

Solar Wind Physics

Marty Lee



USA

University of Chicago 1966

- Gene Parker's solar wind was a story of mythical proportions: it changed and defined "Heliophysics"
- Most rejected the concept of a wind replacing a static atmosphere
- The Solar Breeze was an alternative
- Treating the "solar corpuscular radiation" as a fluid was challenging
- Gene taught my E&M course: tough!
- Parker Solar Probe

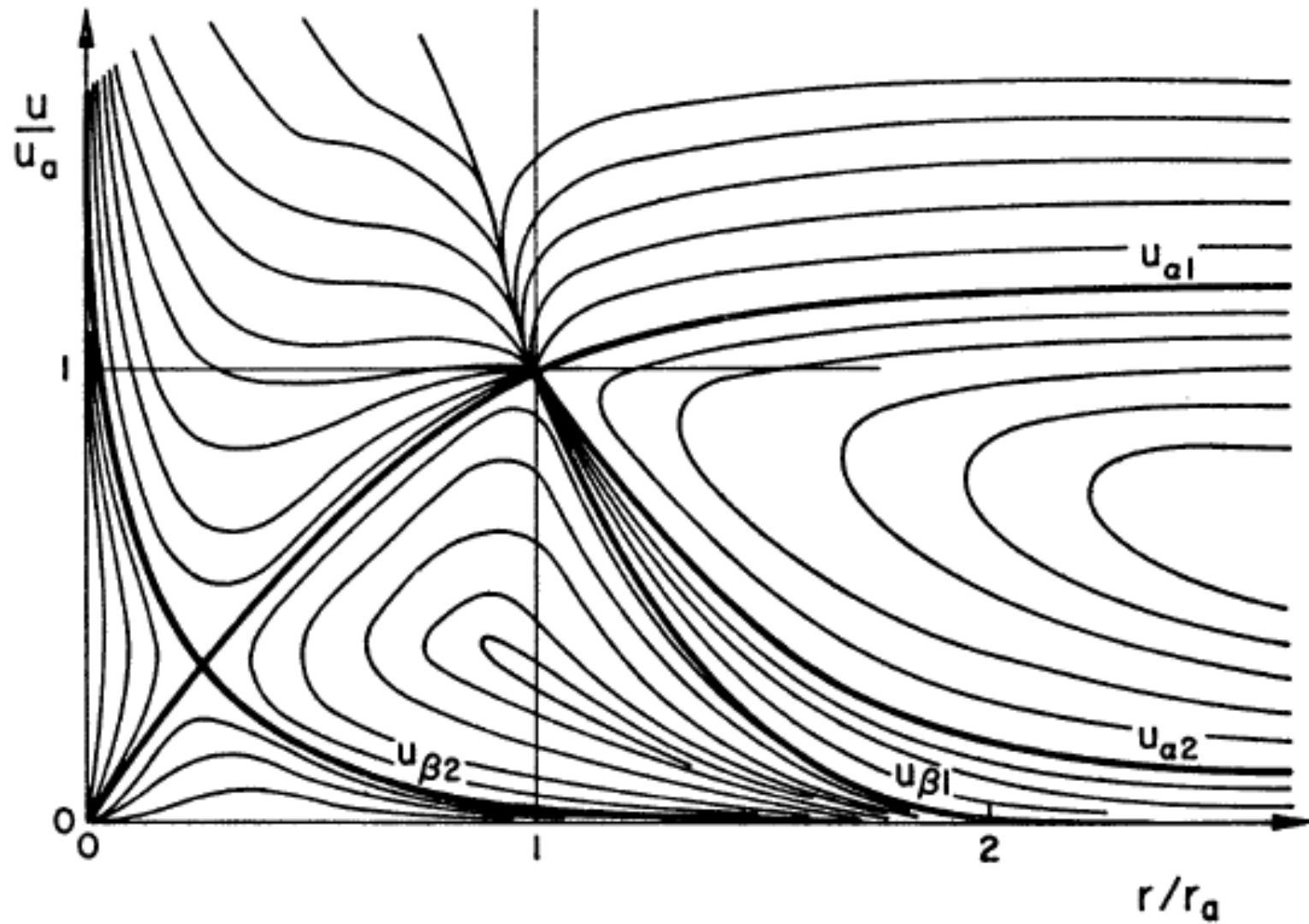
Gene Parker



Some Other Issues

- The concept of the “critical point”
- Bondi accretion
- Magnetic Field and Angular Momentum
- Solar Wind Heating: The weak dependence of $T(r)$
- Complexities of fluid equations – particularly for early space physicists (trained in physics)!

Magnetic Field



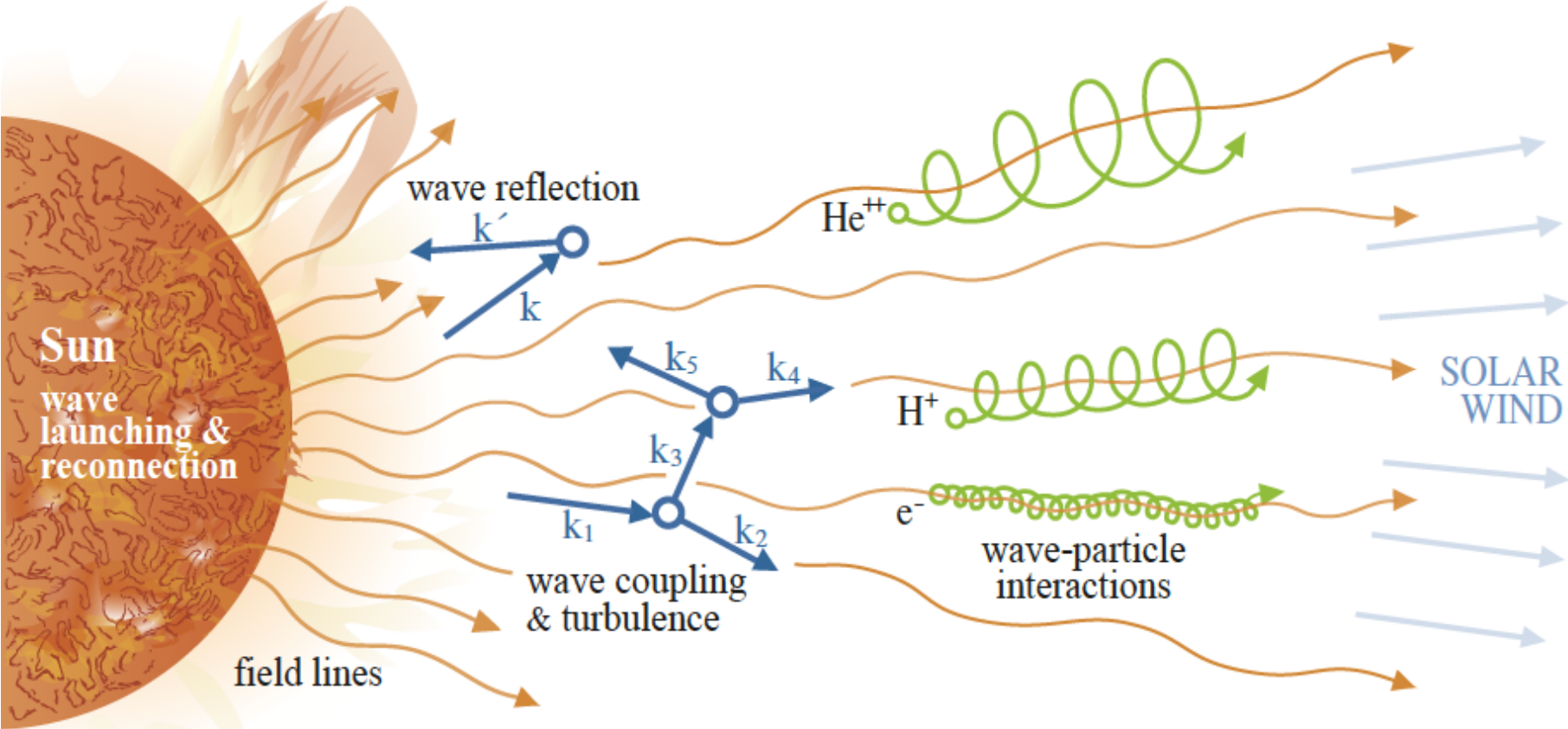
Weber & Davis, 1967

Loss of Angular Momentum

$$L = \Omega r_A^2$$

per unit mass

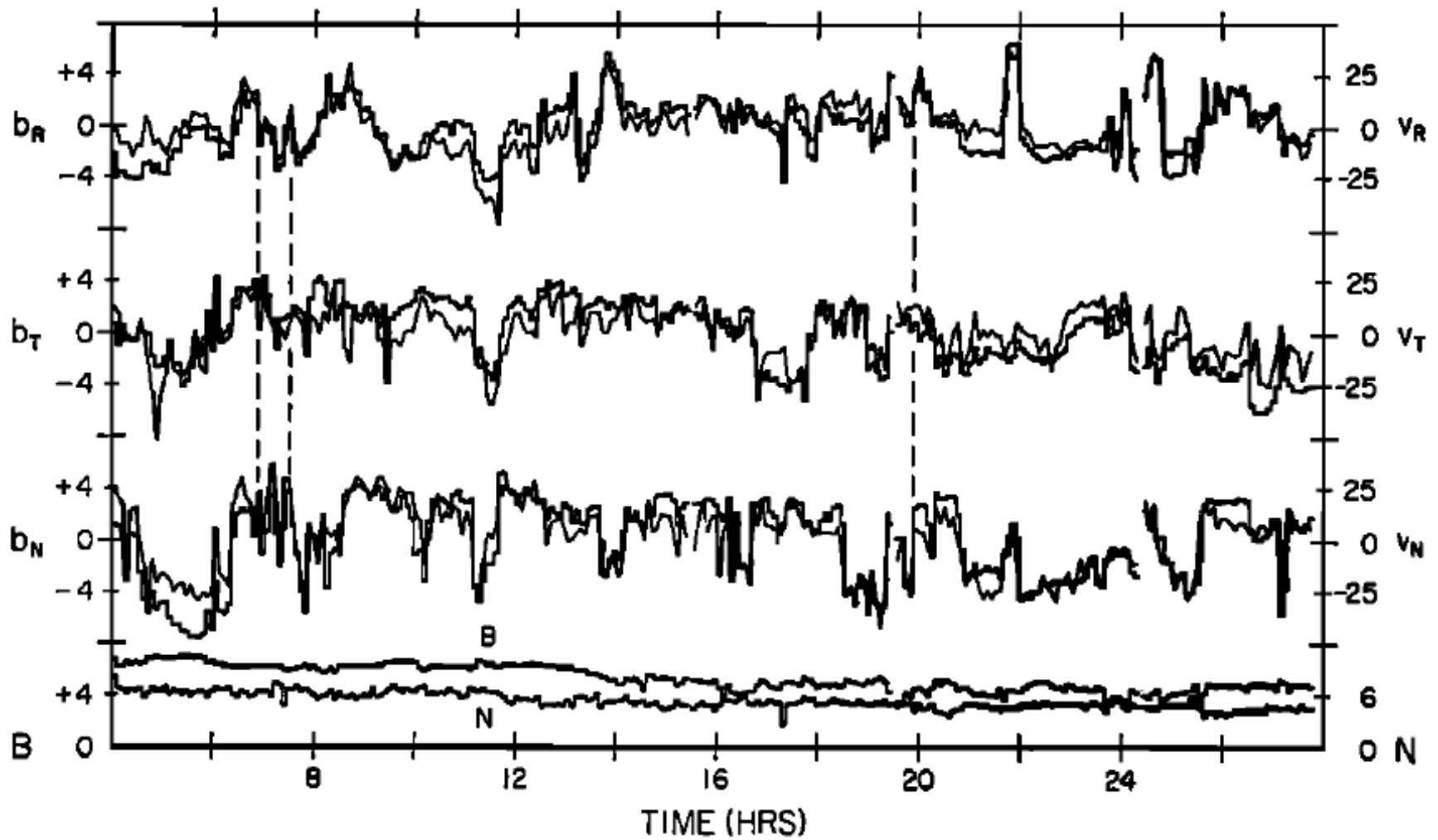
Extended Heating of the Solar Wind



Heating/Acceleration processes in SW

- Turbulent perpendicular heating of ions
- Outward mirroring of ions
- Turbulent wave interactions
- Wave dissipation at higher frequencies
- Electron heat flux (charge separation)
- Magnetic reconnection
- Complex structure of ion distributions
- Spatial structures?

