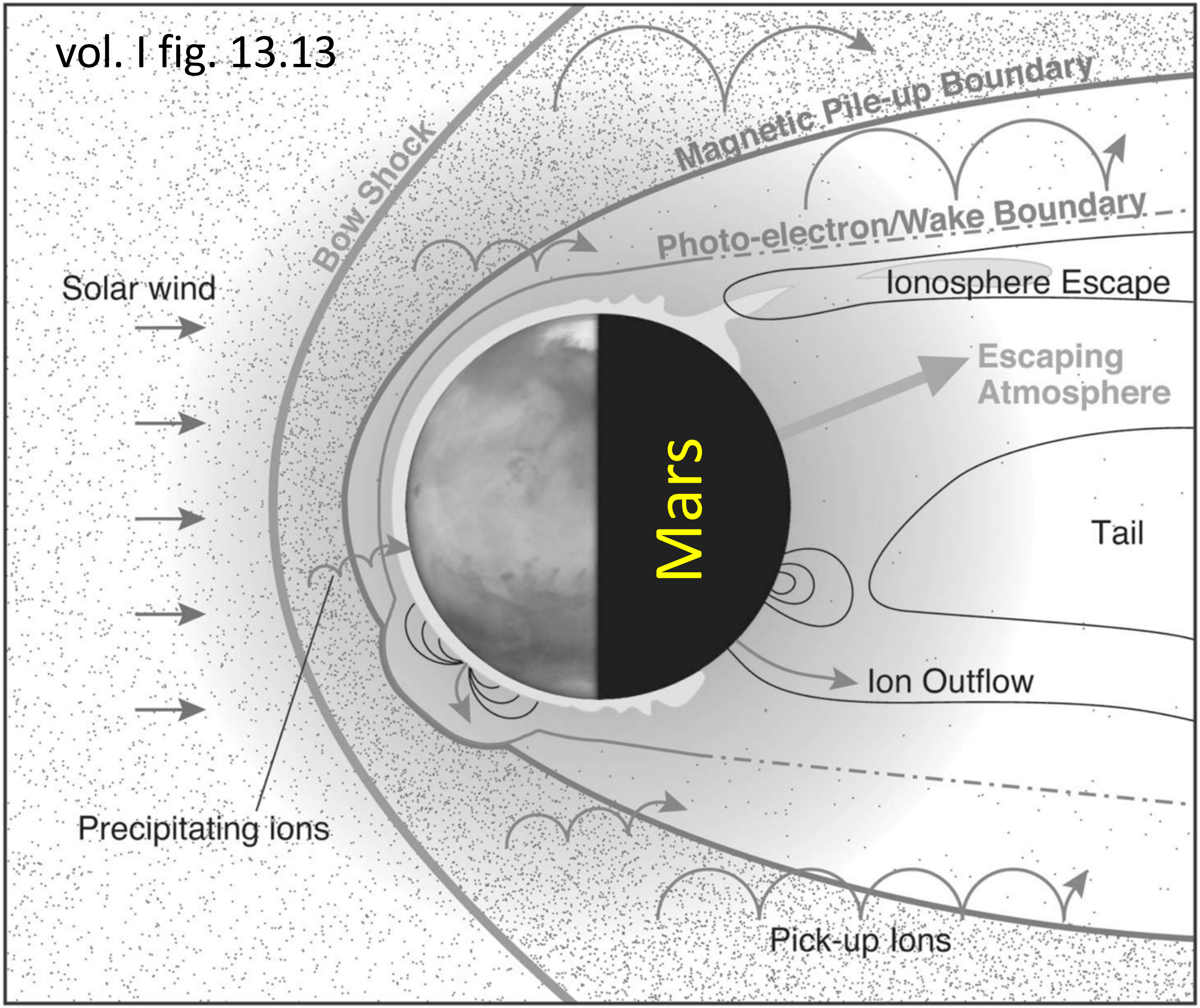
- pressure

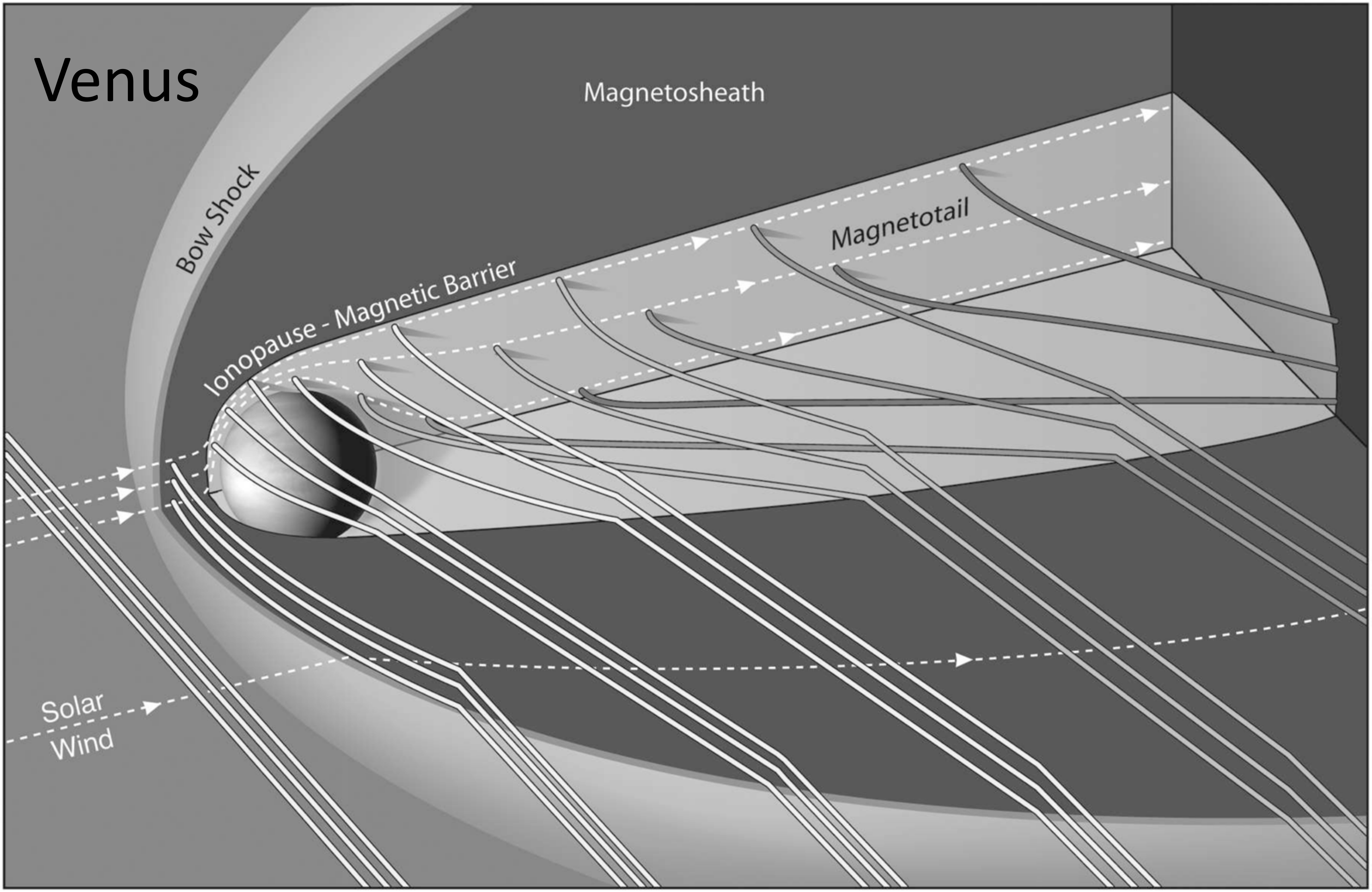
$$p_s = \frac{4}{5} \rho_\infty u_\infty^2$$