ALFVÉN WAVES IN SOLAR WIND



Belcher and Davis, 1971

Electron Strahl

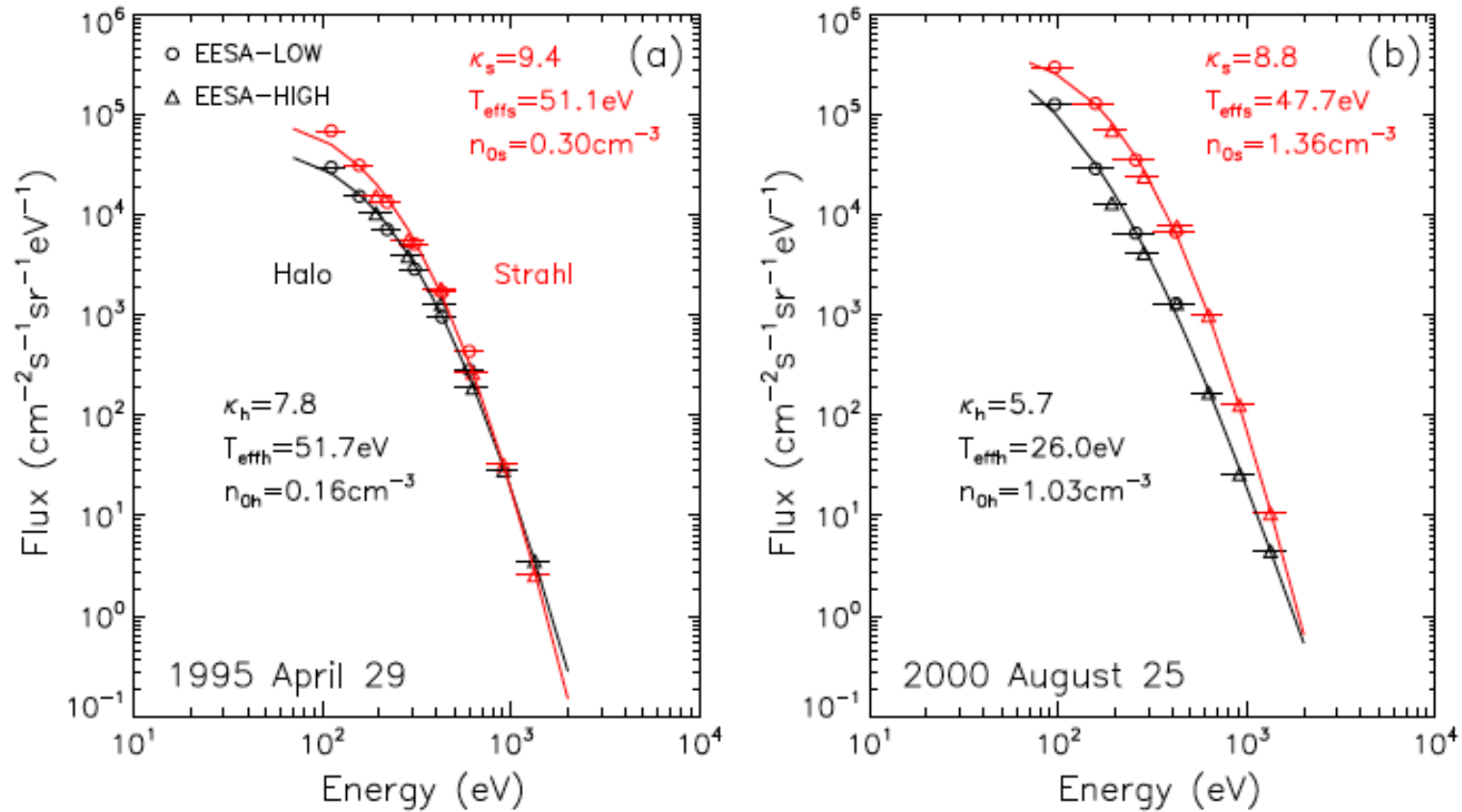
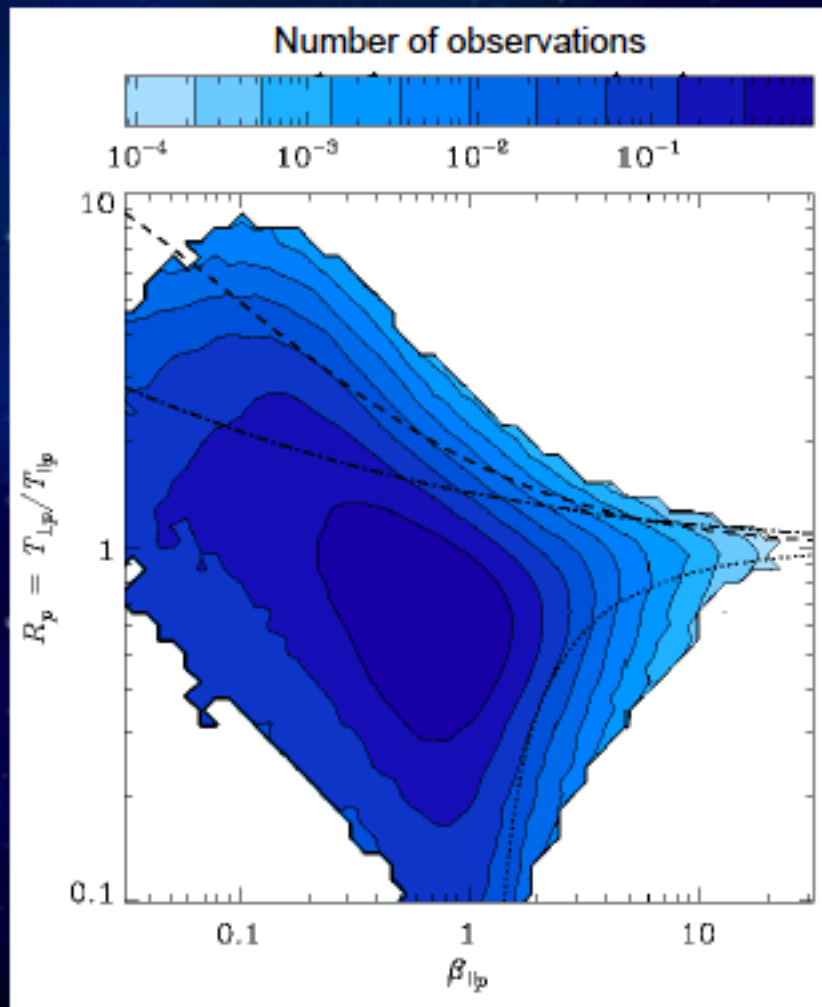


Figure 2. Examples of the energy spectrum of the 0.1–1.5 keV strahl (red) and halo (black) electrons observed by EESA-L (circles) and -H (triangles) on 1995 April 29 (left) and 2000 August 25 (right).

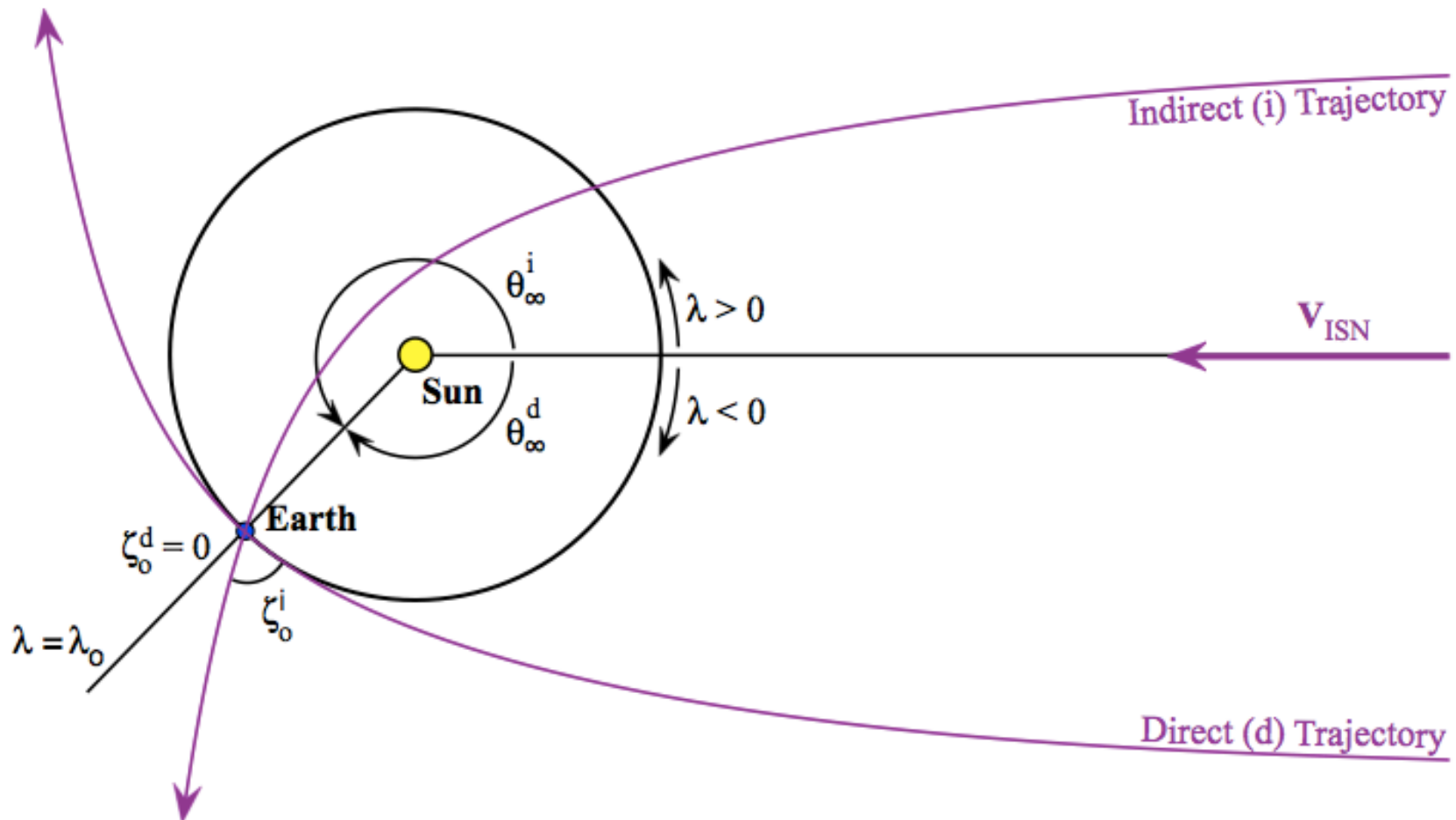
Temperature anisotropy



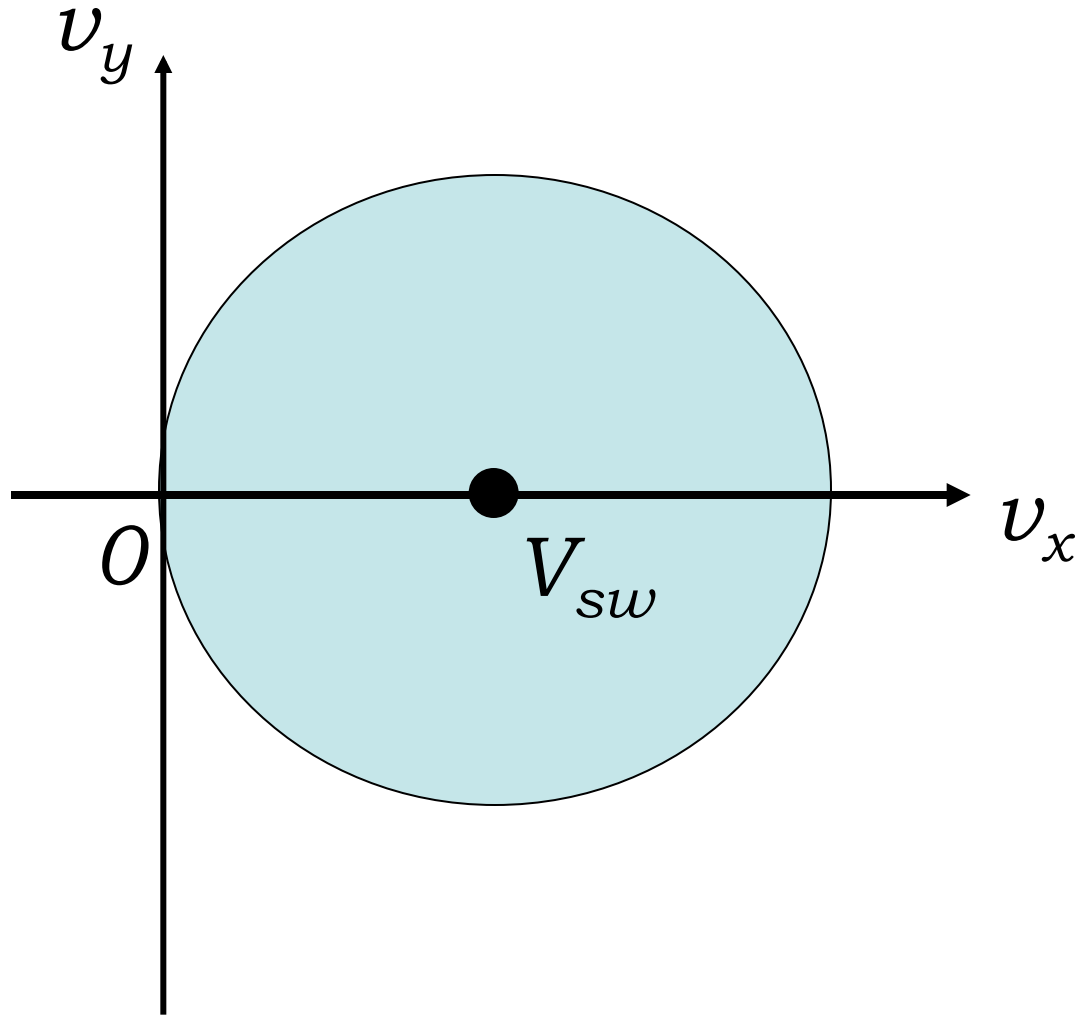
- In the heliosphere
 - Large in the corona
 - Limited by instabilities
- Significance
 - Generates EM fluctuations
 - Modifies particle accel., transport
 - Limits anisotropic heating
- Applied in astrophysics
 - Accretion disks
 - Heat diffusion in galaxy clusters
- Laboratory exp
 - Drives improvements to numerical simulations

Pickup Ions

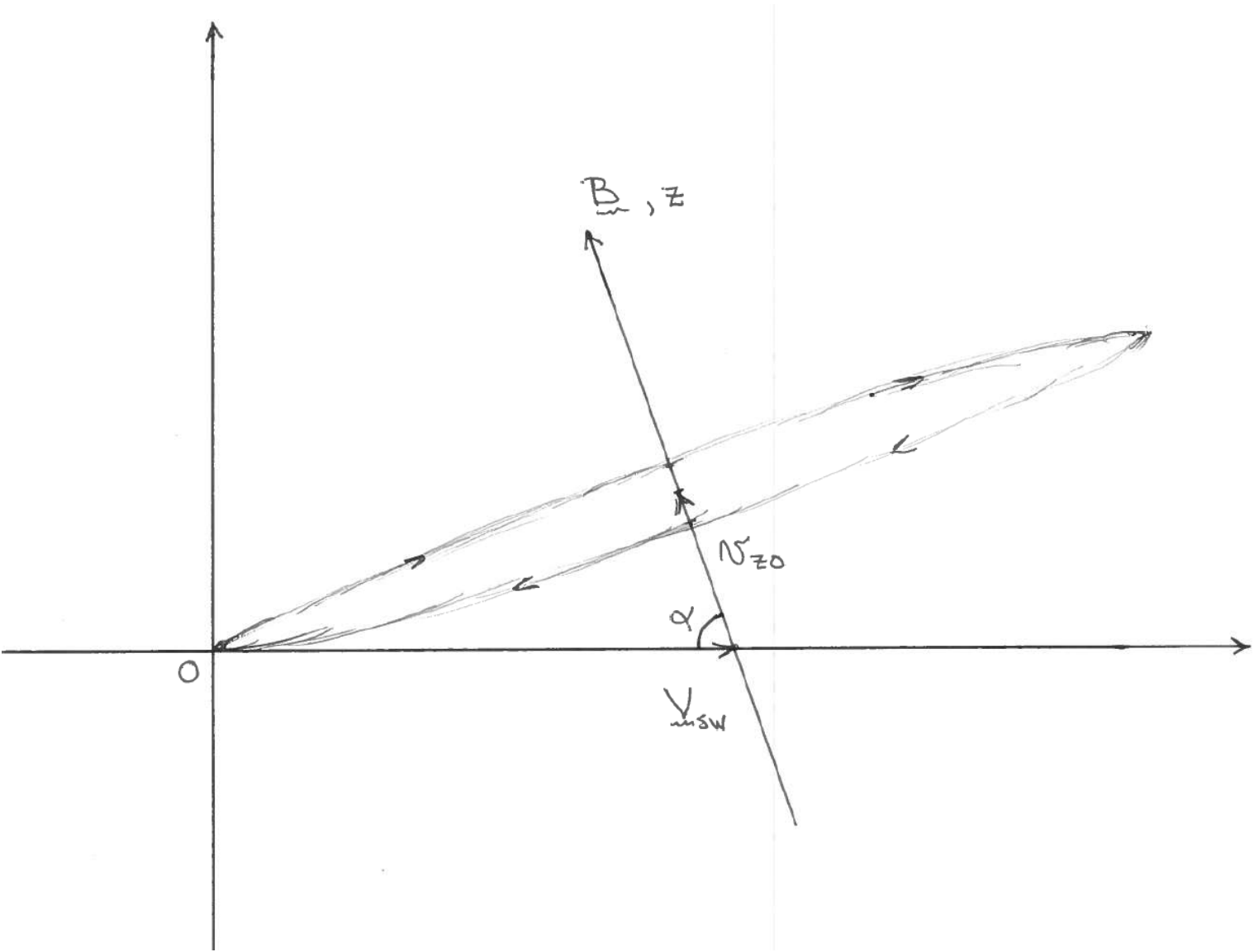
Direct and Indirect Trajectories



Pickup Ion Distribution in SW



Pickup Ion Ring-beam Configuration



Pickup Ion Loading of the Solar Wind I

$$\frac{1}{r^2} \frac{d}{dr} (r^2 \rho V) = m_p \left(\frac{r_0}{r} \right)^2 (v_H^0 n_{H, \infty} + 4v_{He}^0 n_{He, \infty})$$

$$\frac{1}{r^2} \frac{d}{dr} (r^2 \rho V^2) = -\frac{dP}{dr} - \sigma n_{H, \infty} \rho V^2$$

$$\frac{1}{r^2} \frac{d}{dr} \left[r^2 V \left(\frac{1}{2} \rho V^2 + \frac{\gamma}{\gamma - 1} P \right) \right] = -\sigma n_{H, \infty} \frac{1}{2} \rho V^3$$

Pickup Ion Loading of the Solar Wind II

$$P = \frac{1}{2} \frac{\gamma - 1}{2\gamma - 1} \frac{r_0^2}{r} \rho_0 V_0 \alpha$$

$$\alpha \equiv \sigma n_{H, \infty} V_0 + m_p \rho_0^{-1} (v_H^0 n_{H, \infty} + 4v_{He}^0 n_{He, \infty})$$

$$C_s^2 = \frac{1}{2} \gamma \frac{\gamma - 1}{2\gamma - 1} V_0 \alpha r$$

$$V = V_0 - \left[1 - \frac{1}{2} \frac{\gamma - 1}{2\gamma - 1} \right] \alpha r$$

Interstellar Pickup Ion Transport

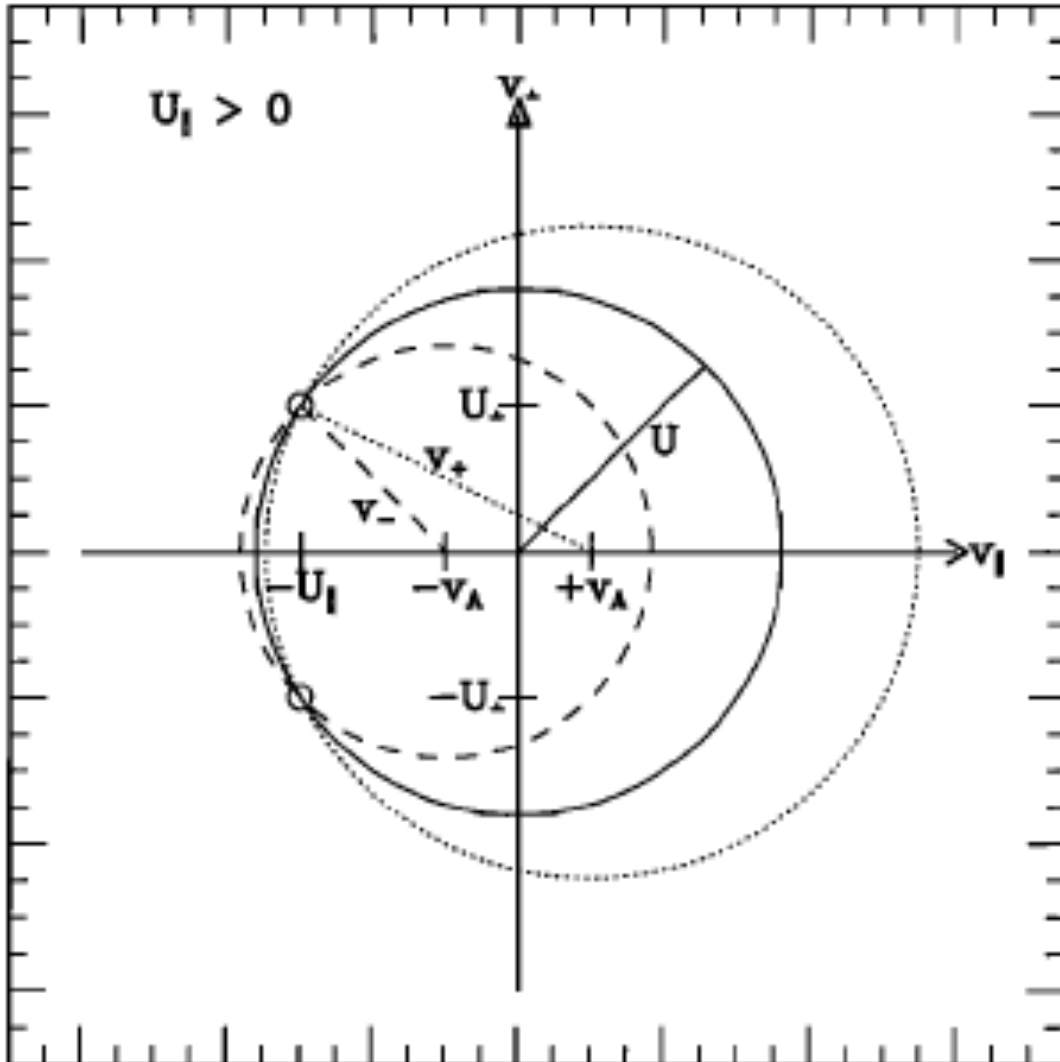
$$\mathbf{V}_D \cong 0, \mathbf{K} \cong 0, \mathbf{V} = \mathbf{e}_r V, \partial / \partial t = 0$$

$$\frac{\partial f}{\partial t} + \mathbf{V} \cdot \nabla f - \frac{1}{3} \nabla \cdot \mathbf{V} v \frac{\partial f}{\partial v} = \beta_0 \left(\frac{r_0}{r} \right)^2 n_g(\mathbf{x}) \frac{\delta(v - V)}{4\pi v^2}$$

$$f(r, v < V) = \frac{3\beta_0 r_0^2}{8\pi V^{5/2}} \frac{1}{rv^{3/2}} n_g \left[r(v/V)^{3/2}, \theta, \phi \right]$$

Exercise 4 in Heliospheric Problems (M. Lee)

Why are waves excited?