Spreiter *et al.* 1966

vol. I fig. 13.13

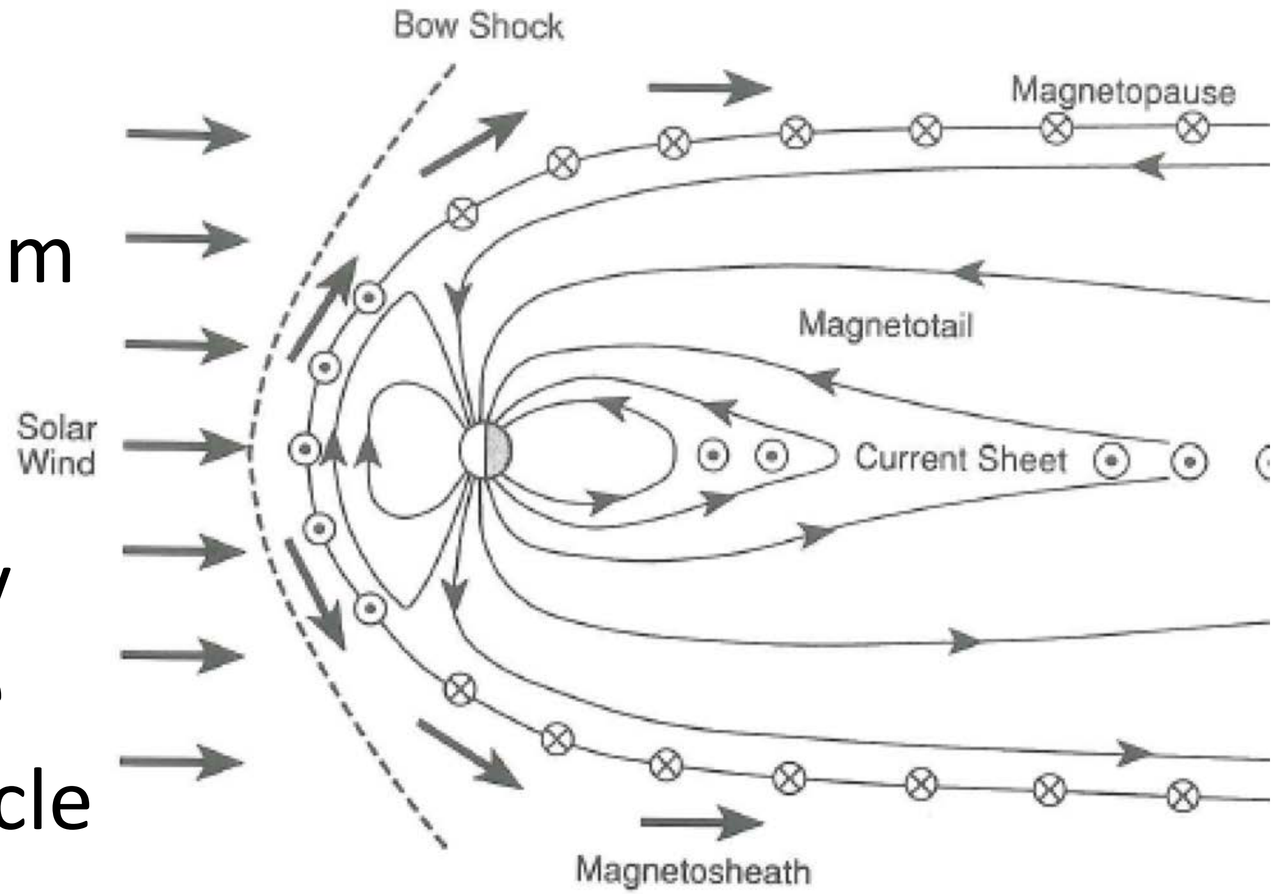




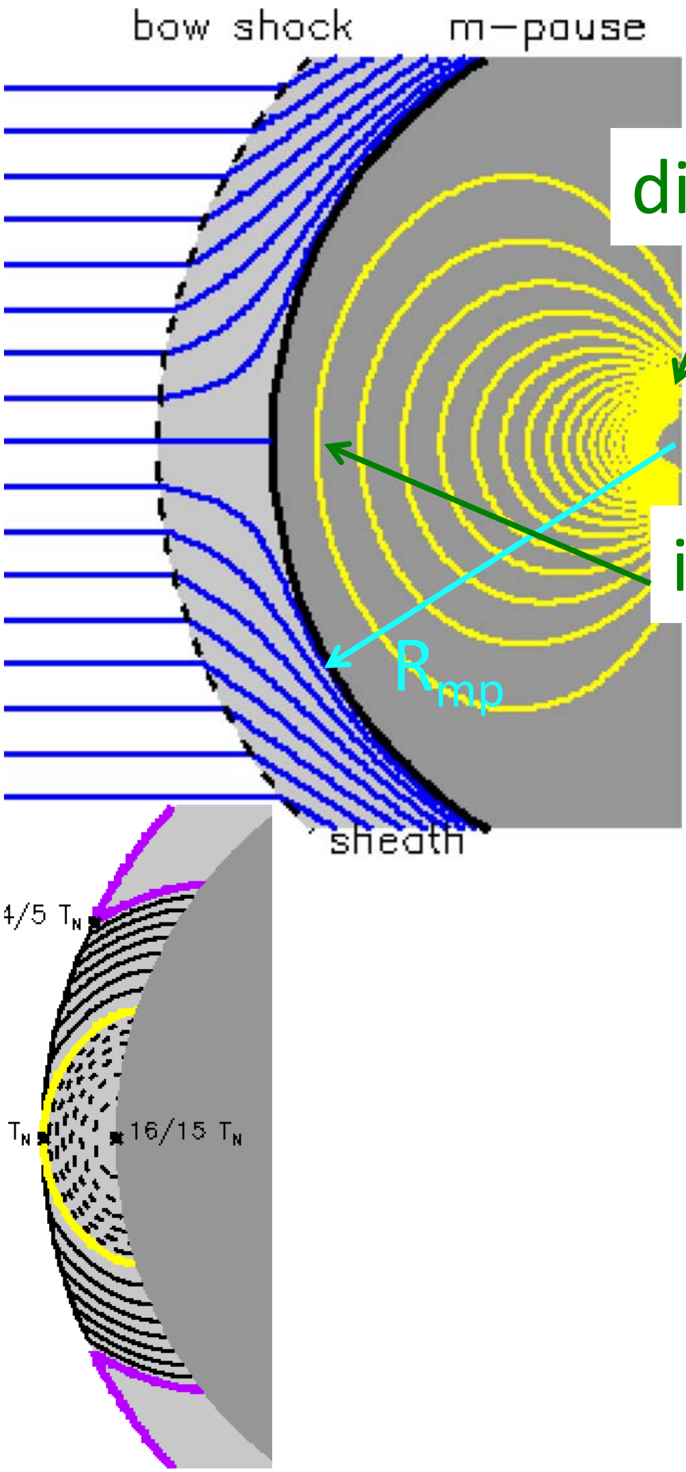
vol. I fig. 13.12

Wind @ Magnetized Planets Earth, Jupiter, Saturn, …

- Planetary **B** prevents SW from reaching ionosphere
- SW deflected by **magnetosphere**
- “squishy” obstacle



Hughes (*cf.* vol. I fig. 10.1)



Shock & sheath: similar to before

- Stagnation point (SP) @ r=R_{mp}
- Plasma pressure: $p_s = \frac{4}{5}\rho_\infty u_\infty^2$
- Inside (r < R_{mp}): $\mathbf{B} = -\nabla\chi$

$$\chi(r,\theta) = \frac{B_\oplus R_\oplus^3}{R_{mp}^2}\left(\frac{R_{mp}^2}{r^2} + \frac{2r}{R_{mp}}\right)\cos\theta$$

- Magnetic pressure @ SP

$$\frac{1}{8\pi}\left|\mathbf{B}(R_{mp},0)\right|^2 = \frac{1}{8\pi}\left(\frac{1}{R_{mp}}\frac{\partial\chi}{\partial\theta}\right)^2 = \frac{9R_\oplus^6}{8\pi R_{mp}^6}B_\oplus^2$$

- Ignore inner plasma – balance

$$R_{mp} = \left(\frac{45}{32\pi}\right)^{1/6}\left(\frac{B_\oplus^2}{\rho_\infty u_\infty^2}\right)^{1/6} R_\oplus$$

Chapman-Ferraro Distance

Intuition break

$$R_{mp} = \left(\frac{45}{32\pi}\right)^{1/6}\left(\frac{B_\oplus^2}{\rho_\infty u_\infty^2}\right)^{1/6} R_\oplus \sim 12\ R_\oplus$$

$\rho_{sw} = 10^{-23}$ g/cm

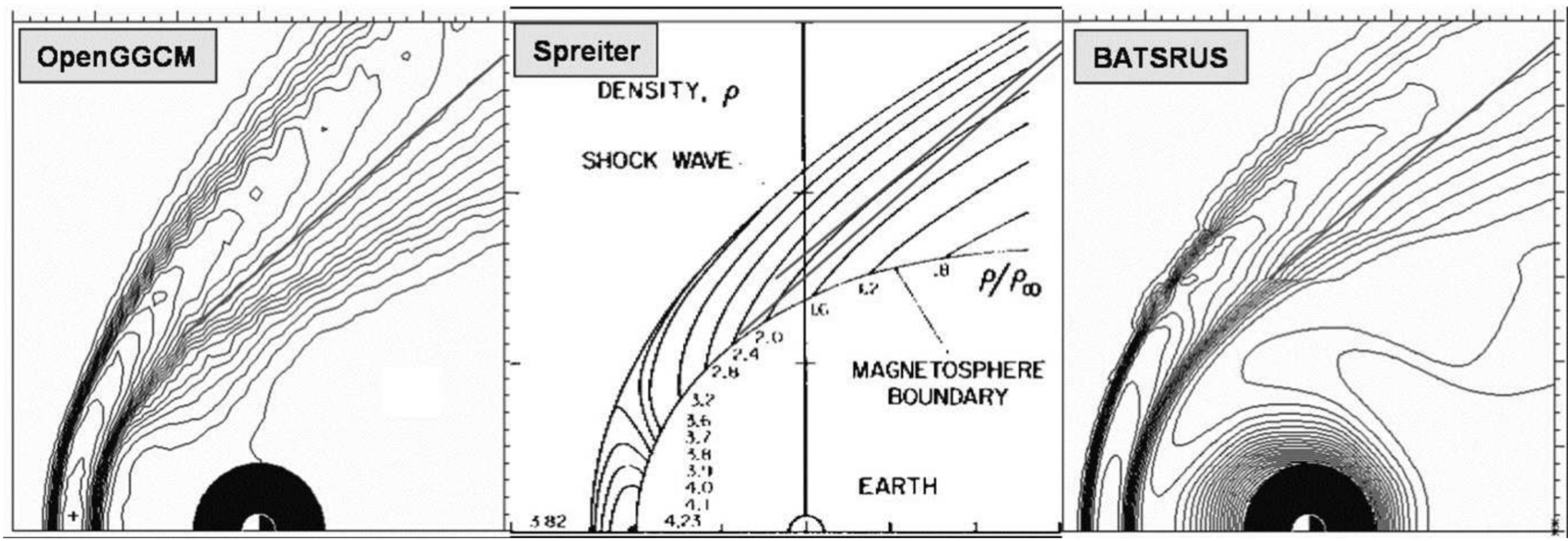
$u_{sw} = 400$ km/s

- At what distance do geostationary satellites orbit?

- Is the moon inside or outside the magnetopause?

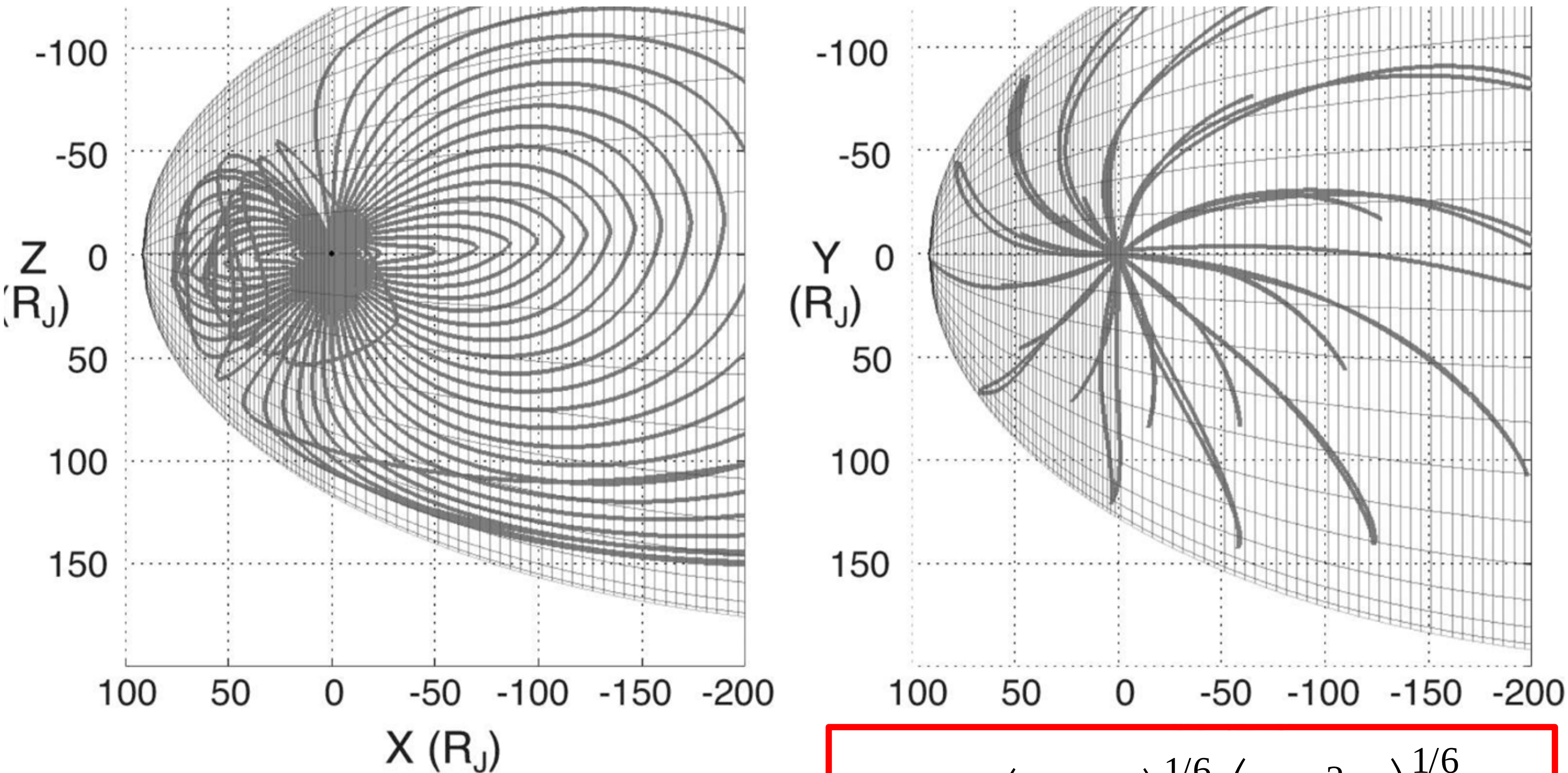
- What happens to R_{mp} during fast SW: $u_{sw} = 800$ km/s

Similar picture from high-powered codes



vol. I fig. 11.2

Other planets… same story



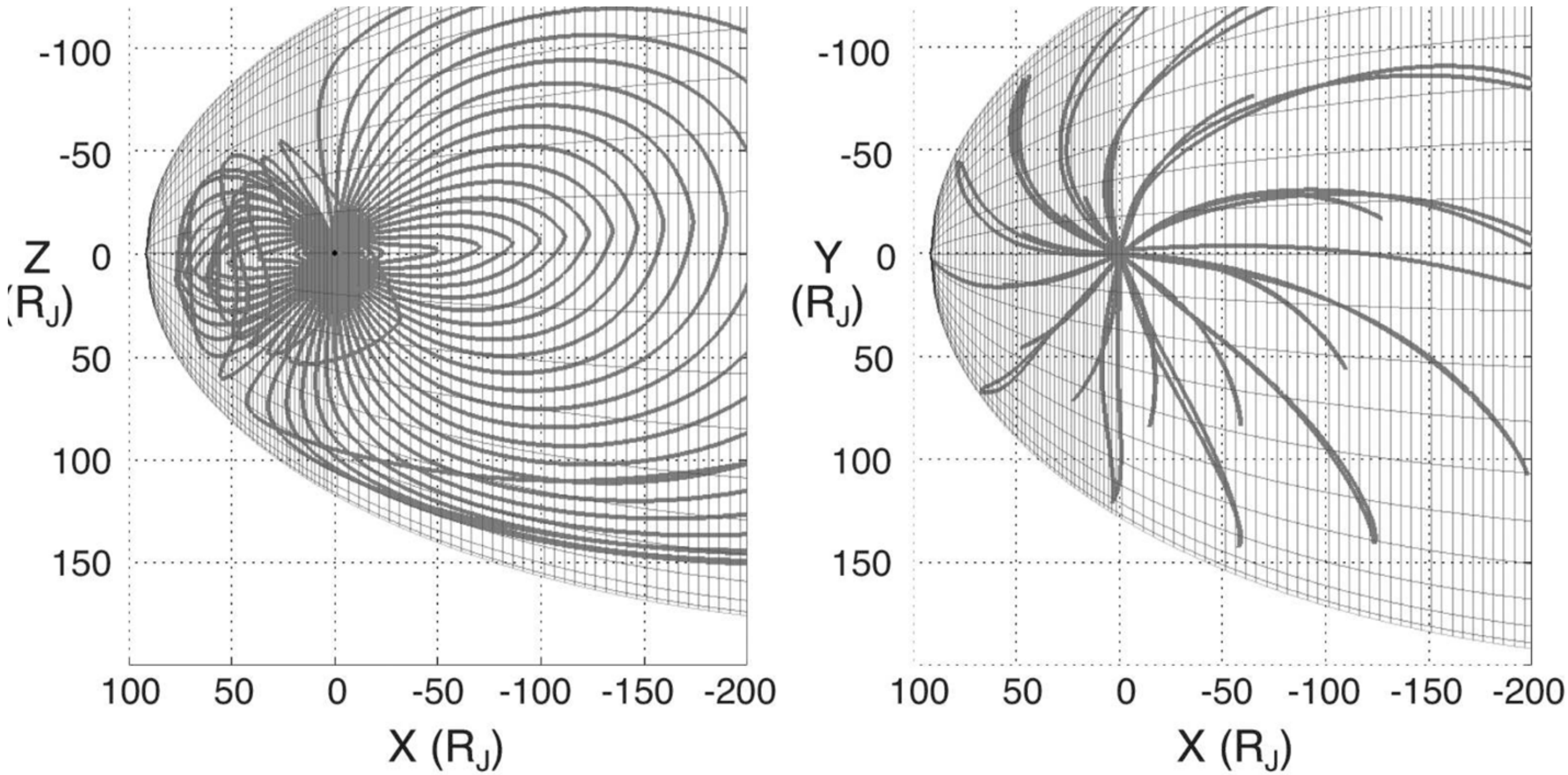
vol. I fig. 13.6

$B_J \sim 15\, B_{\oplus} \sim 5$ G

$$R_{mp} = \left(\frac{45}{32\pi}\right)^{1/6} \left(\frac{B_{\oplus}^2}{\rho_\infty u_\infty^2}\right)^{1/6} R_{\oplus}$$

Q: how do u_∞ & ρ_∞ @ Jupiter compare to @ Earth?