Williams & Zank, 1994

Quasilinear Theory

$$V_A / v \ll 1$$

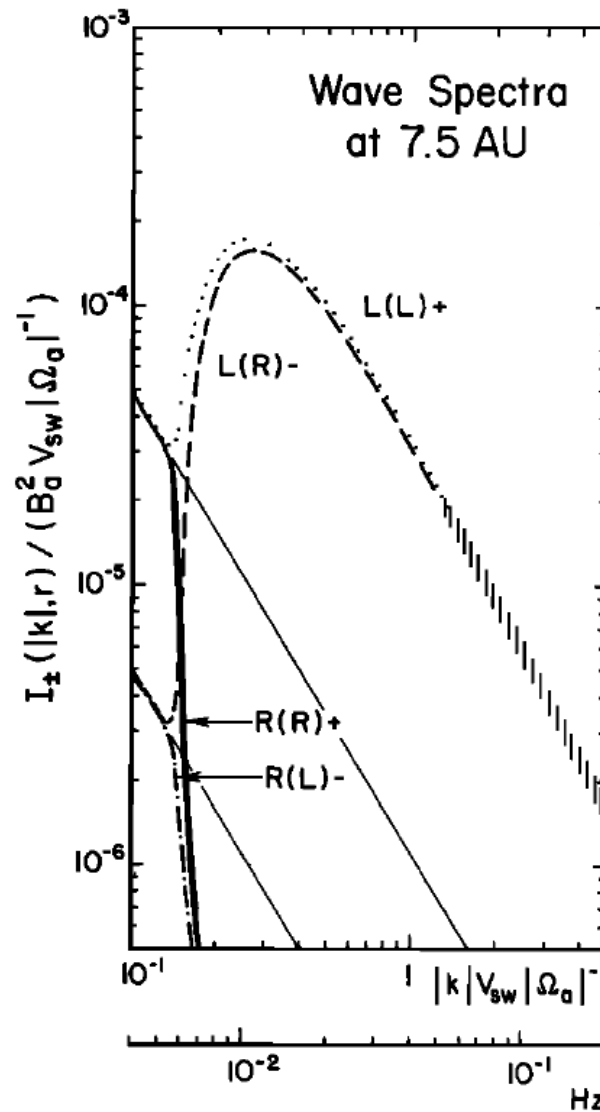
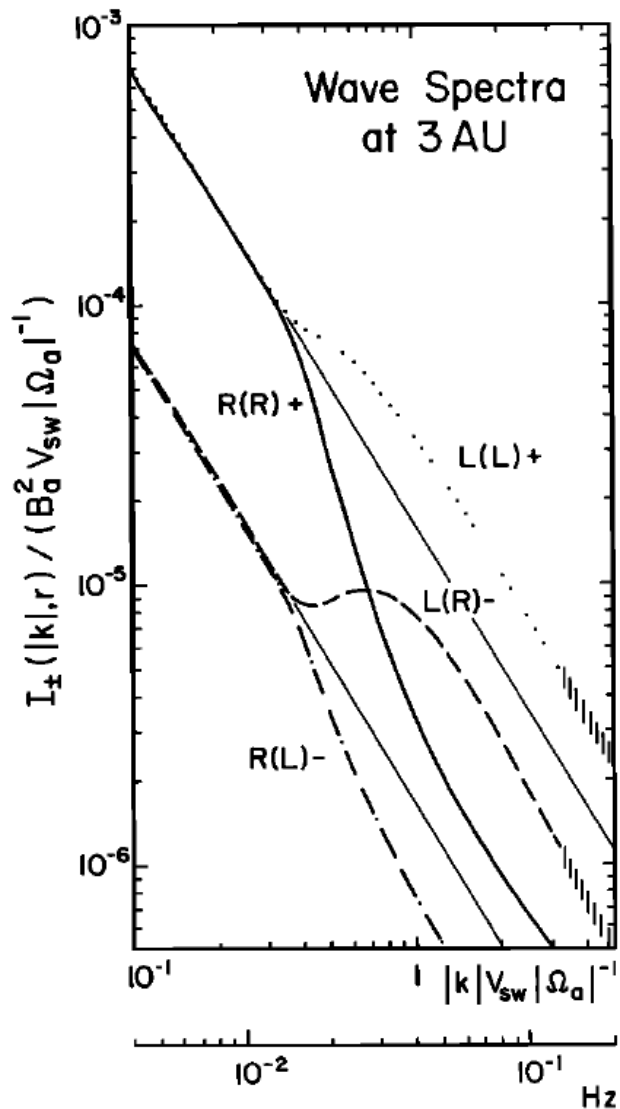
$$\frac{\partial F}{\partial t} = \frac{\partial}{\partial \mu} \left\{ A(\mu) \left[I_+ \left(\frac{\Omega}{v\mu} \right) + I_- \left(\frac{\Omega}{v\mu} \right) \right] \frac{\partial F}{\partial \mu} \right\} \quad A(\mu) = \frac{\pi}{2} \left(\frac{q}{mc} \right)^2 (1 - \mu^2) \frac{1}{v_0 |\mu|}$$

$$\frac{\partial I_{\pm}}{\partial t} = \pm I_{\pm} B(|k|) \int dv d\mu (v/v_0)^2 (1 - \mu^2) \delta \left(\mu - \frac{\Omega}{kv} \right) \quad B(|k|) = \frac{4\pi^3 V_A q^2 v_0^2}{|k| c^2 m}$$

$$\Rightarrow I_{\pm}(k, \infty) = (1/2) \left\{ [C(k)]^2 + 4I_+(k, 0)I_-(k, 0) \right\}^{1/2} \pm (1/2)C(k)$$

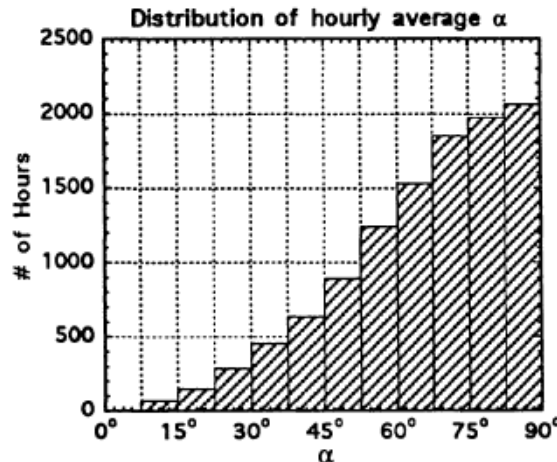
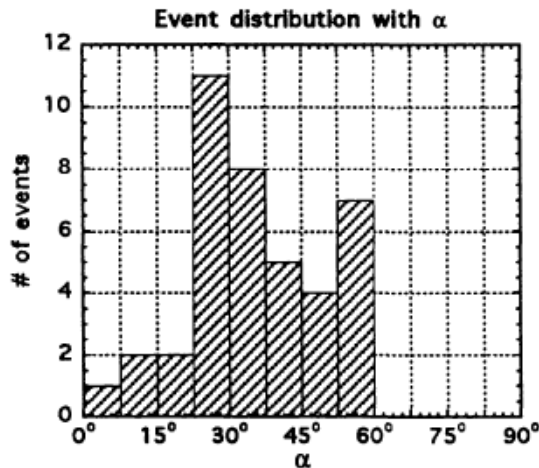
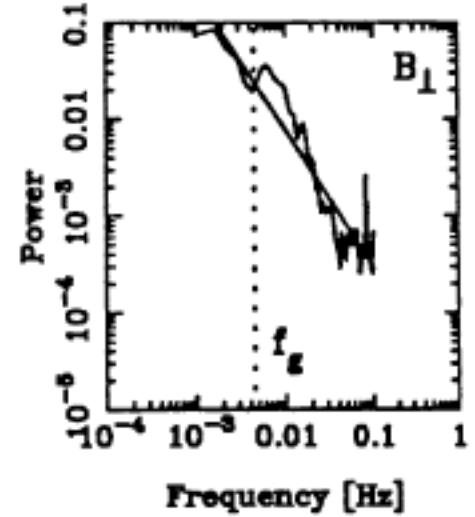
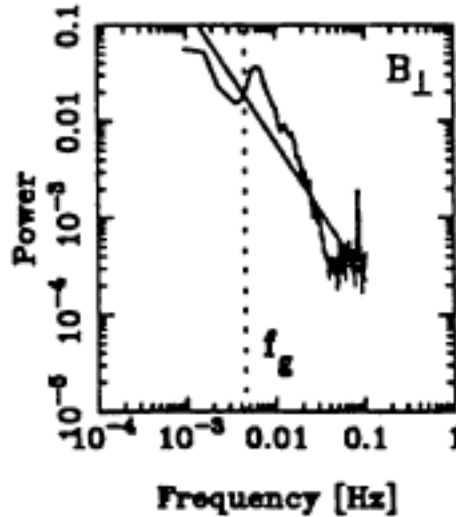
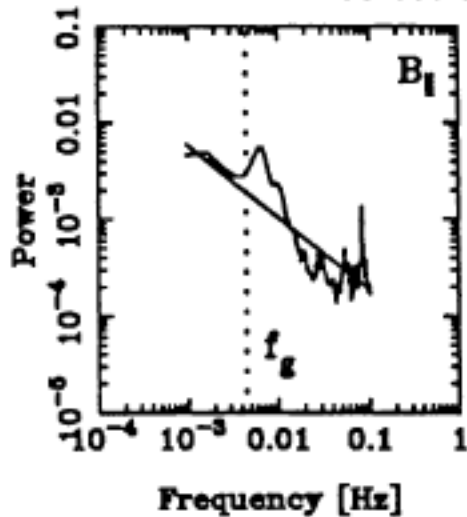
$$C(k) = I_+(k, 0) - I_-(k, 0) + 2\pi N V_A \frac{|\Omega|}{k^2} \left[\frac{\Omega}{kv_0} - \frac{\Omega k^{-1} v_0^{-1} - \mu_0}{|\Omega k^{-1} v_0^{-1} - \mu_0|} \right]$$

PUI Excited Waves With WKB Evolution

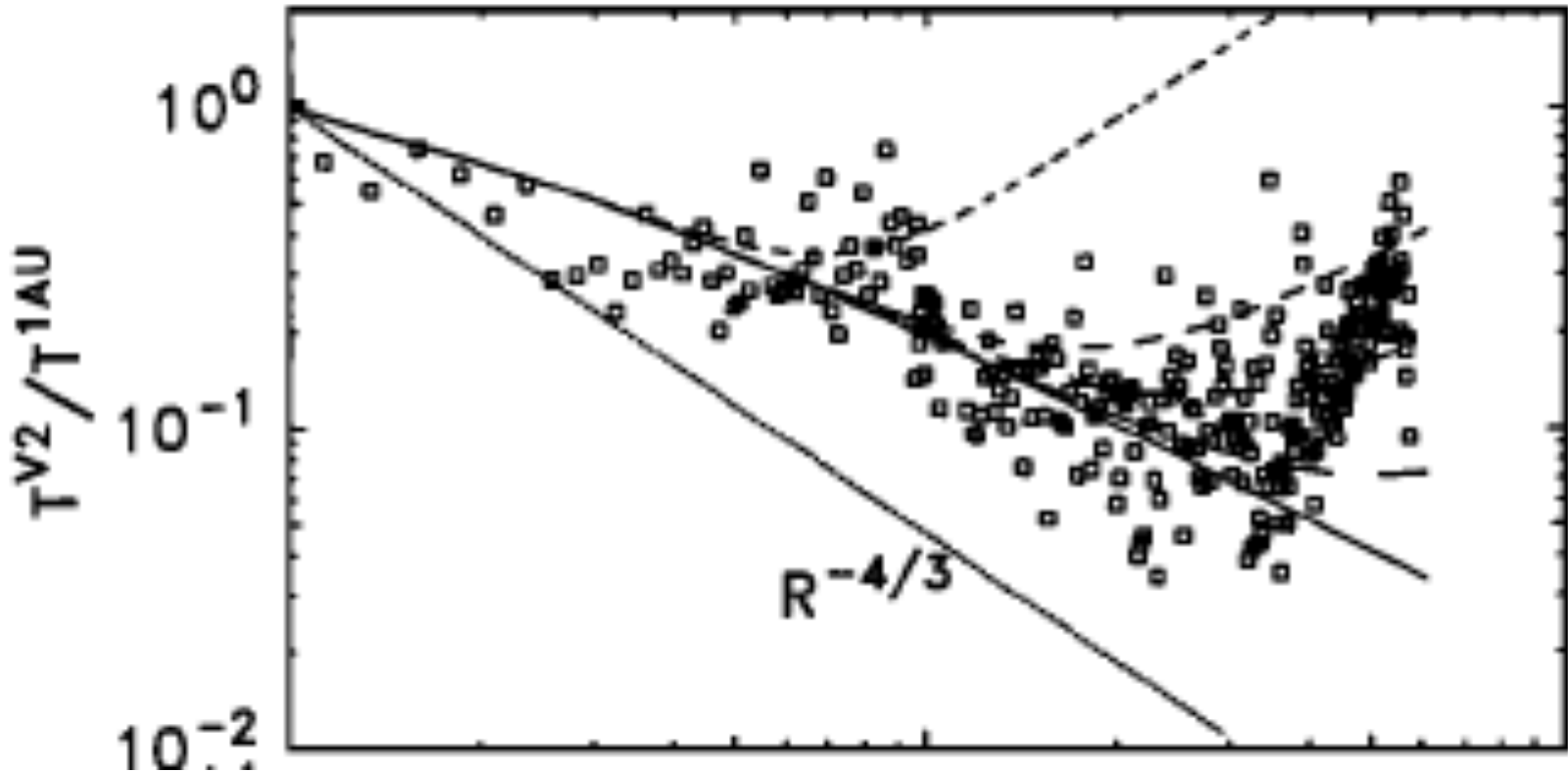


31 Ulysses PU-Proton Events Finally Detected!

93 006 JAN 6 09:00:00.000 to 93 006 JAN 6 11:00:00.000 3.0 #



The Role of Turbulence



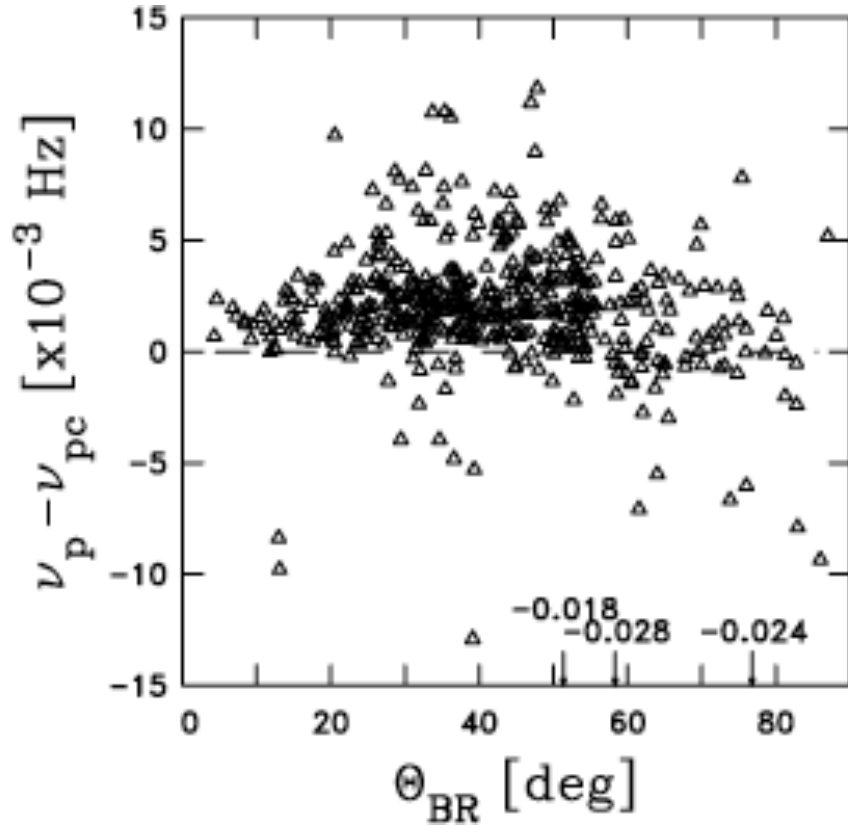
Smith, Matthaeus, Zank, Ness, Oughton & Richardson, 2001

Isenberg, Smith & Matthaeus, 2003

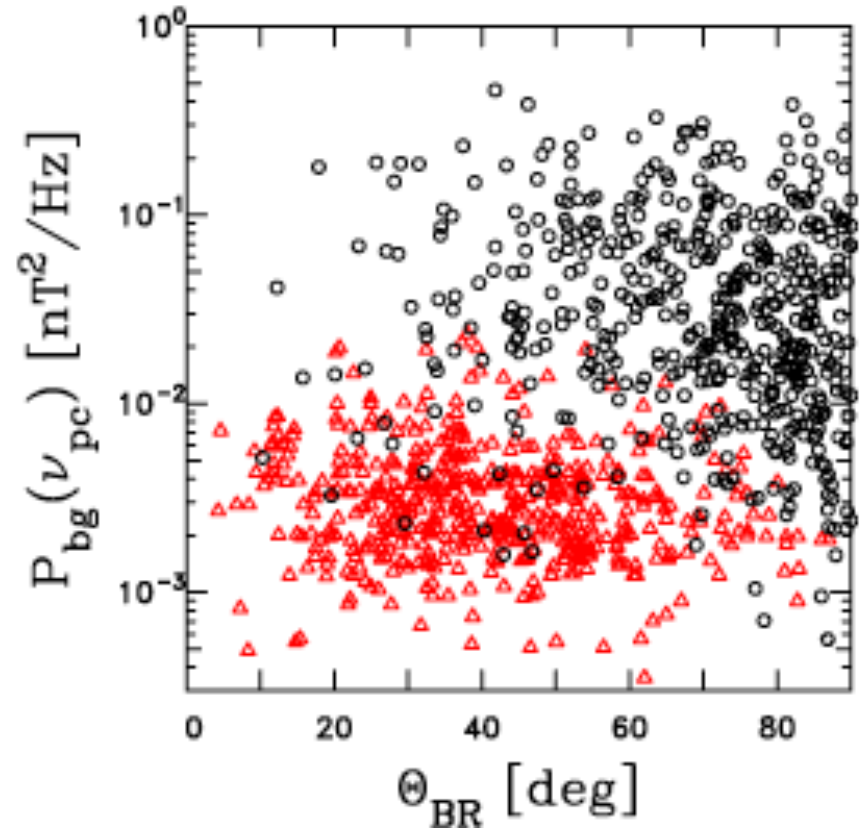
Isenberg, 2005

Isenberg, Smith, Matthaeus, & Richardson, 2005

PUI Wave Properties II



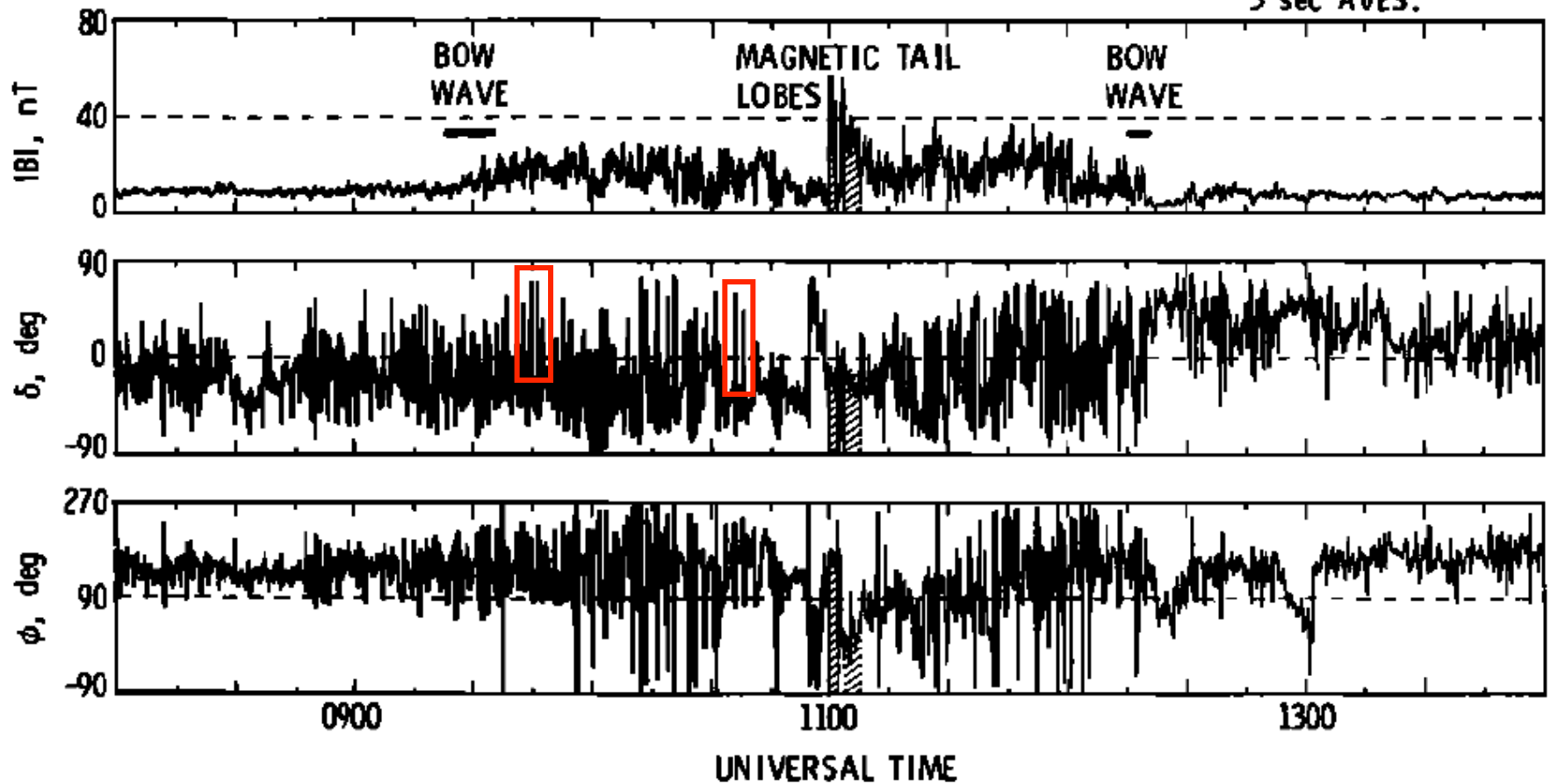
Cannon et al., 2014a



Cannon et al., 2014b

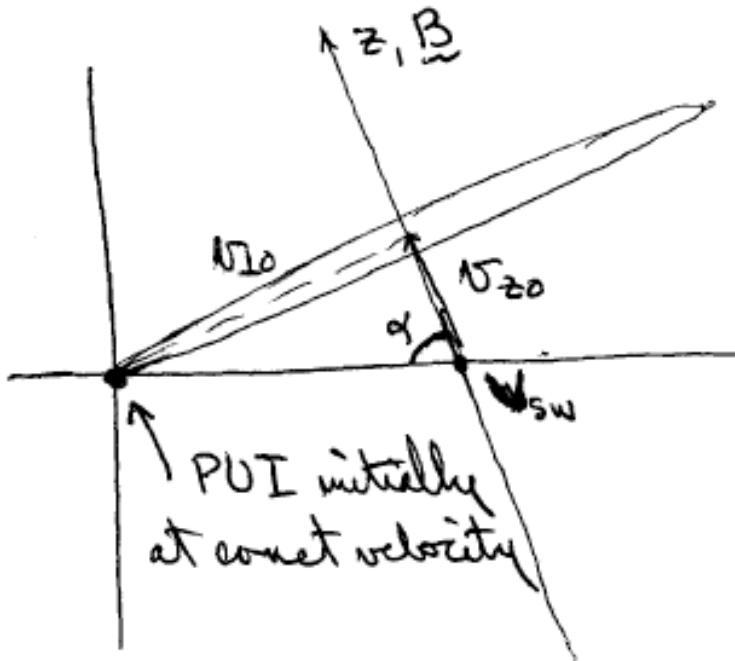
Pickup Ion Excited Waves at Comet G-Z

SEPTEMBER 11, 1985
GSE COORDINATES
3 sec AVES.