Other planets… same story



vol. I fig. 13.6

$B_J \sim 15\ B_\oplus \sim 5$ G ; $\rho_\infty \sim 0.04\ \rho_{\infty,\oplus}$

➔ Jupiter's magnetopause:

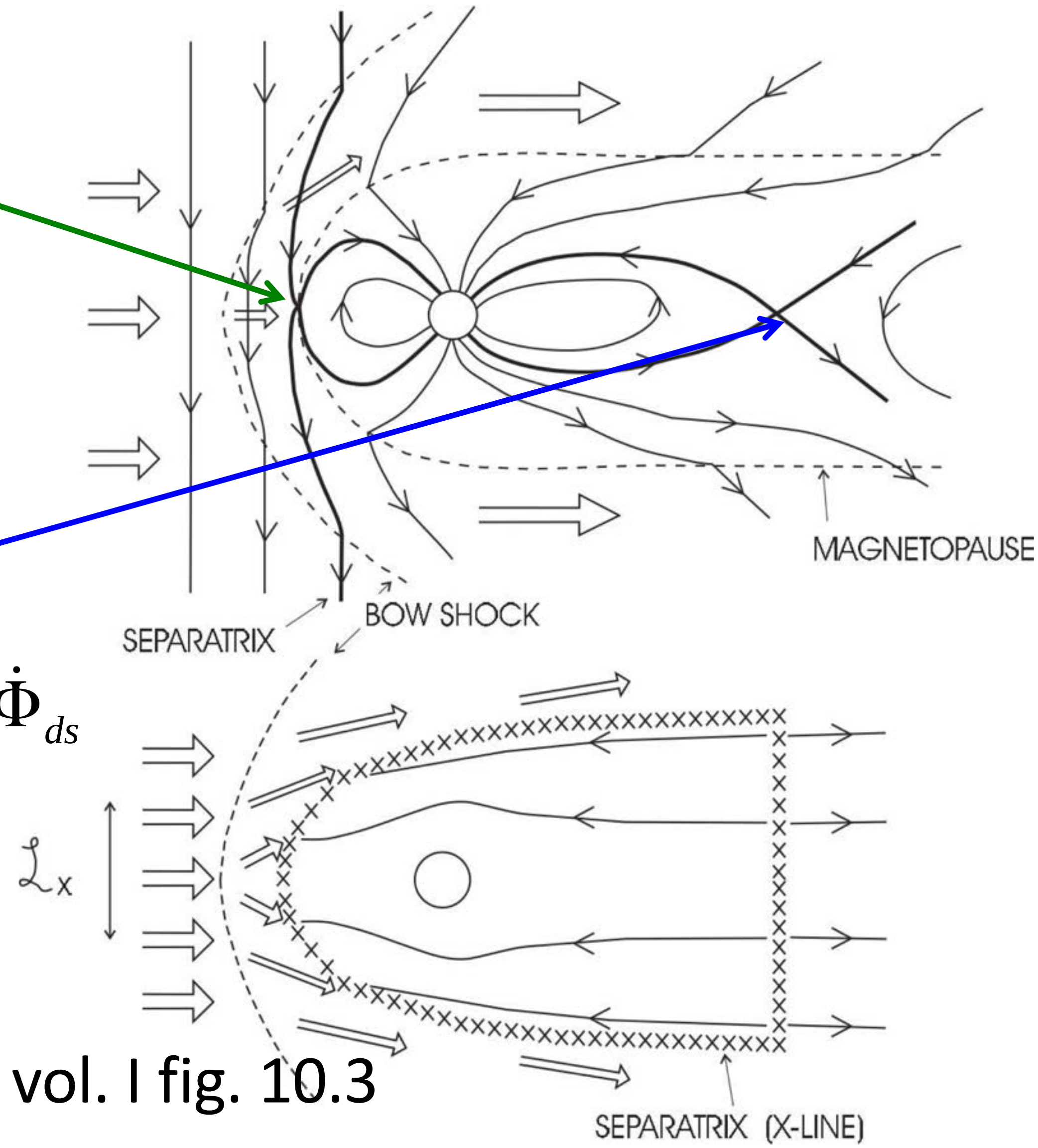
$R_{mp,J} \sim 50\ R_J = 3.5 \times 10^{11}$ cm

$$R_{mp} = \left(\frac{45}{32\pi}\right)^{1/6} \left(\frac{B_J^2}{\rho_\infty u_\infty^2}\right)^{1/6} R_J$$

But not all of Earth's field stays confined to m-sphere

Reconnection with SW field (consider southward IMF)

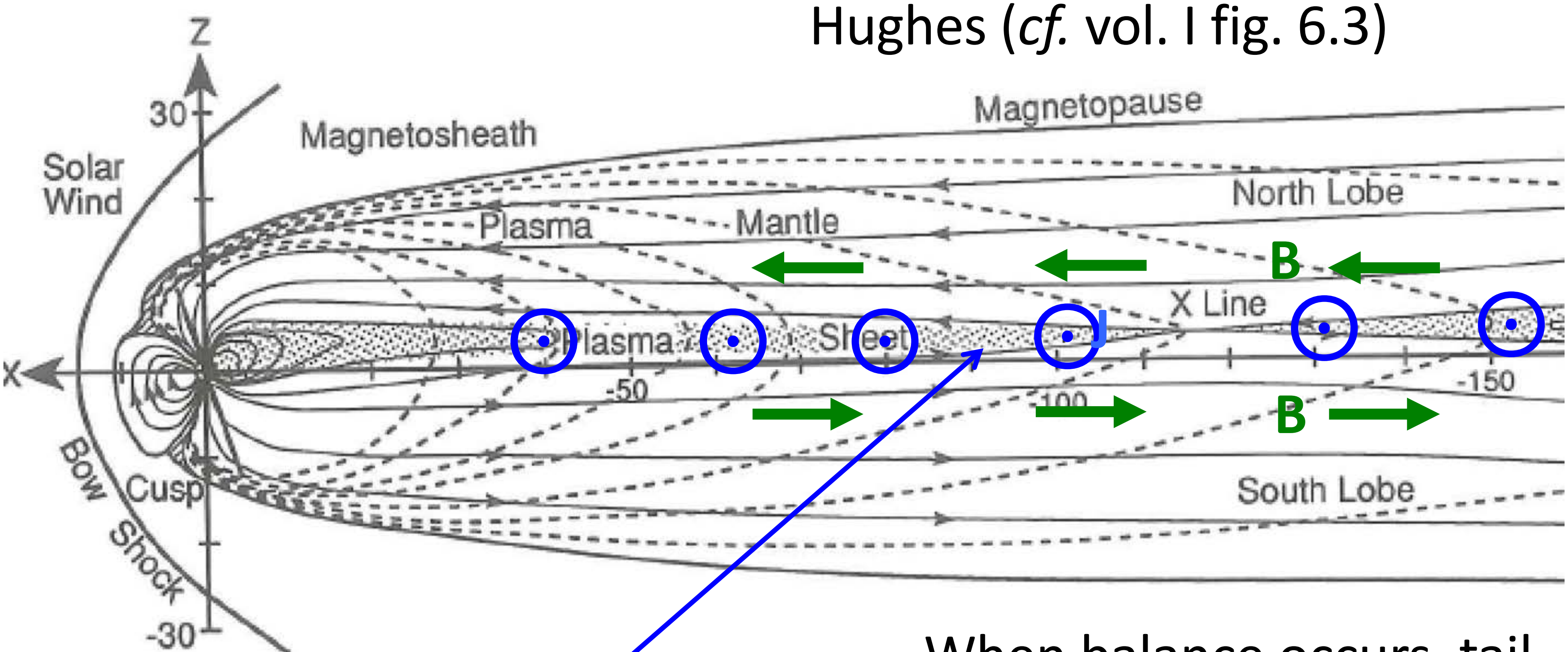
- Creates "open" flux connected to poles @ $\dot{\Phi}_{ds}$
- SW sweeps flux downstream – into **magnetotail**
- Steady state only when reconnection in tail "closes" flux at rate $\dot{\Phi}_{n} = -\dot{\Phi}_{ds}$
- Requires long & strong **neutral sheet** in magnetotail



vol. I fig. 10.3

But not all of Earth's field stays confined to m-sphere

Hughes (*cf.* vol. I fig. 6.3)



"closes" flux at rate $\dot{\Phi}_n = -\dot{\Phi}_{ds}$

- Requires long & strong **neutral sheet** in magnetotail

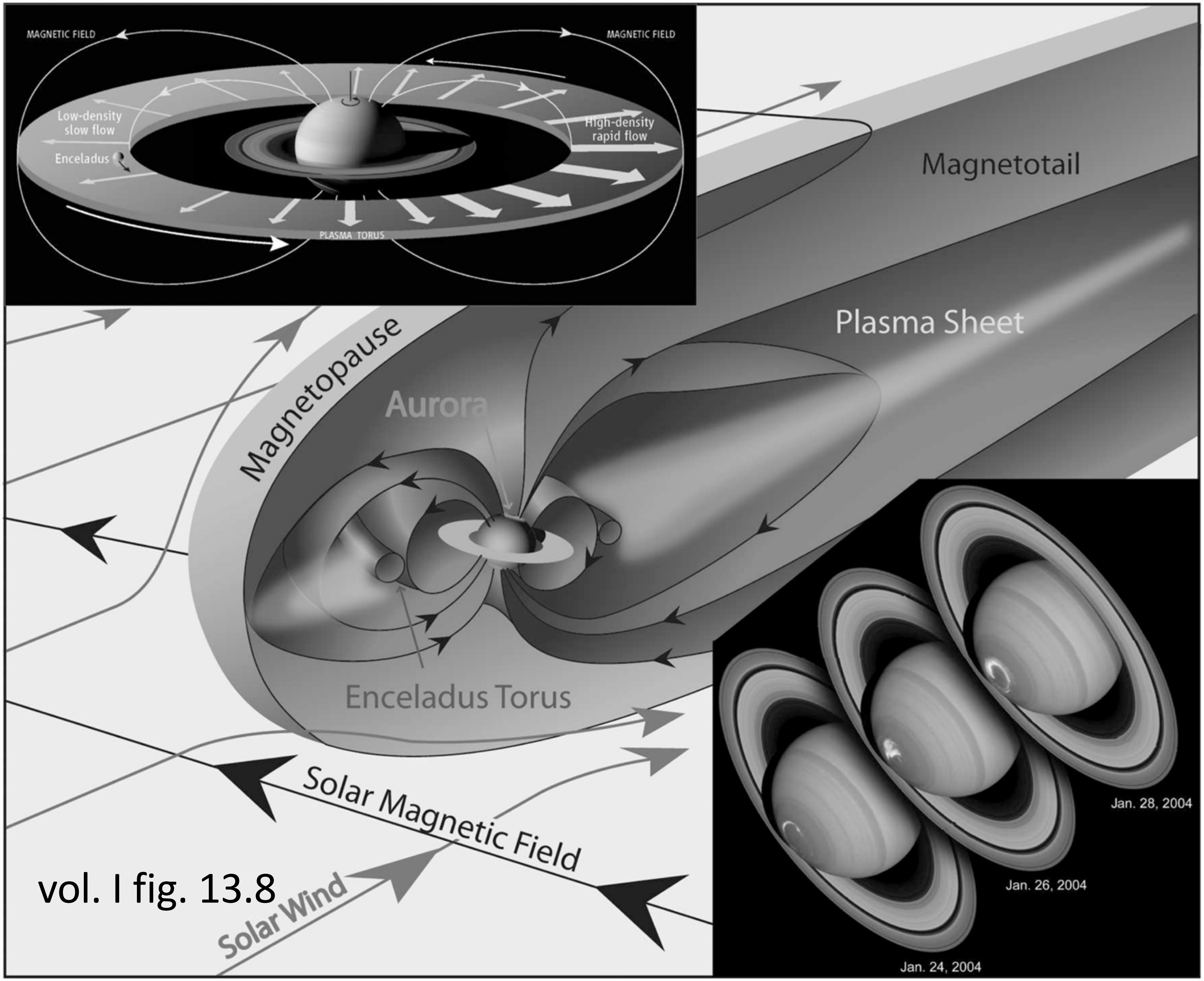
When balance occurs, tail…

- … has some length

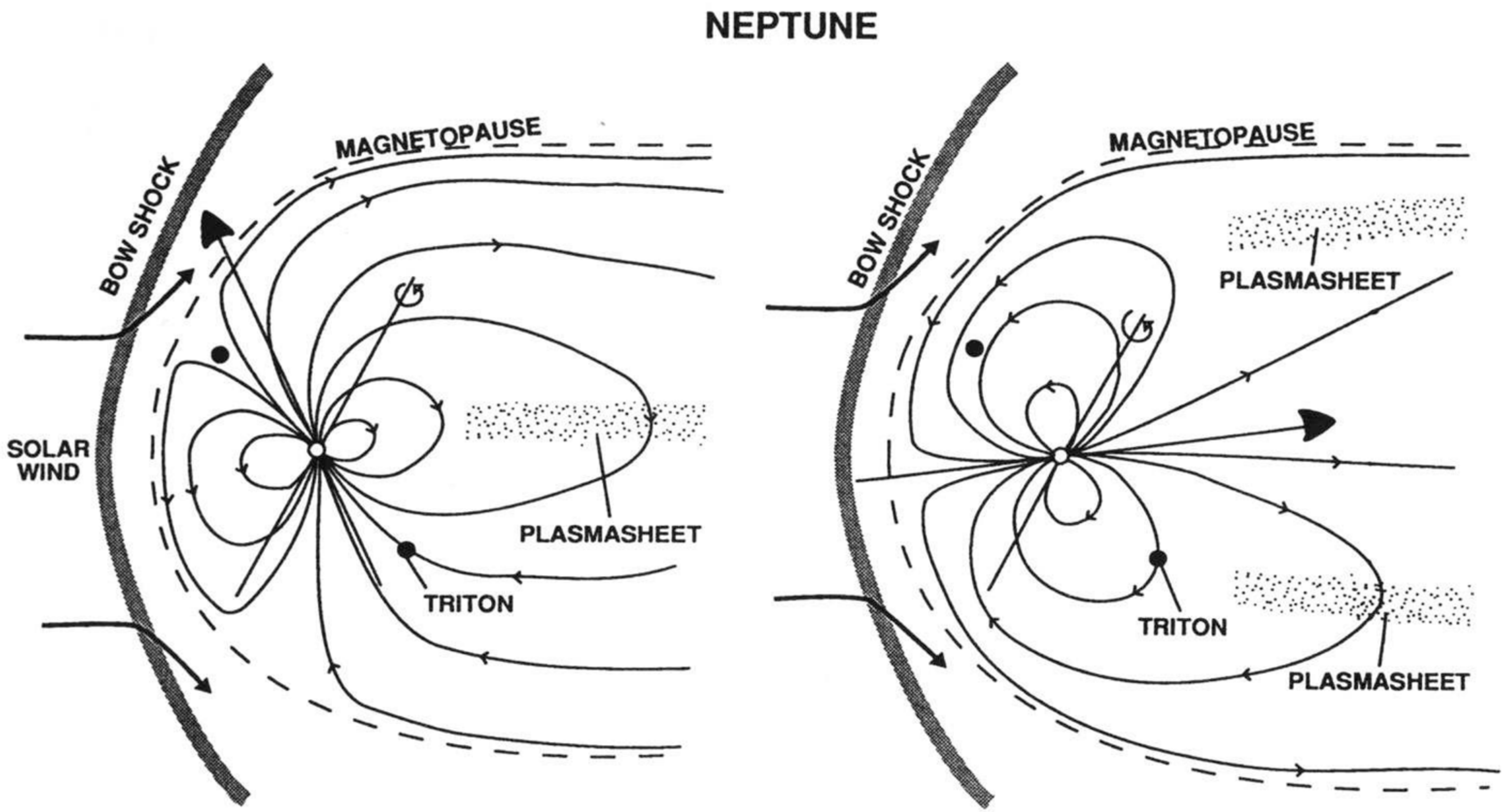
  $L_t \gg R_{mp}$

- … has some open flux

  Φ_t

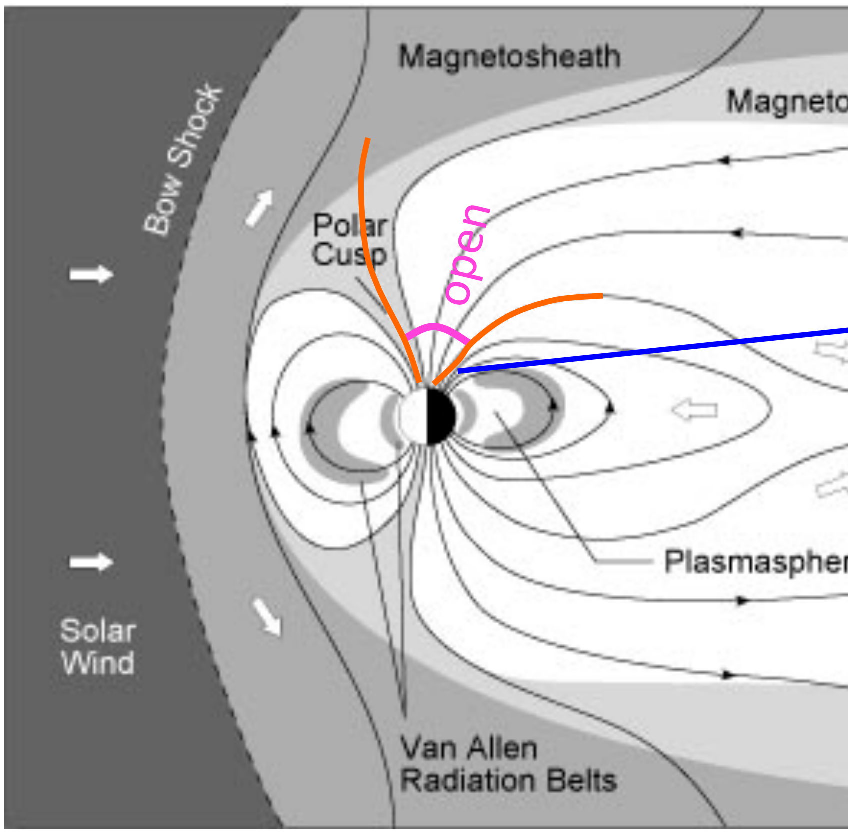

MAGNETIC FIELD
MAGNETIC FIELD
Low-density
slow flow
Enceladus
High-density
rapid flow
PLASMA TORUS
Magnetotail
Plasma Sheet
Magnetopause
Aurora
Enceladus Torus
Solar Magnetic Field
Solar Wind
Jan. 28, 2004
Jan. 26, 2004
Jan. 24, 2004