Tsurutani and Smith, 1986

Cometary Ion Pickup



$$\omega - k_z v_z + \Omega = 0$$

$$\omega - k_z V_{SW} \cos \alpha + \Omega = 0$$

$$\omega + \mathbf{k} \cdot \mathbf{V}_{SW} + \Omega = 0$$

$$\omega_{sc} = -\Omega$$

$$T_{sc} = \frac{2\pi}{|\omega_{sc}|} = \frac{2\pi}{|\Omega_p|} A$$

$$B = 8 \text{ nT}$$

$$T_{sc} = 8 A \text{ sec} \approx 120 \text{ sec} \Rightarrow A \approx 15$$

**Shock Waves:
Difficult to Avoid in a Supersonic Flow**

Water Ripples



Porthtowan © Michael Regan 2011

Fundy Tidal Bore



Fundy Tidal Bore



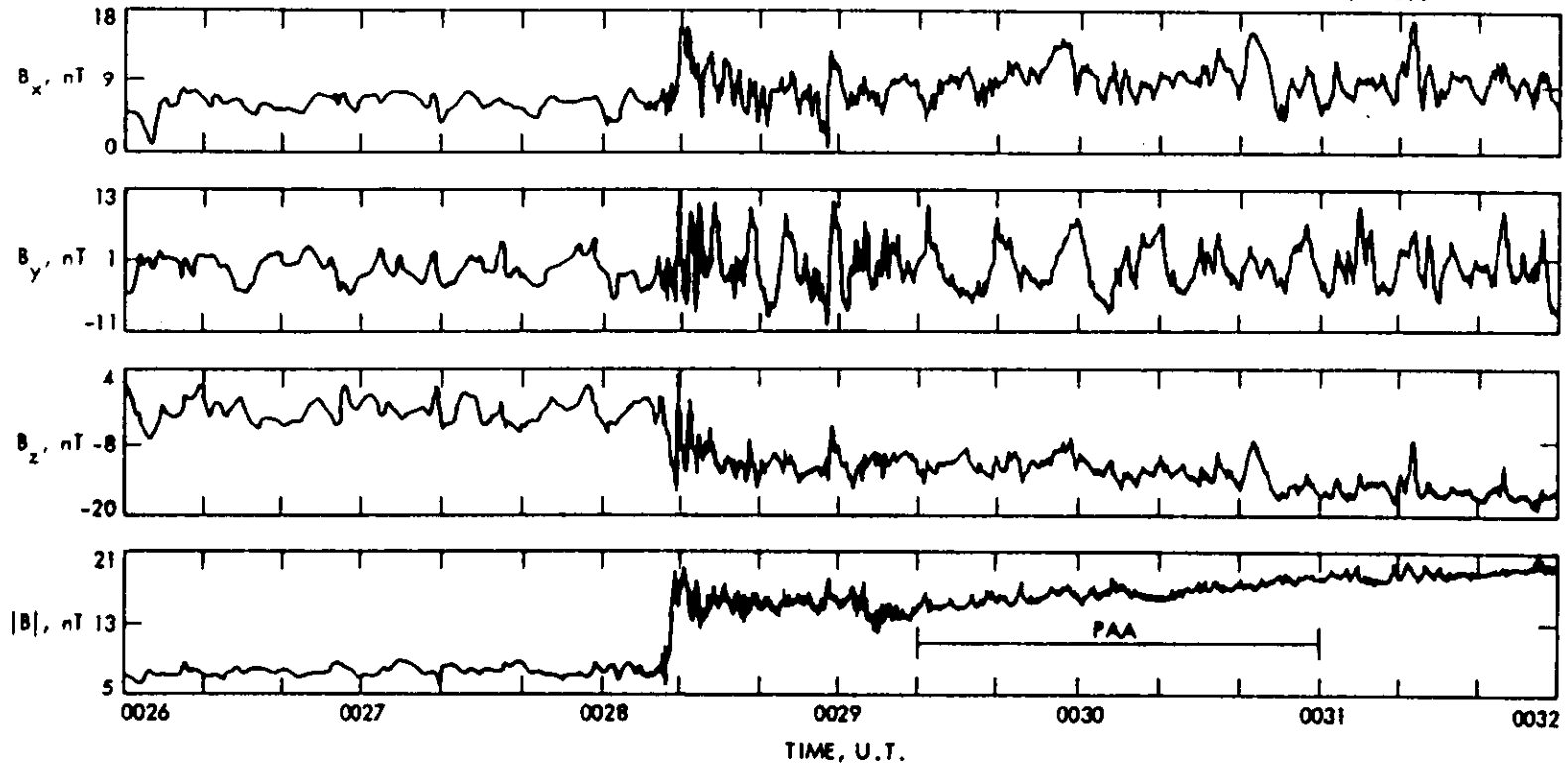
Tsunami Hits Japan, 2011



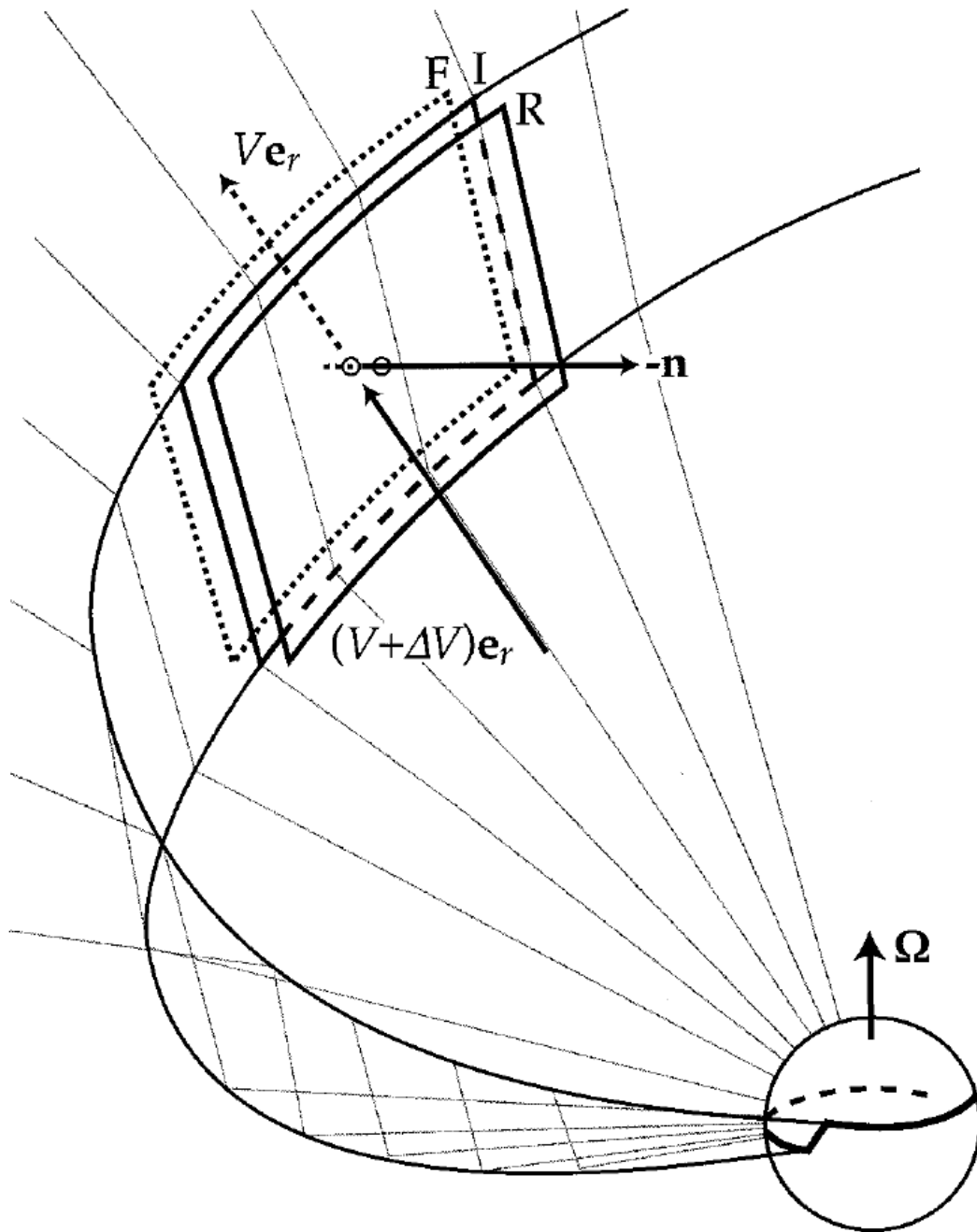
Upstream Waves I

DAY 316, 1978
NOVEMBER 12
ISEE-3

$\hat{n} = (-.96, .28, .09)$
 $\theta_B = 22^\circ$
 $M_s = 4.7$
 $\beta = 0.5$



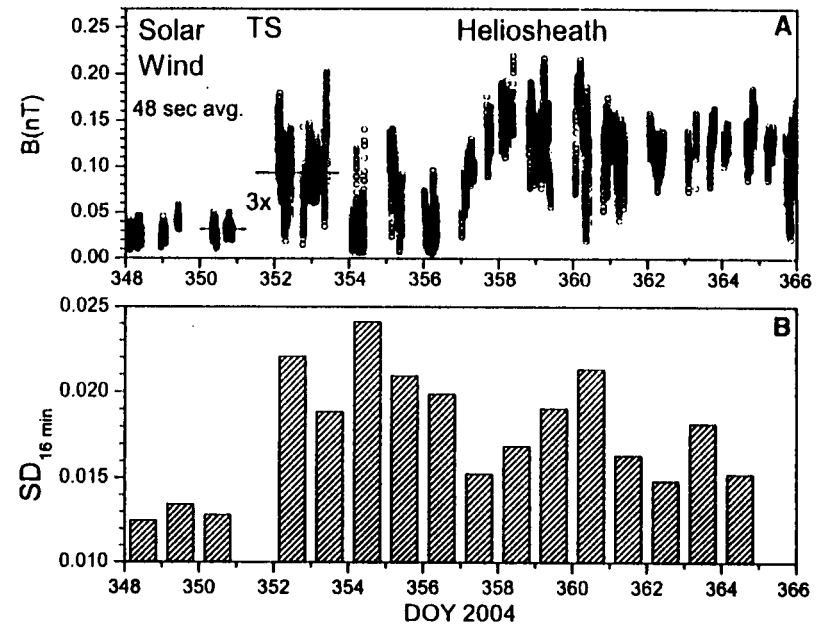
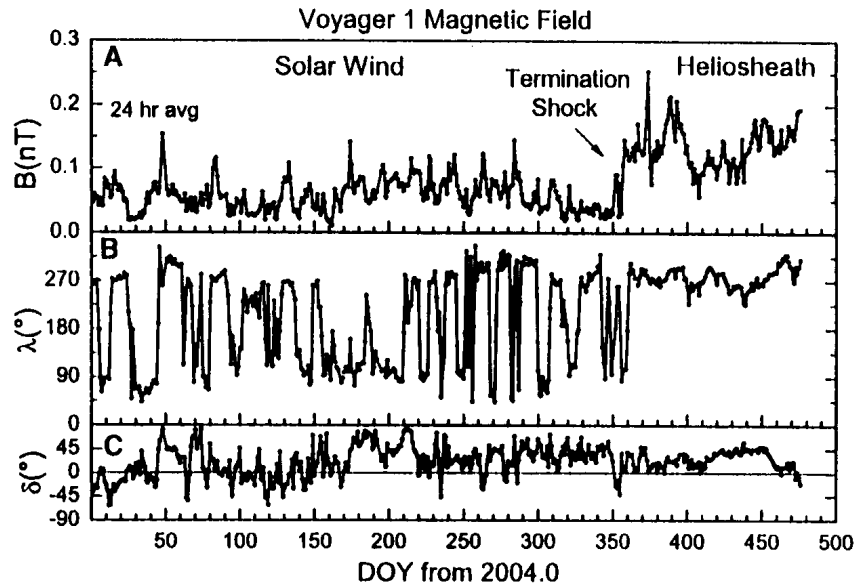
Tsurutani et al., 1983