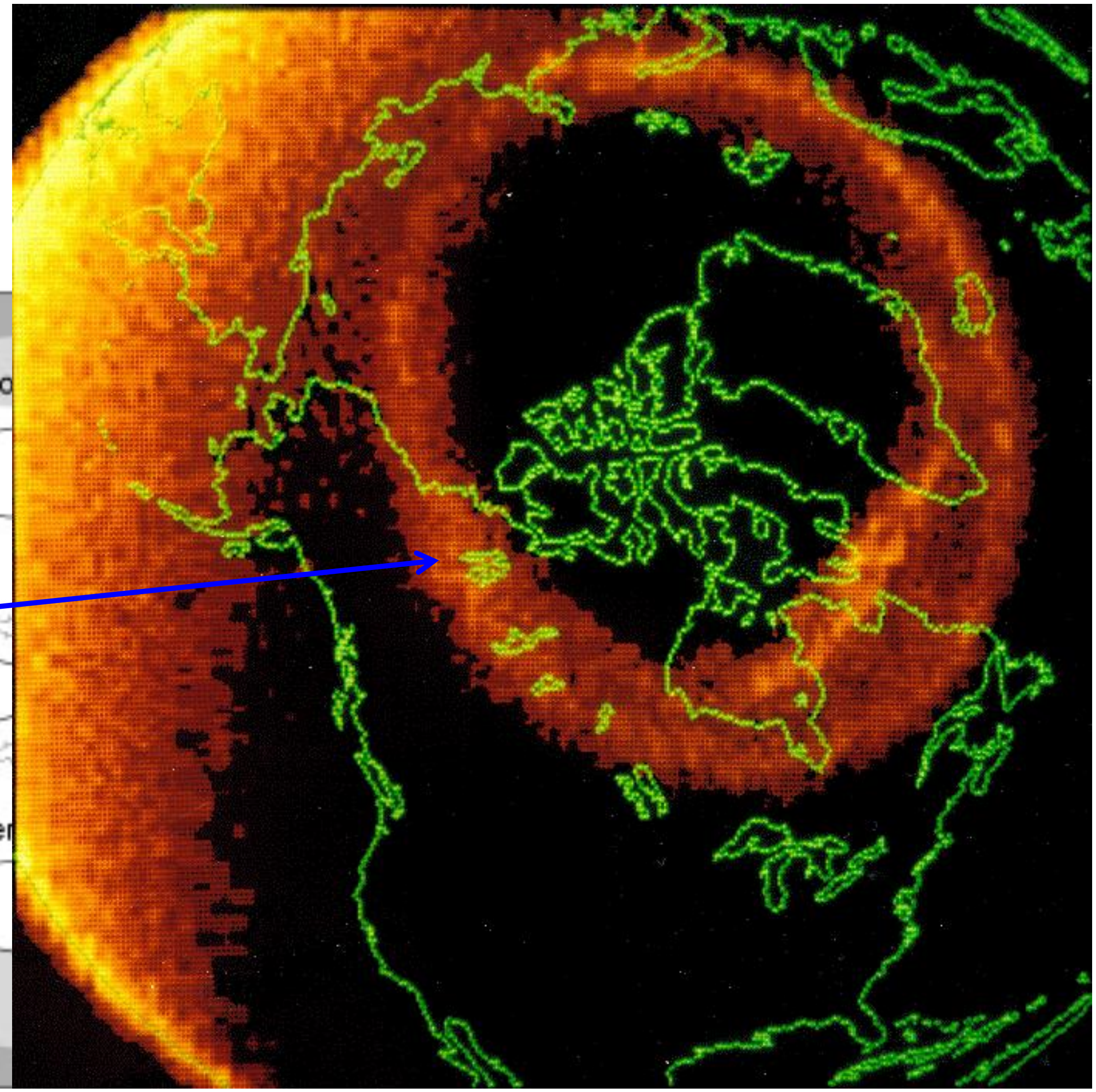
vol. I fig. 13.8

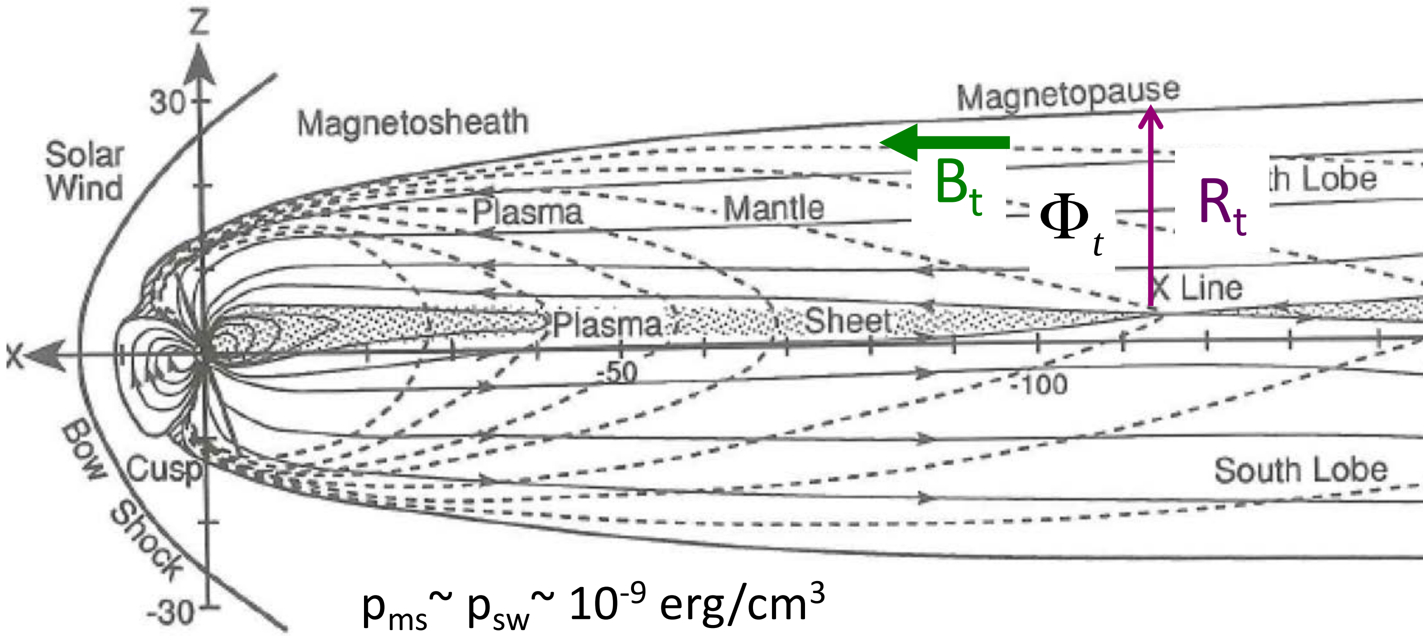


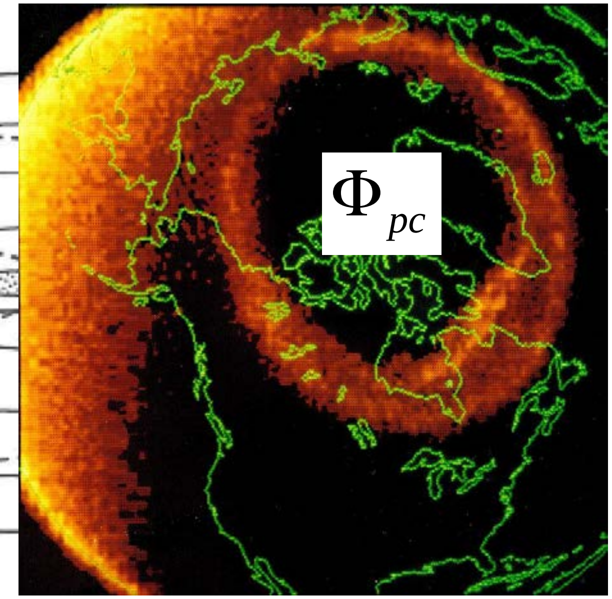
vol. I fig. 13.10

closed/open boundary maps down to "auroral oval"









$p_{ms} \sim p_{sw} \sim 10^{-9}$ erg/cm^3

$$\Phi_t = \Phi_{pc} = \pi \left(R_\oplus \sin\theta_{pc}\right)^2 B_{np} \sim \pi R_\oplus^2 \theta_{pc}^2 B_{np} \sim 10^{17}\,\mathrm{Mx}$$

$$\Phi_t = \frac{\pi}{2} R_t^2 B_t$$

mag. pressure In tail:

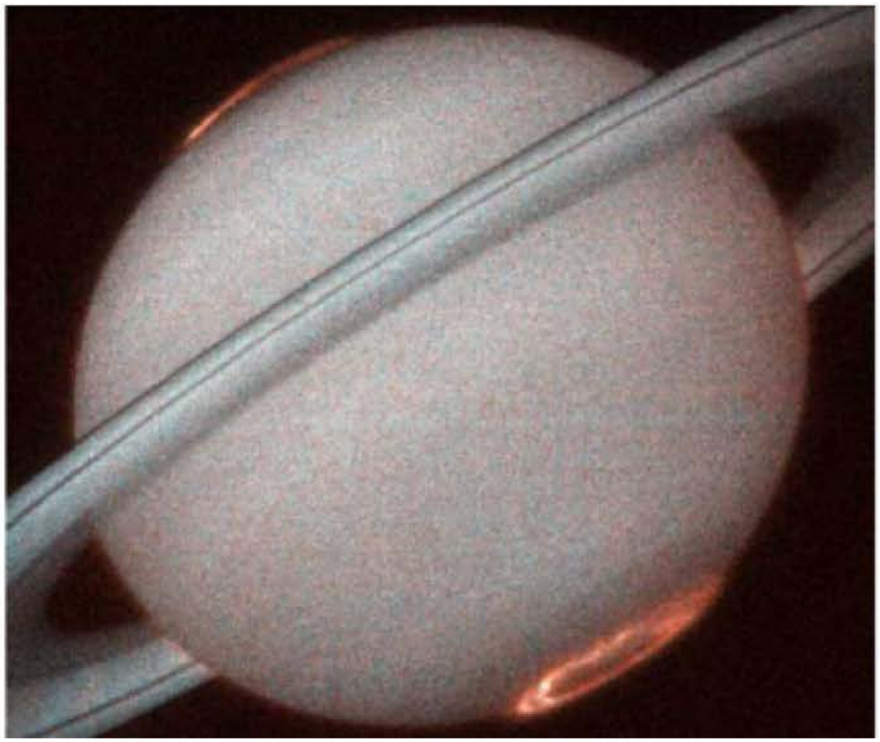
$$\frac{1}{8\pi} B_t^2 = \frac{1}{2\pi^3} \frac{\Phi_t^2}{R_t^4} = \frac{1}{2\pi} \left(\frac{R_\oplus}{R_t}\right)^4 \theta_{pc}^4 B_{np}^2$$

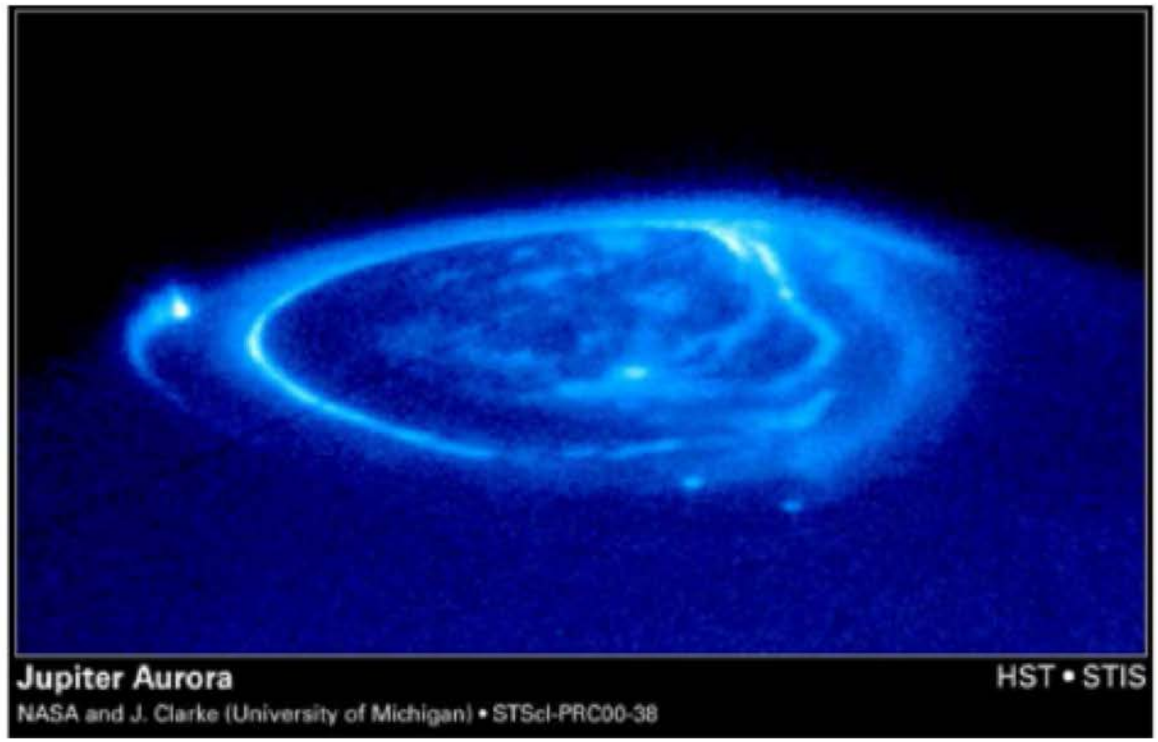
Pressure balance @ m-pause:

$$\frac{R_t}{R_\oplus} = (2\pi)^{-1/4} \frac{B_{np}^{1/2}}{p_{sw}^{1/4}} \theta_{pc} \sim 25$$

$B_t \sim 10^{-4}$ G ~ 10 nT

Other auroral ovals

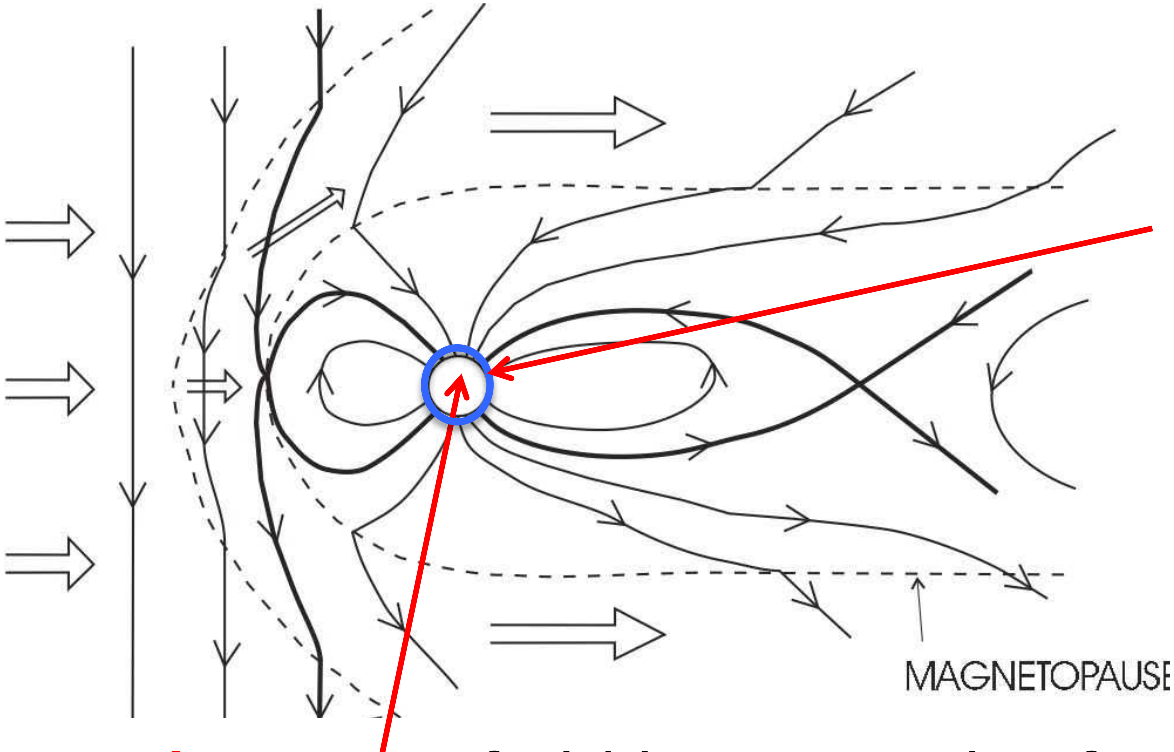




vol. I fig. 2.9

Convection: magnetosphere meets ionosphere

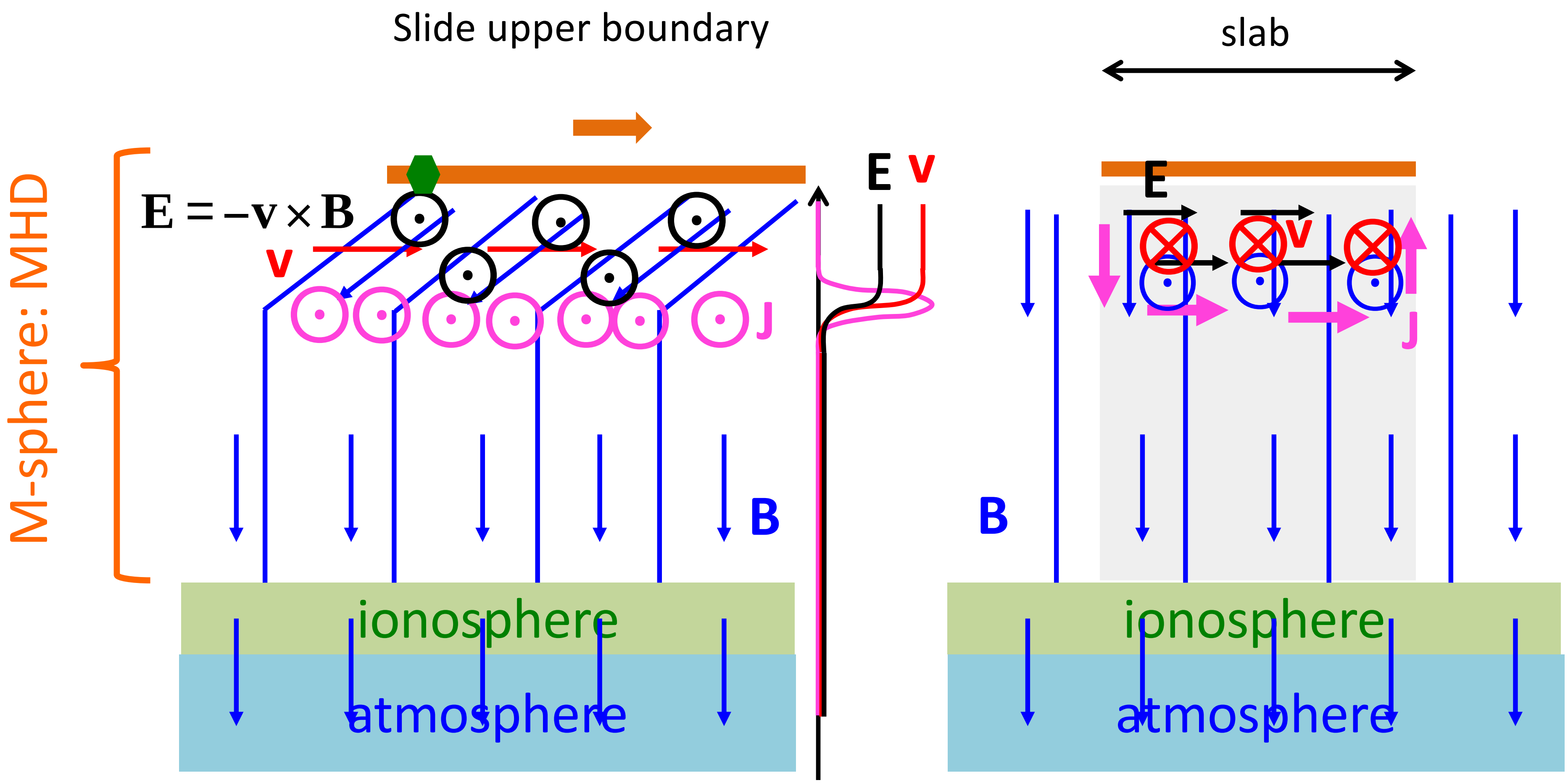
field lines are frozen to M-spheric plasma.
motion sweeps filed lines back



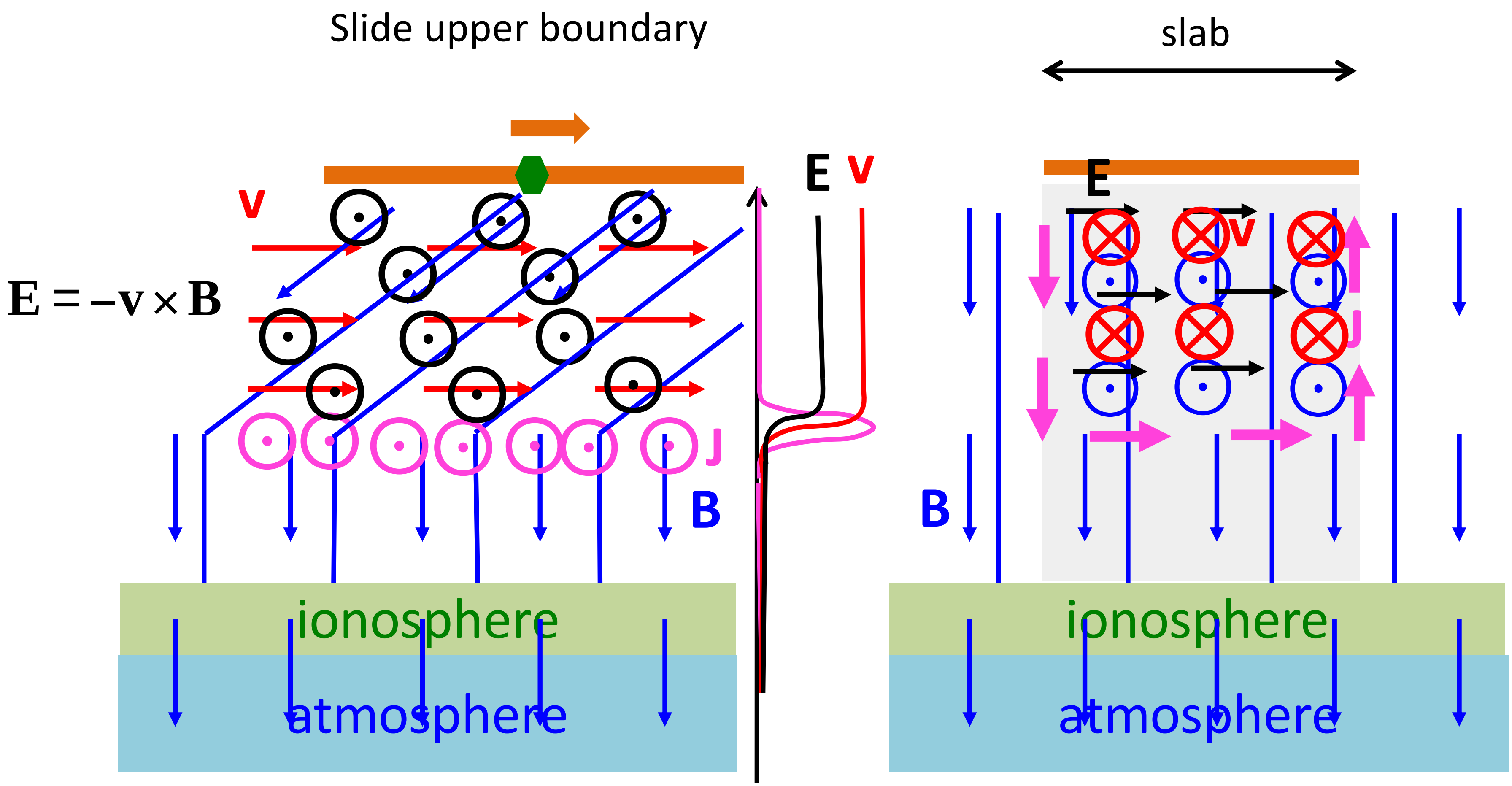
BUT: atmosphere & solid crust are insulators – field lines are **imaginary** there