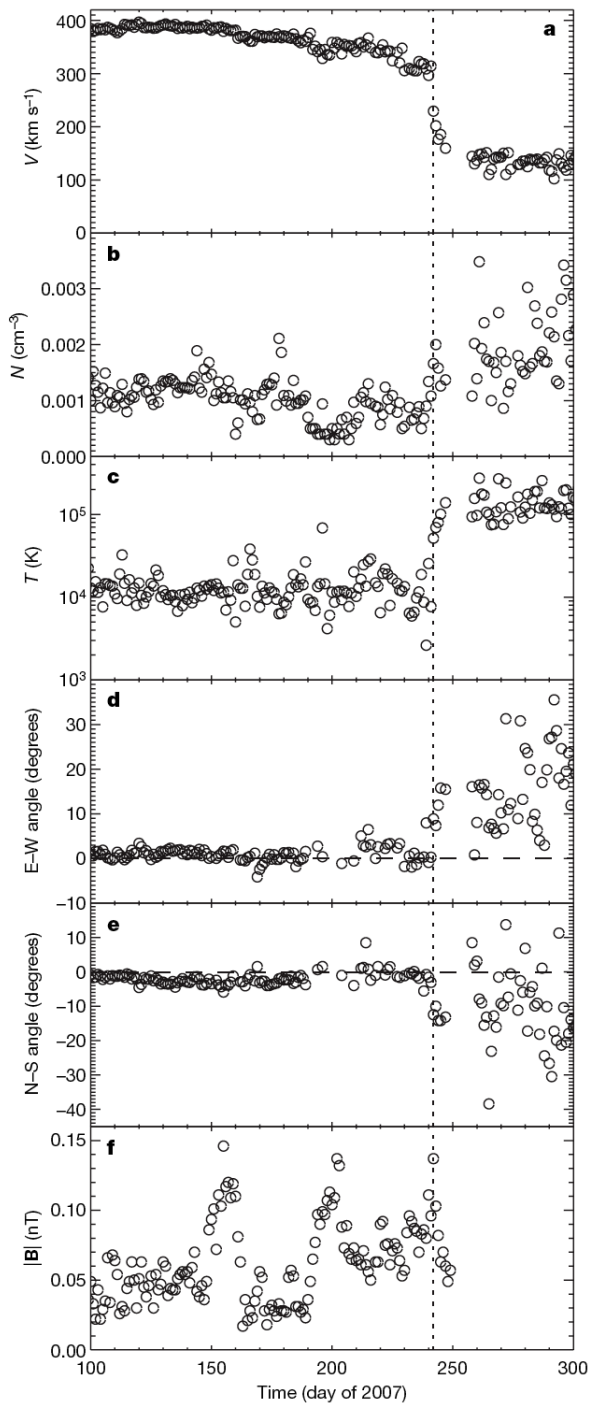
CIR Geometry

Termination of the Solar Wind

Voyager 1 at the Termination Shock



Termination Shock Plasma Data At Voyager 2

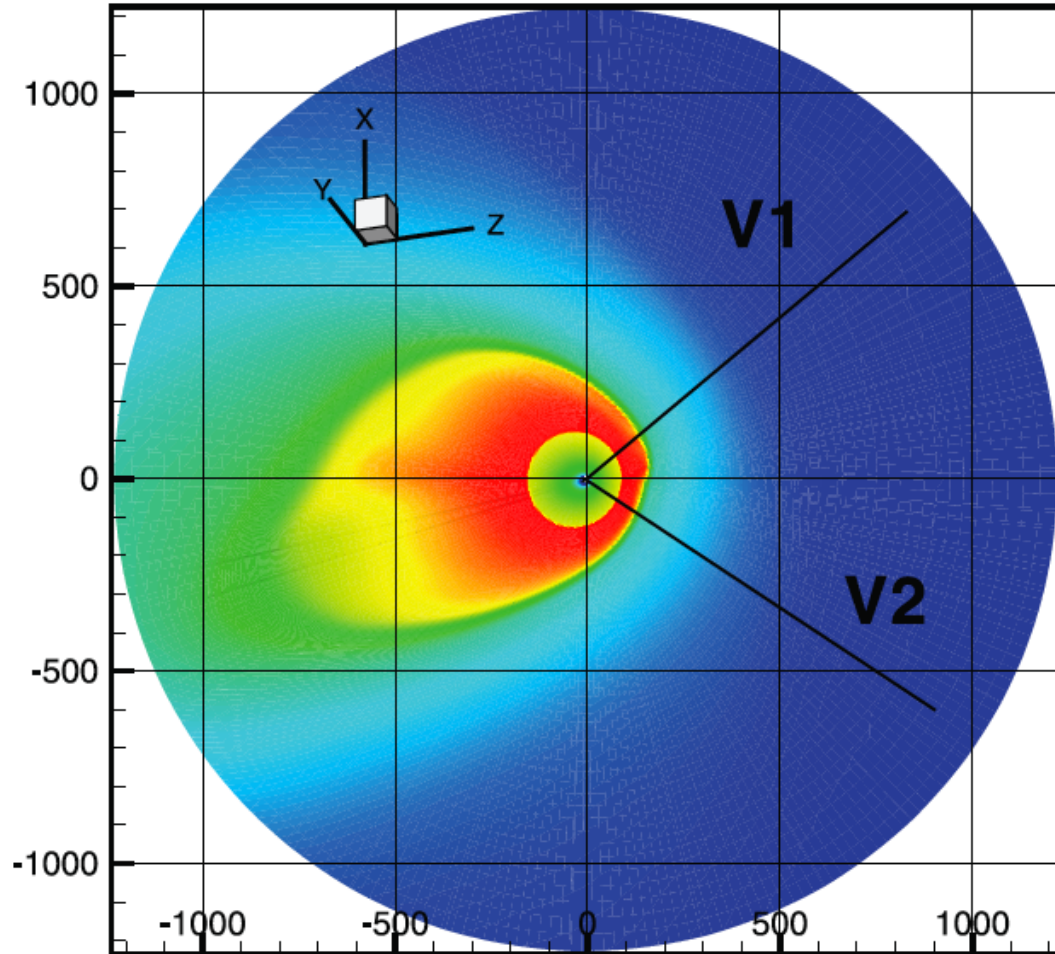


Richardson et al., 2008

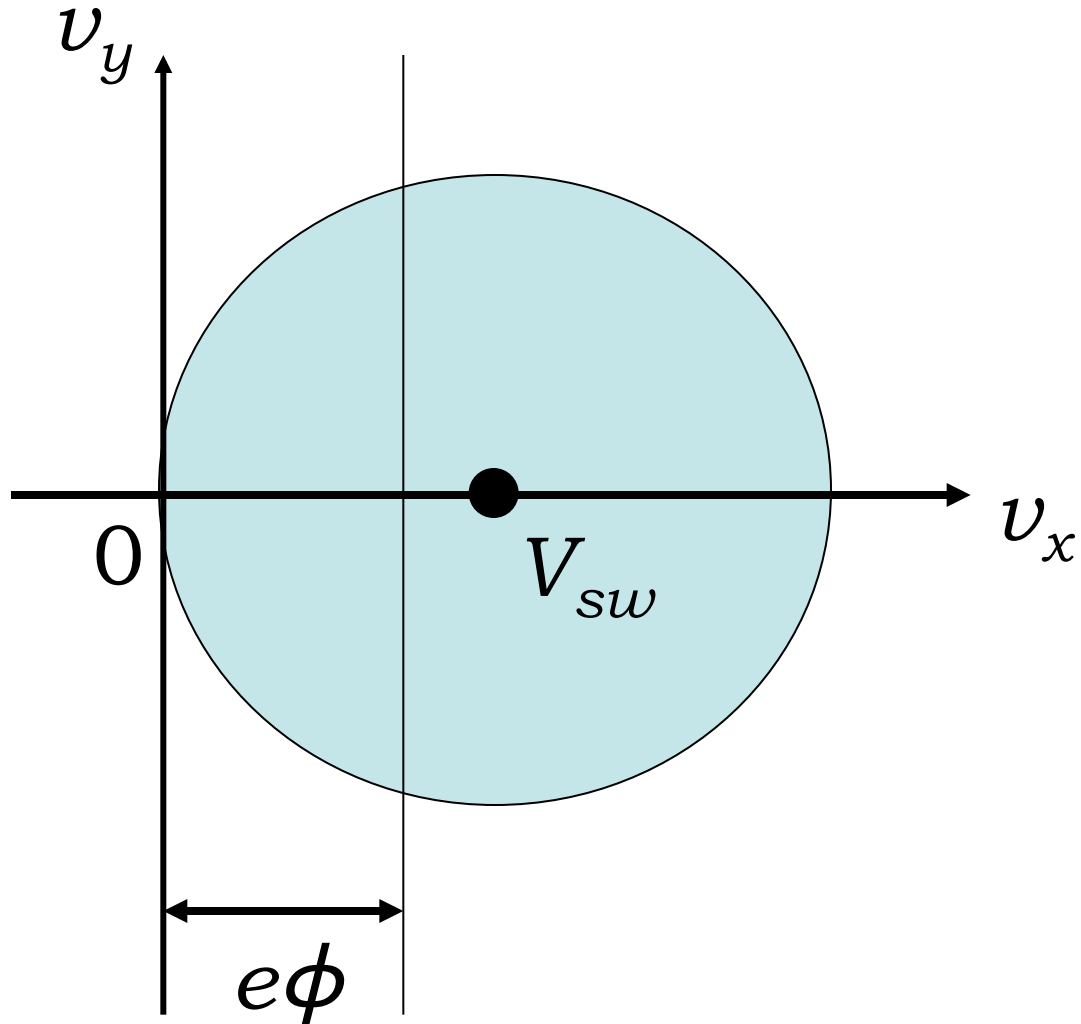
“Termination Shock” in Your Sink



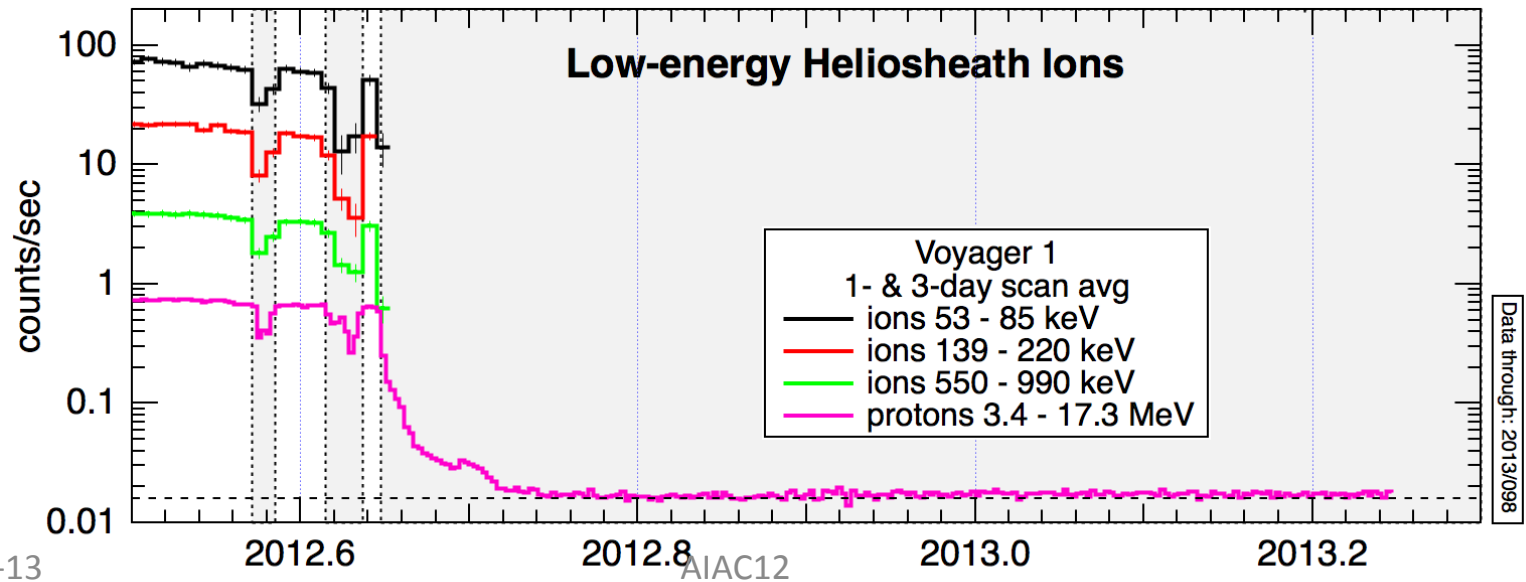
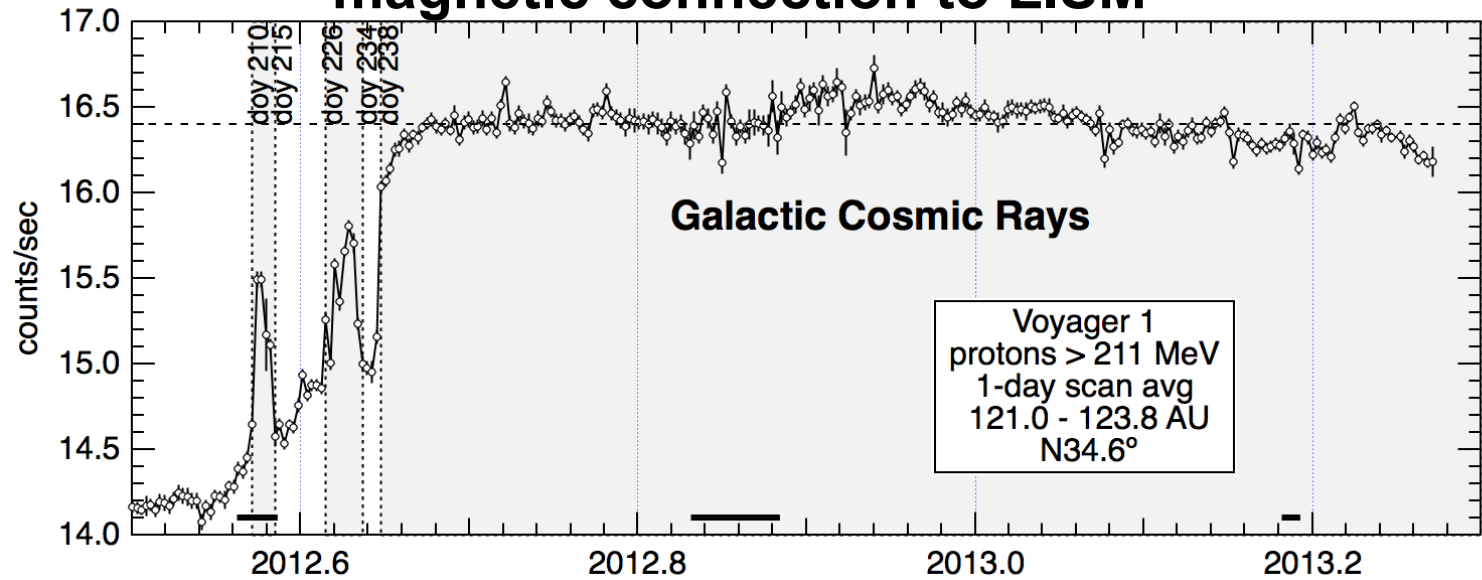
The Heliosphere