Objection: field lines are also frozen into liquid core – ends cannot be moved

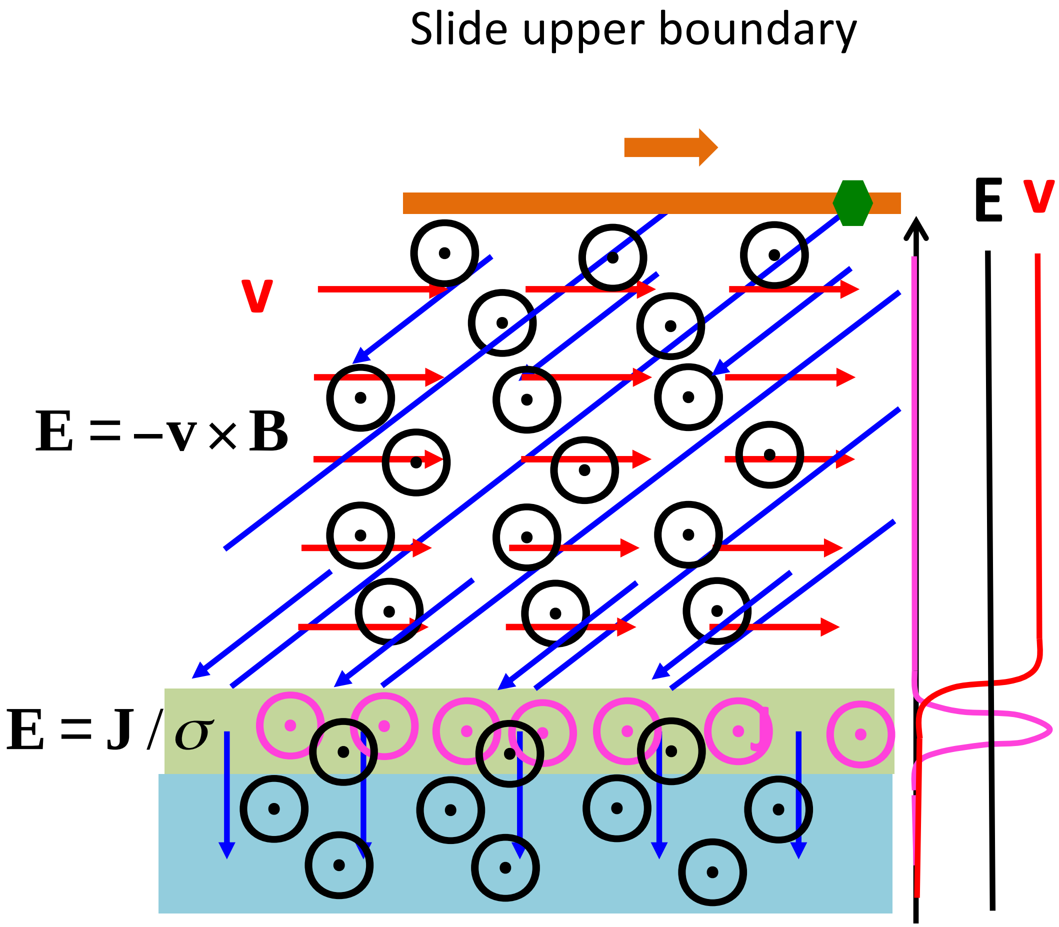
Example of how the motions meet

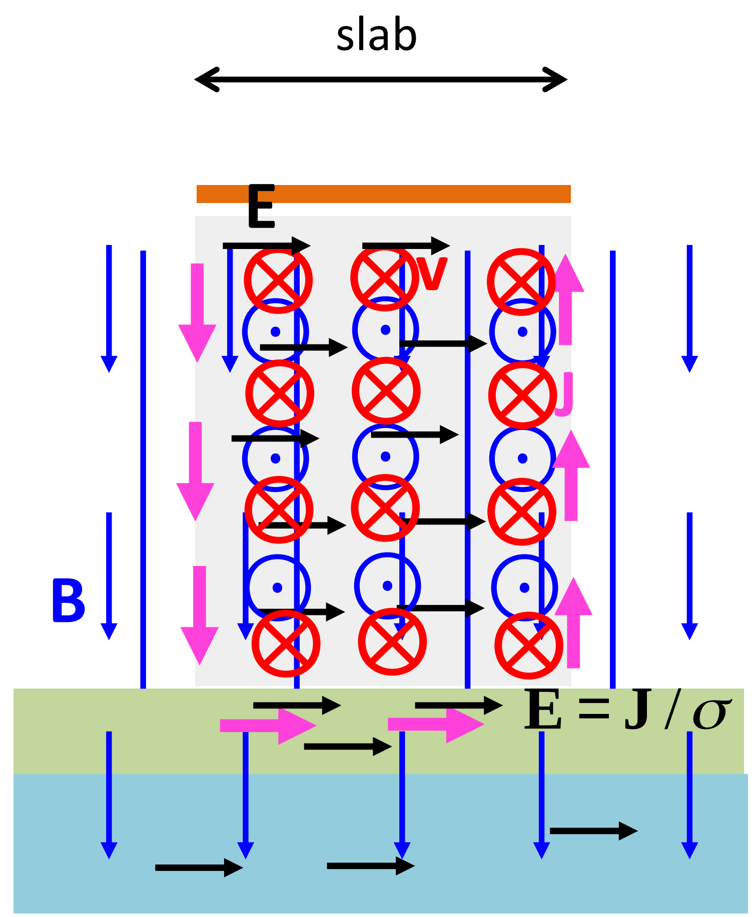


Example of how the motions meet

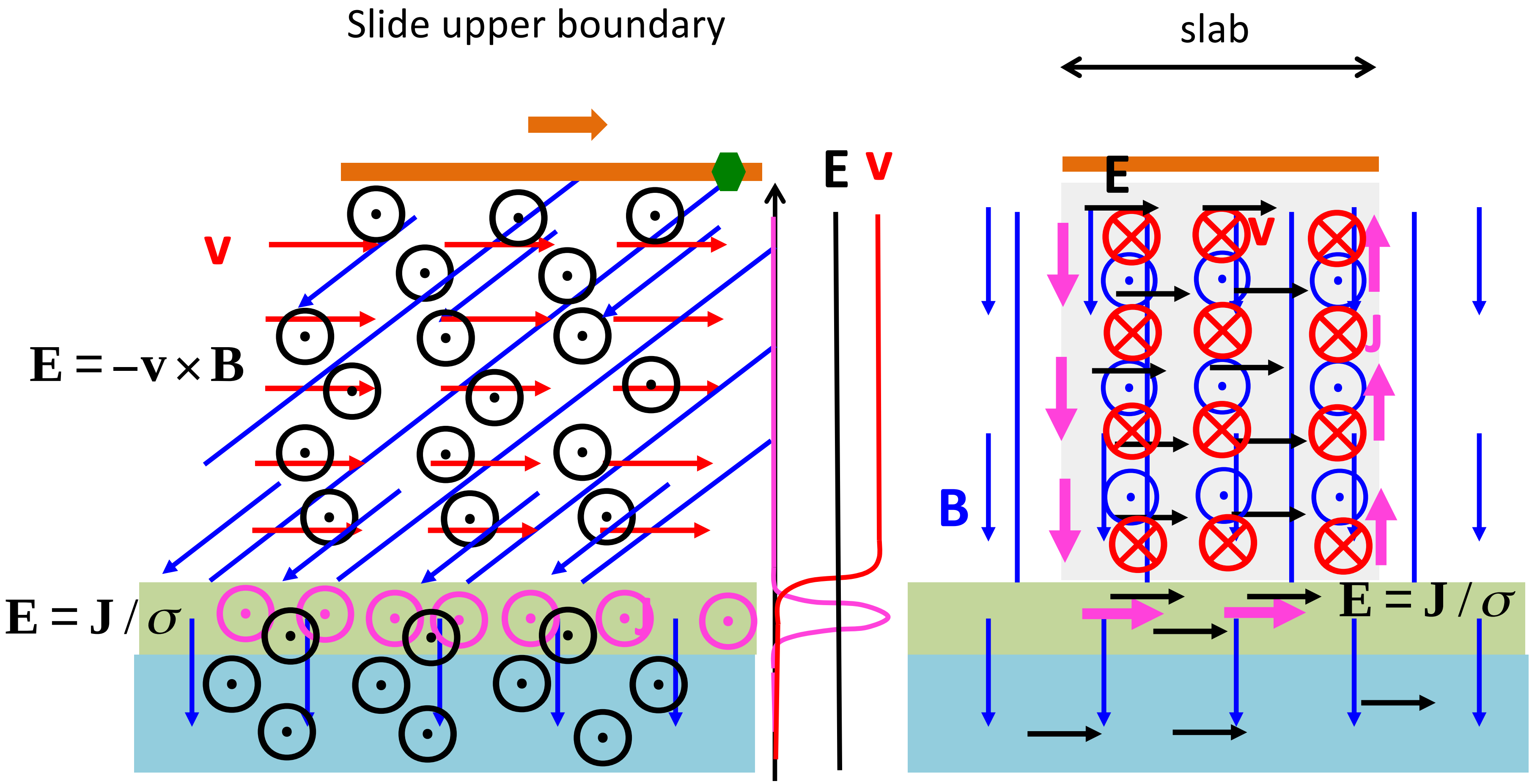


Example of how the motions meet

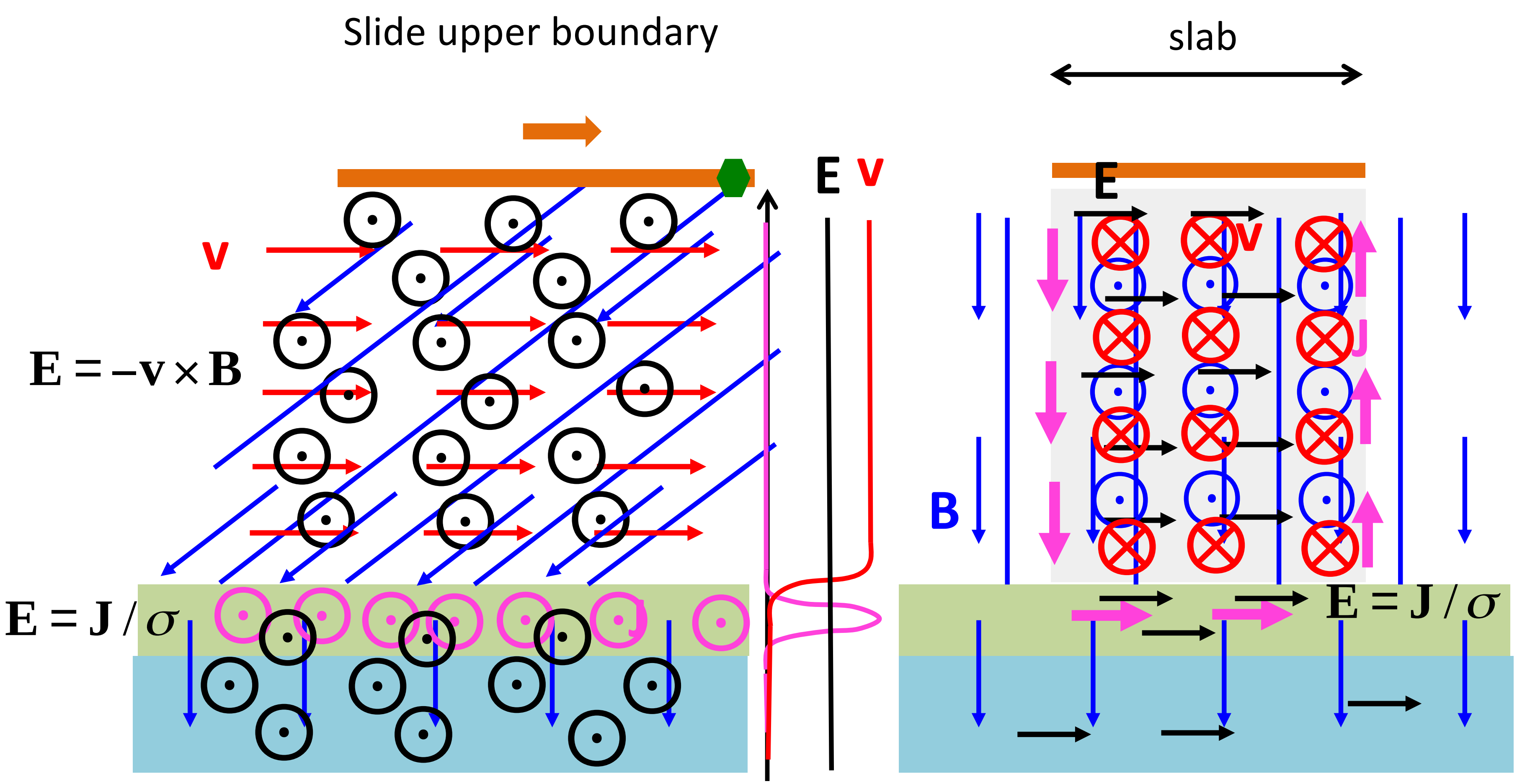




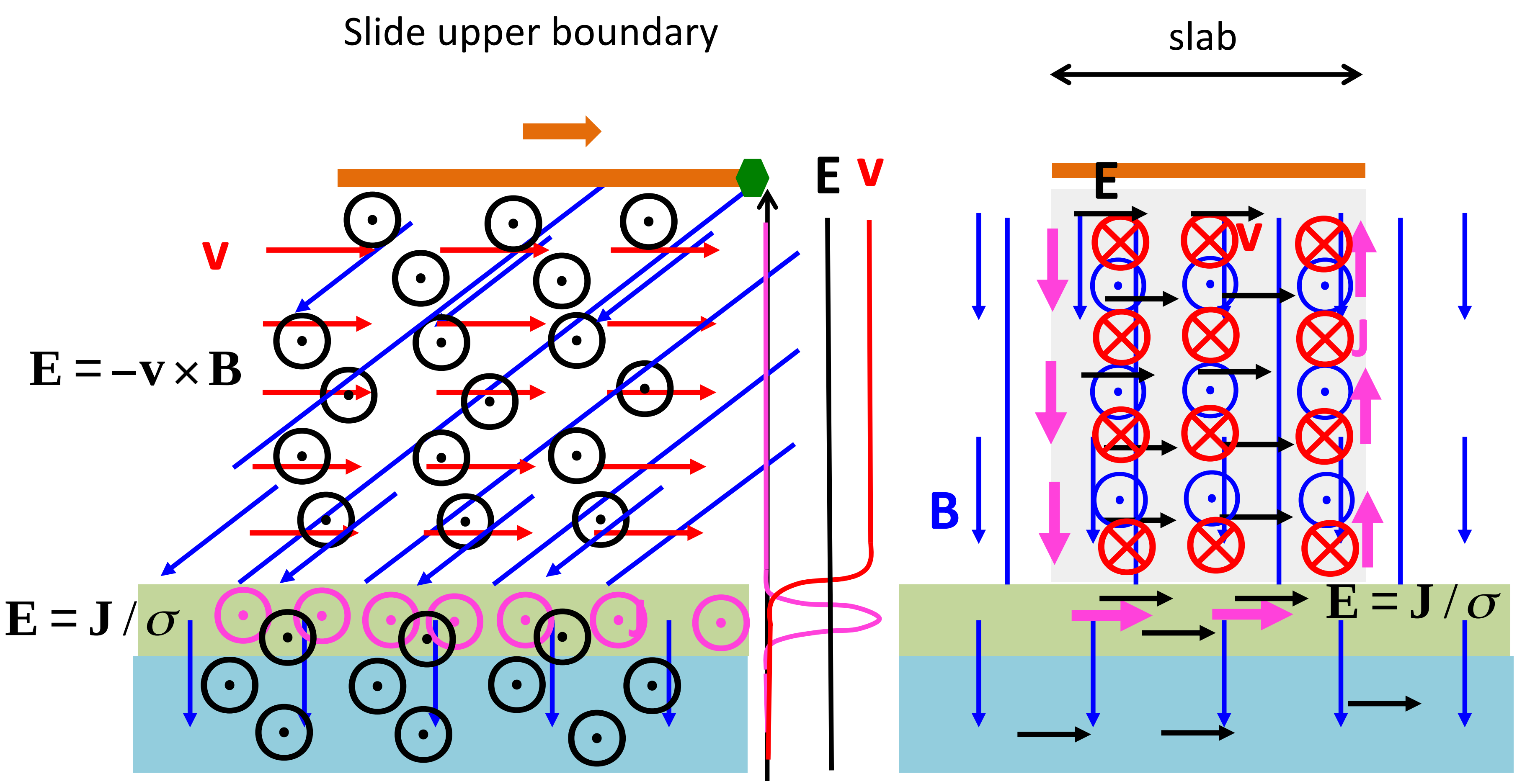
Example of how the motions meet



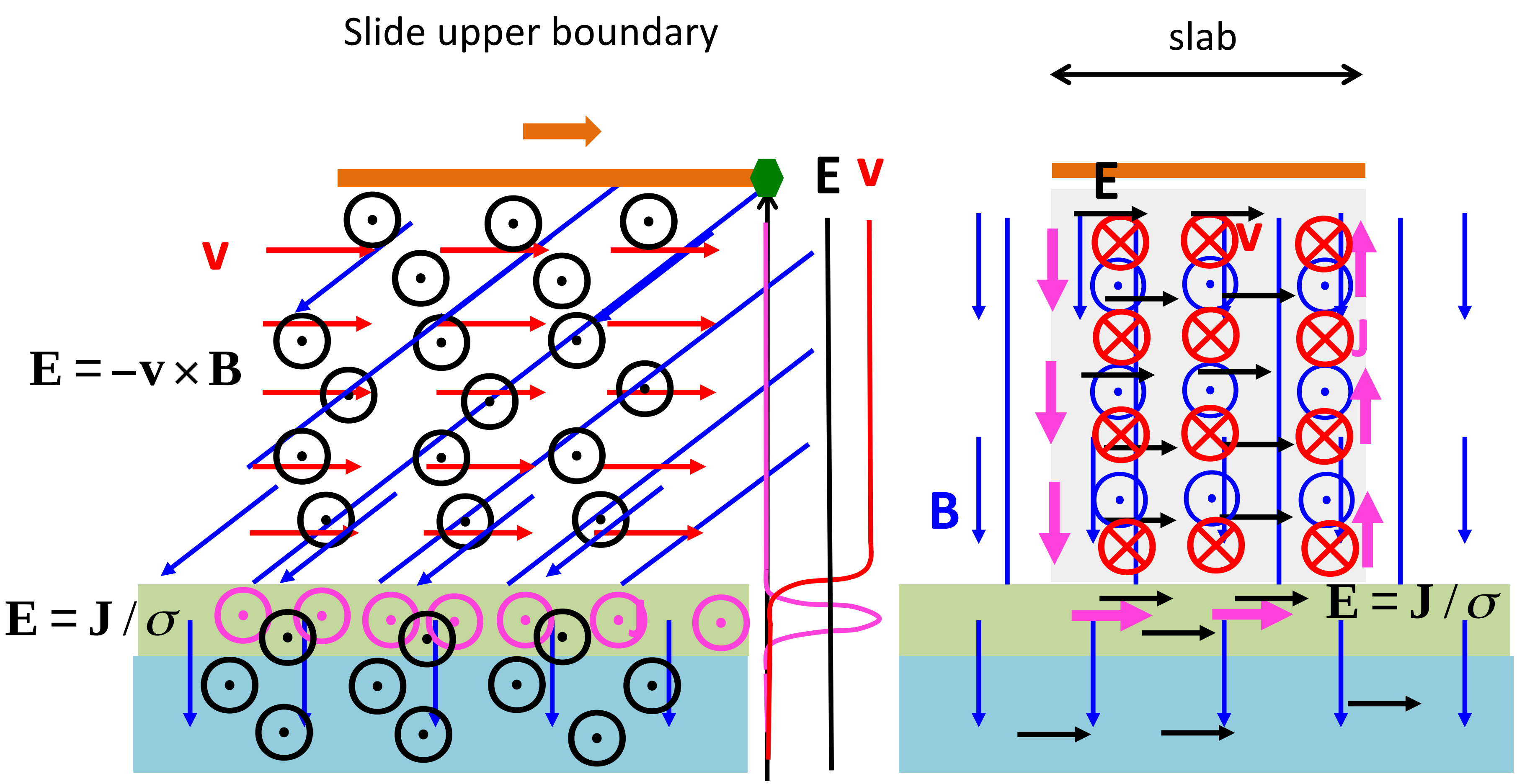
Example of how the motions meet



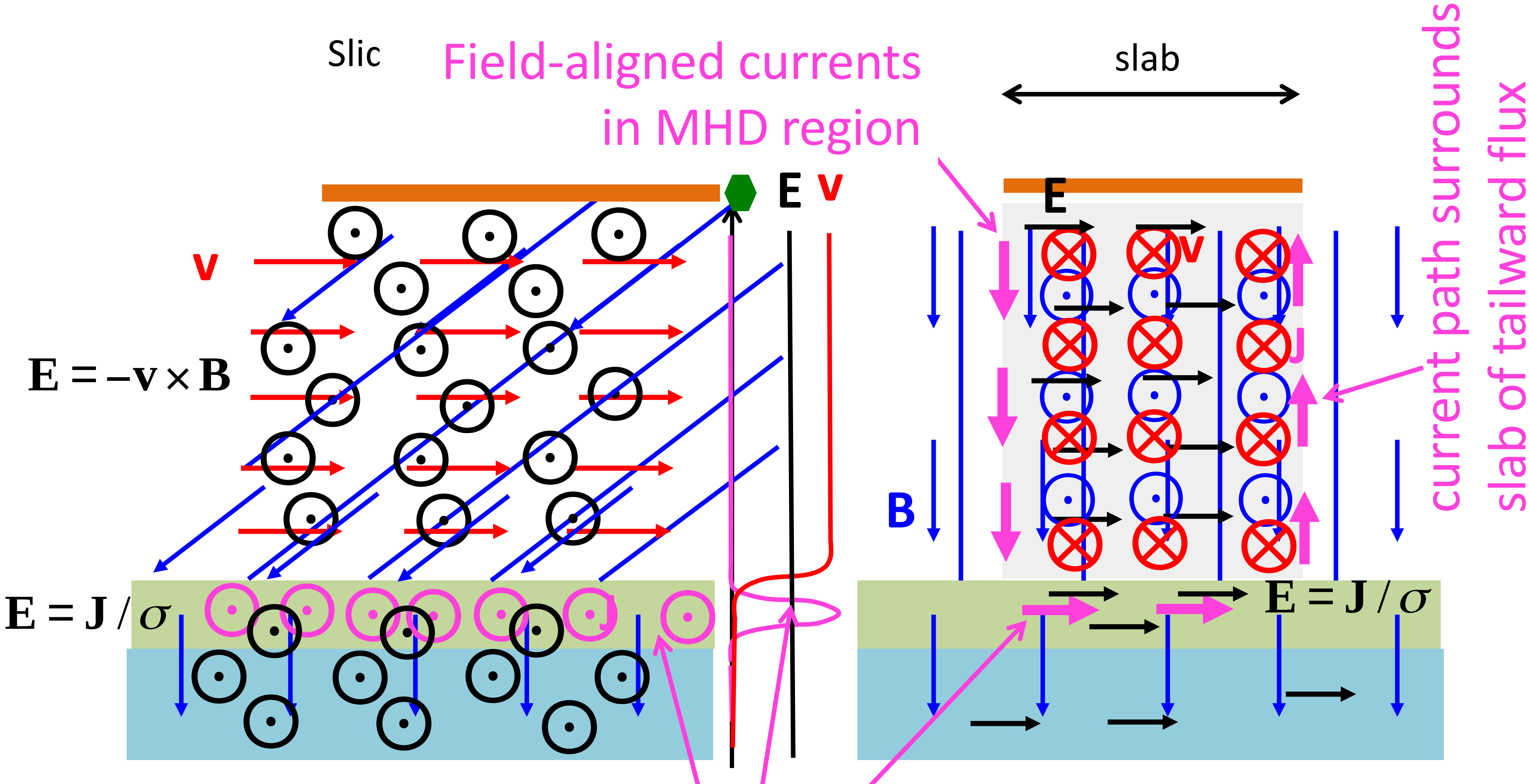
Example of how the motions meet



Example of how the motions meet



Example of how the motions meet



MHD Field line motion creates current in ionosphere – accompanied by **E**

Convection: magnetosphere meets ionosphere

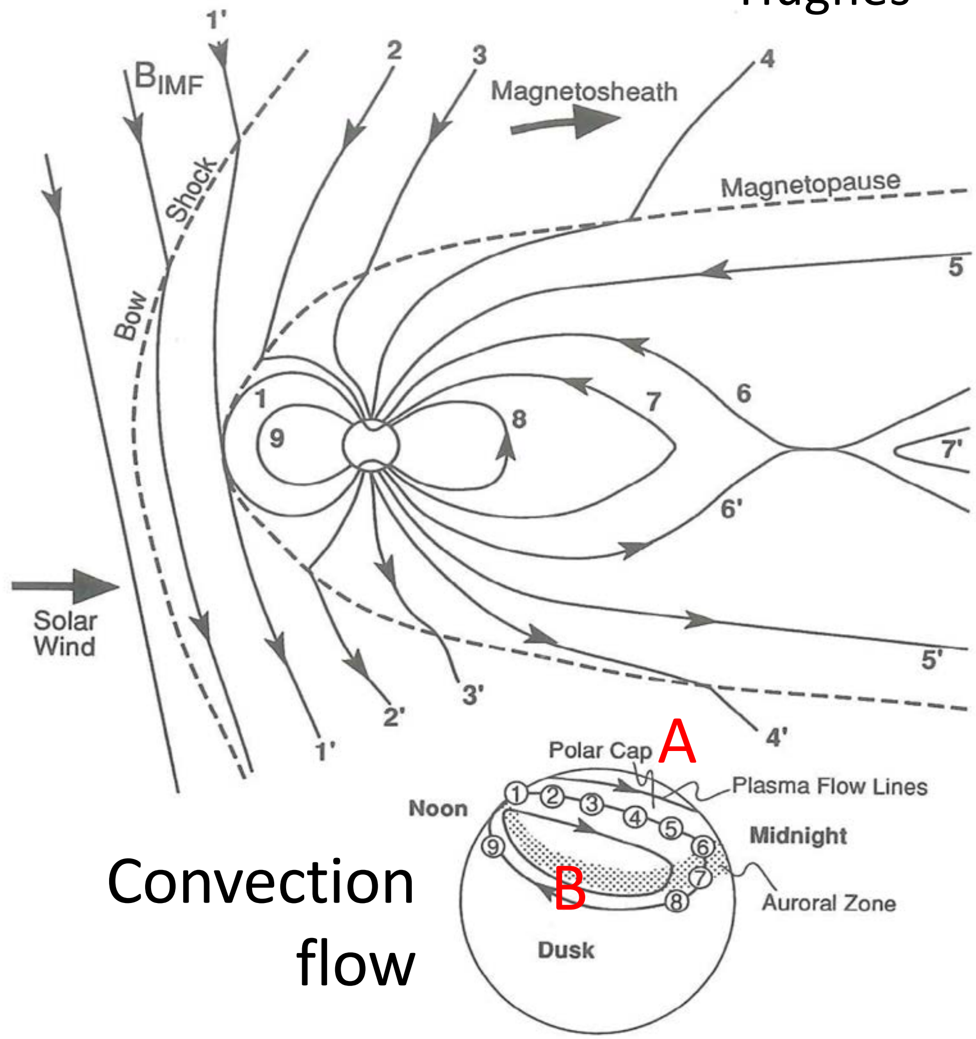
Hughes