Pickup Ion Mediated Termination Shock

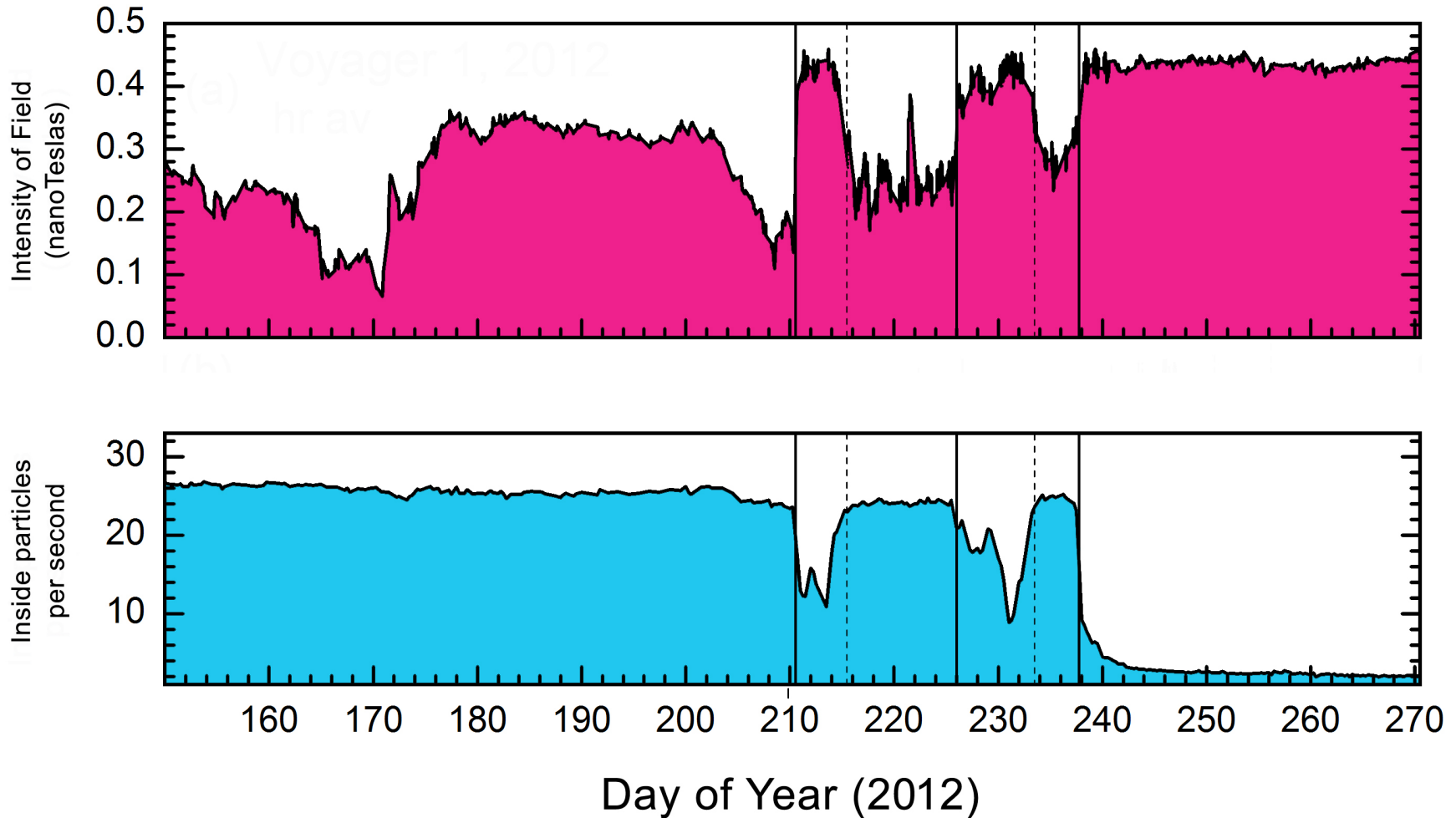


Voyager 1, 2012/183-2013/98: Galactic cosmic rays enter the high-B regions, while low-energy heliosheath particles leave via magnetic connection to LISM



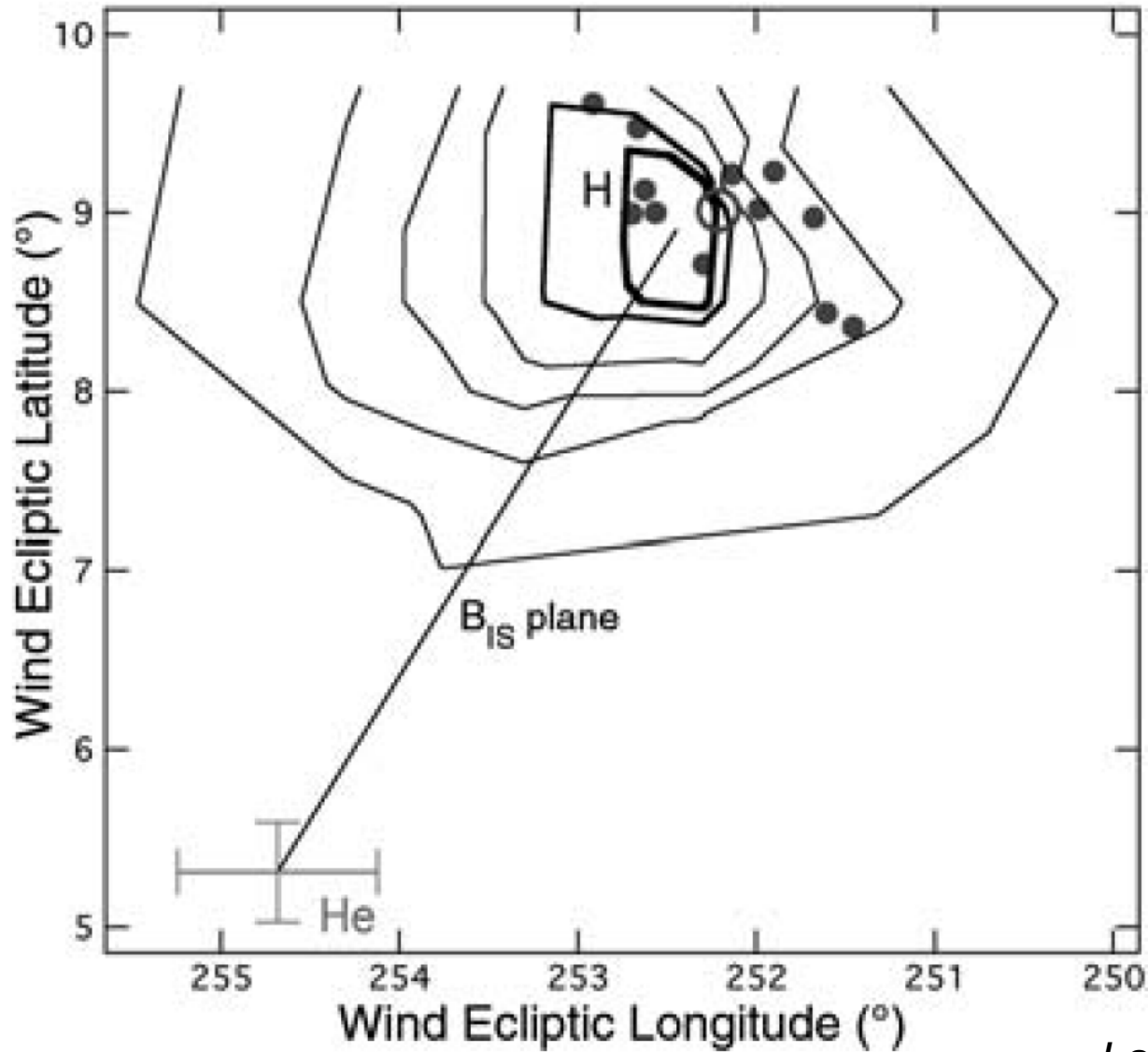
Voyager in the Heliosheath

Voyager 1 Magnetic Field and Charged Particles



Burlaga et al., 2013 Image credit: NASA/JPL-Caltech/GSFC/
University of Delaware

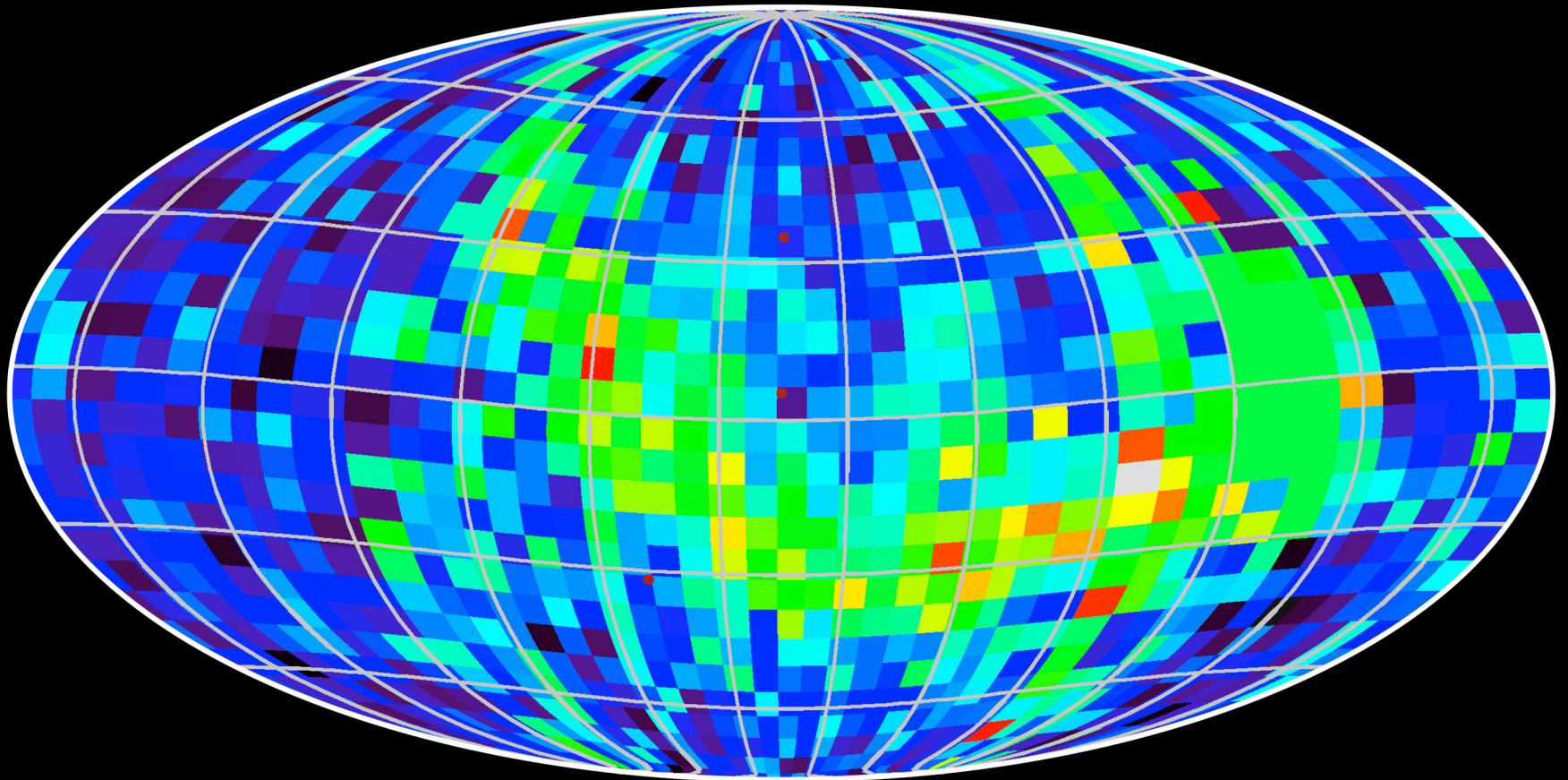
Lyman- α Backscatter



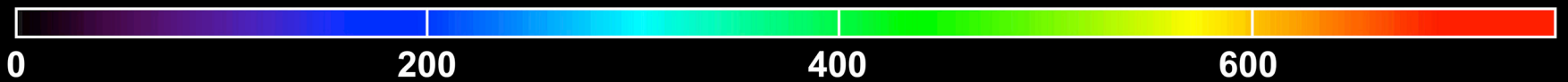
Lallement et al., 2005

The Termination of the Neutral Solar Wind

IBEX-Hi (0.6-1.0 keV)



Differential Flux [ENAs/(cm² s sr keV)]



Pickup Ions Beyond Heliosphere