MHD motions drag footpoints across polar caps and back around to day side

Integrate* **E** across polar cap:

$$\int_A^B \mathbf{E}\cdot d\mathbf{l} = \varphi_{pc} = \dot{\Phi}_{ds}$$

Really an EMF – but called "cross polar cap potential"



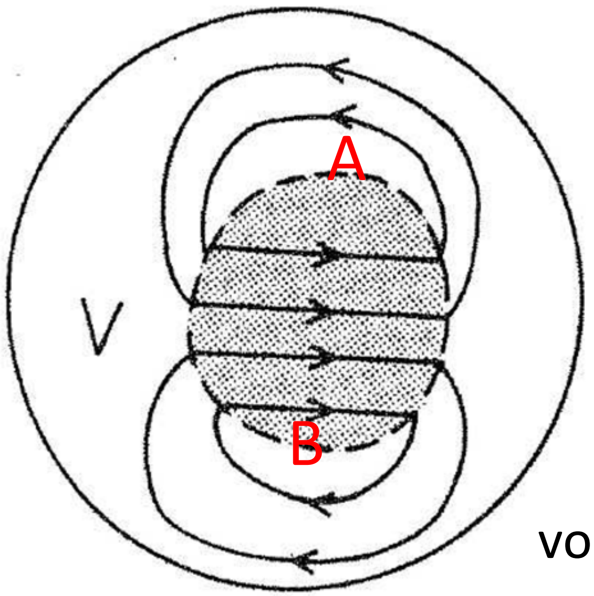
Convection flow

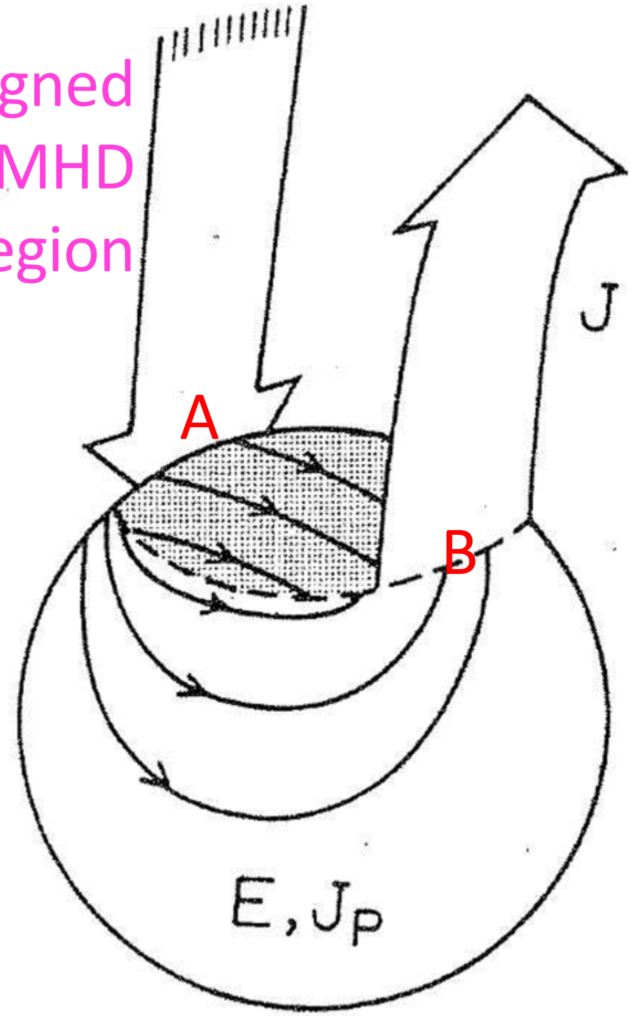
* use MKS here

$$\int_A^B \mathbf{E}\cdot d\mathbf{l} = \varphi_{pc} = \dot{\Phi}_{ds}$$

Field-aligned currents in MHD region

Convection flow





vol. I fig. 10.5

ϕ_{pc} = 50 kV

= 5 × 10^{12} Mx/s

recycle in Φ_t in ~ 5 hours

Summary

- Ionospheres created by EUV & X-rays from Sun's TR and corona
- Diminish during night – lower during solar minimum
- SW deflected by ionospheres of unmagnetized planets (Venus & Mars)
- SW deflected by magnetospheres
- Magnetotail created by reconnection with solar wind magnetic field